

Gravitational singularities and Holographic Complexity¹

Shubho Roy

(Indian Inst. of Technology Hyderabad)

BIRS workshop on “Quantum Information Theory in Quantum Field Theory and Cosmology”

June 9, 2023

¹w/ J. Ren (SY-S U.) & G. Katoch (IITH) (2303.02752 [hep-th])

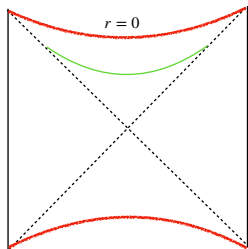
w/ E. Rabinovici (Racah) & S. Bolognesi (Pisa), 1802.02045 [hep-th]

Introduction

- ▶ Holography: Bulk geometry \equiv Boundary State Entanglement structure (Ryu-Takayanagi '06, Maldacena-Susskind '13 "*ER=EPR*", Raamsdonk '10)

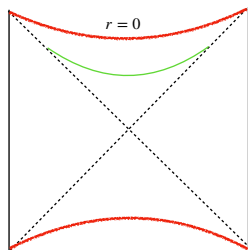
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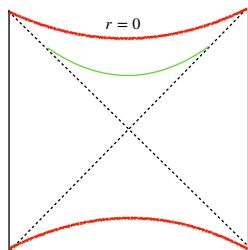
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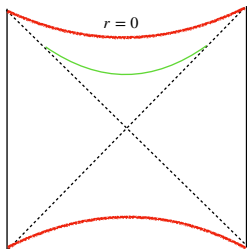
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- ▶ EAdS-BH at late times:

$$C \sim \text{"ERB volume"}; \quad \frac{dC}{dt} \sim T S,$$

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- ▶ Conclusions and Outlook

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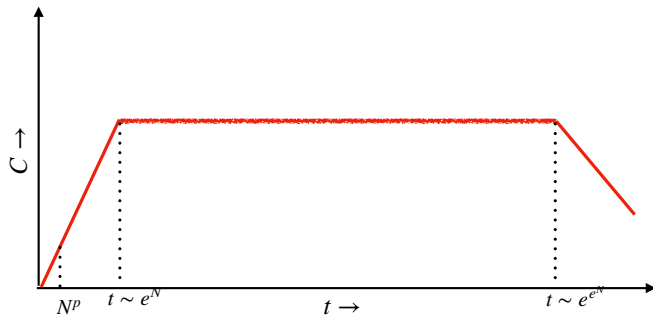
- ▶ Quant. mech., $|\psi\rangle = \sum_1^{2^N} \alpha_i |i\rangle$

$$C_{max} \sim 2^N !$$

Computational Complexity in Quantum systems

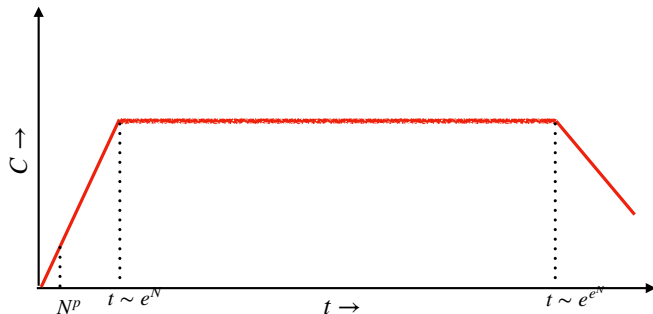
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► Time evolution of Complexity



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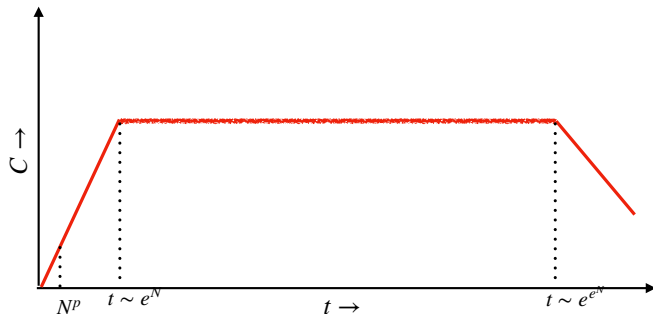
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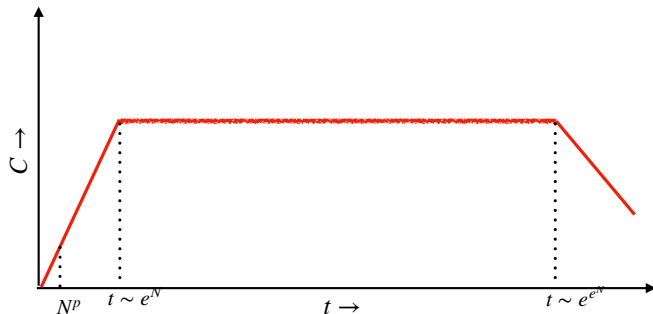
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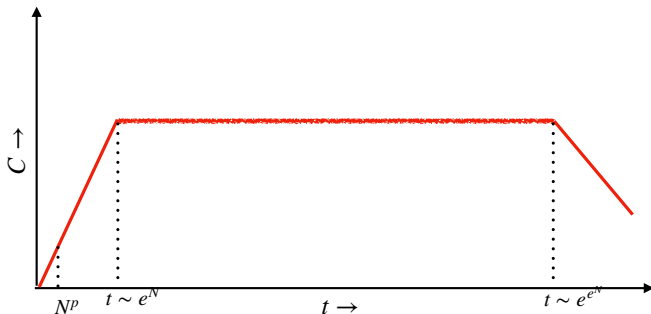
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- Complexity decrease by $t \sim \beta e^{e^N}$ (Poincaré recurrences)

