

Differentiability of shape functions  
for directed polymers  
in continuous space

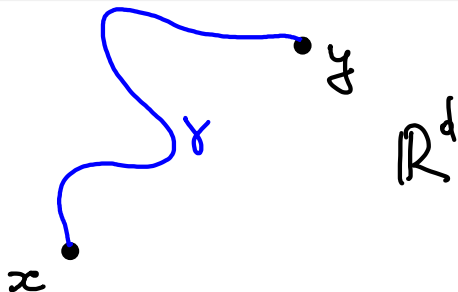
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joint work with Douglas Dow

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BIRS 2023

# Optimal paths in random environments (zero temperature)



$$A_\omega(x, y) = \inf \left\{ A_\omega(\gamma) \mid \text{admissible } \gamma : x \rightsquigarrow y \right\}$$

Shape function

$$\Lambda(v) = \lim_{T \rightarrow \infty} \frac{1}{T} A_\omega(0, Tv)$$

or

$$\Lambda(v) = \lim_{n \rightarrow \infty} \frac{1}{n} A_\omega(0, x_n), \quad \frac{x_n}{n} \rightarrow v.$$

## Exact shape functions (including the temperature $> 0$ case)

- Corner growth, i.i.d. exponential weights: Rost (1981)
- (generalized) Hammersley process: Hammersley (1972), Aldous, Diaconis (1995), Cator, Pimentel (2011)
- Euclidean FPP: Howard, Newman (1997)
- O'Connell–Yor polymers: Baryshnikov (2001), Gravner, Tracy, Widom (2001), Hambly, Martin, O'Connell (2002), Moriarty, O'Connell (2007)
- Log-gamma polymers: Seppäläinen (2012)
- Burgers equation, quadratic  $L$ : Bakhtin, Cator, Khanin (2014), Bakhtin (2016), Bakhtin, Li (2019)
- KPZ equation: Janjigian, Rassoul-Agha, Seppäläinen (recent)

# Shape function

- Always convex
- In all explicit examples, differentiable and strictly convex
- Strict convexity would imply existence-uniqueness of one-sided geodesics and infinite volume polymer measures (thermodynamic limits) with a given slope, and even just differentiability allows to make pretty strong claims in that direction [Janjigian, Rassoul-Agha, Seppäläinen 2020,2022]
- Differentiability at the edge of the percolation cone for lattice FPP, LPP [Auffinger, Damron 2013].

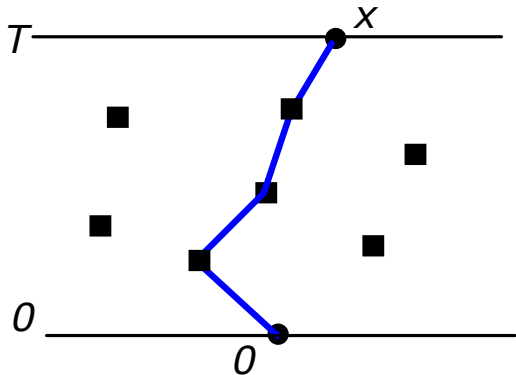