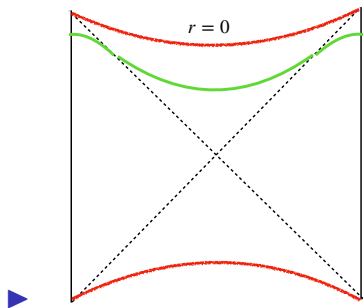
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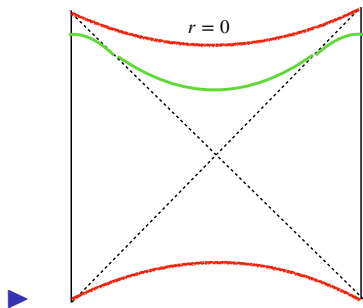


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- ▶ Initial Growth Slope: $\frac{dC}{dt} \sim TS$

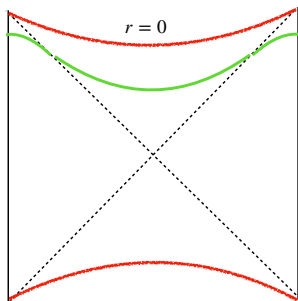
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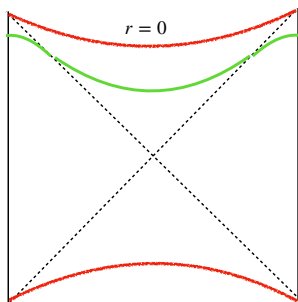
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$$C = \frac{\text{Vol}(\Sigma_{max})}{G_N R_c}$$

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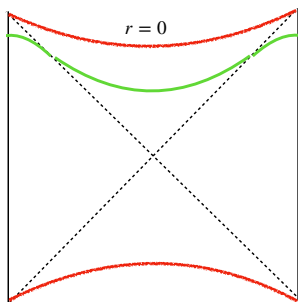


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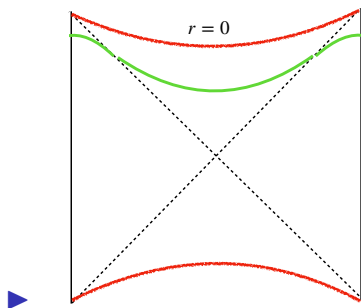


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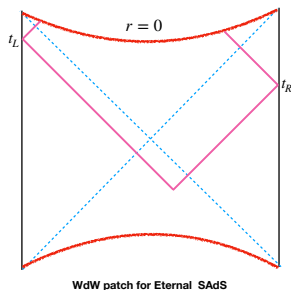
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- ▶ However, lesson from BH: lack of entanglement \Rightarrow singular spacetime (firewalls)

Holography: Complexity and WdW Action: CA

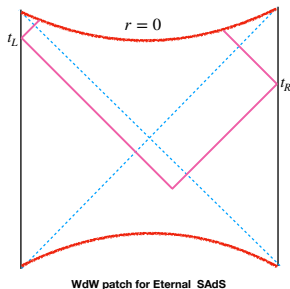
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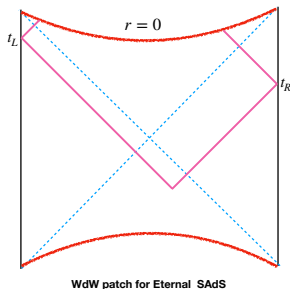


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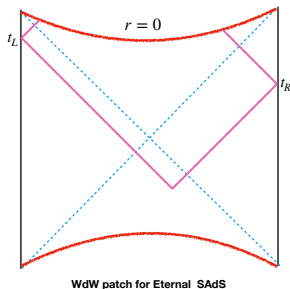


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- ▶ Still CV and CA matches perfectly!

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- ▶ Marginal: Coupling or boundary metric gains time-dependence (Kasner, Topological Crunch)

$$ds^2 = \frac{l^2}{z^2} (dz^2 - dt^2 + h_{ij}(t, x) dx^i dx^j), \quad i, j = 1, \dots, d$$

$$h_{ij}^K(t, x) = \text{diag} \left(\left(\frac{t}{l} \right)^{2p_1}, \dots, \left(\frac{t}{l} \right)^{2p_d} \right),$$

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- ▶ Relevant: Time dependent Mass scale, $M(t) = M \sec t$ (dS/Crunch)

$$ds_{bulk}^2 = d\rho^2 + f^2(\rho, M) ds_{dS_d}^2$$

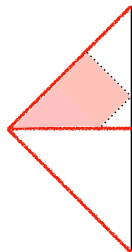
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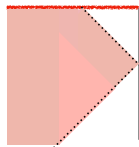
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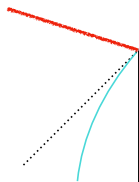
Cosmological Singularities in the bulk³



▶ *AdS Kasner*



Topological Crunch



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Complexity Estimates CV

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$$C(t) \sim N^2 \Lambda^d V_x \frac{|t|}{l} + \Lambda^{d-2} N^2 \frac{V_x}{l t}, \quad N^2 \sim \frac{l^d}{G_N}.$$

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$$C_\infty \sim N^2 V_{S^d} \Lambda^d \cos\left(\frac{t}{l}\right) + N^2 \frac{V_{S^d}}{l^2} \Lambda^{d-2} \frac{\sin^2 t/l}{\cos t/l}.$$

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$$C \sim N^2 V (\Lambda^{d-1} - M(t)^{d-1}) + N^2 l_- \Omega_{d-1} r(t)^{d-1}$$

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- ▶ Every case: Complexity decreases as we approach the singularity!

Complexity Estimates: CA

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- ▶ dS/Crunch: Subleading terms are also different!

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Complexity of Cosmological Crunches: Universal features

- ▶ Complexity Monotonically decreases due to loss of dof (CFT volume crunches)!
- ▶ Time rate of change of complexity contains a UV divergent time-dependent piece for CFT metric being time-dependent
- ▶ Coefficient of the rate of change determined by the subleading term (YGH term for $C \propto \mathcal{A}$).

Complexity of Cosmological Singularities: Takeaway

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- ▶ Perhaps two distinct bulk geometric constructions are two different CFT complexities as well
- ▶ Universal features for decrease of complexity, contrasts w/ local probes (point probes/strings - blue shifting)
- ▶ Perhaps one can attempt a parallel with the classic BKL work regarding universality

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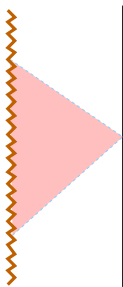
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- ▶ Solutions to effective holographic theories at zero temperature have typically naked *timelike* singularities
- ▶ Such singularities are generically resolved by lifting them to higher dimensions or eventually by the inclusion of the stringy states.
- ▶ Gubser criterion: *Naked singularities allowed in geometries are those which can be obtained as deformations/limits of regular black holes* [Gubser '01, Kiritsis et. al. '10,...]

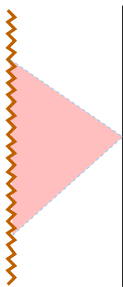
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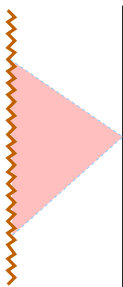


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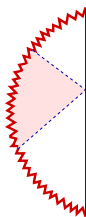
Overall action complexity (also \mathcal{C}_V) is less than empty global AdS! (criterion)

Timelike Kasner AdS

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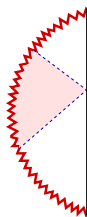
$$ds^2 = \frac{R^2}{z^2} \left(\frac{dz^2}{f(z)} - f^\alpha(z) dt^2 + f^\beta(z) dx^2 + f^\gamma(z) dy^2 \right), \quad f(z) = 1 - \frac{z^3}{z_0^3}$$



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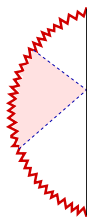
$$\mathcal{C}_A = \frac{l^2}{16\pi^2 G_N} \frac{V_{xy}}{\delta^2} - \frac{l^2}{32\pi G_N} \frac{V_{xy}}{z_0^2} \frac{(3-\alpha)\Gamma\left(\frac{1}{3}\right)\sec\left(\frac{\pi\alpha}{2}\right)}{\Gamma\left(\frac{5-3\alpha}{6}\right)\Gamma\left(\frac{\alpha+1}{2}\right)}.$$

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- ▶ **Action complexity lower than the empty (Poincaré) AdS: in sync with Gubser criterion**

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$$ds^2 = f(r)(-dt^2 + d\mathbf{x}^2) + \frac{dr^2}{f(r)}, \quad f(r) = r^2 \left(1 + \frac{b}{r}\right)^{\frac{2\delta^2}{1+\delta^2}} = e^{\delta\phi},$$

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- ▶ **Overall C_A is positive and larger than pure AdS for $\delta > 1/\sqrt{3}$. For $\delta < 1/\sqrt{3}$, C_A is negative and (IR) divergent! In sync Gubser criterion!**

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- ▶ Volume complexity not a reliable tool to probe timelike singularities.
- ▶ Need to conduct a more comprehensive survey of other nakedly timelike singular geometries in future to confirm \mathcal{C}_A criterion.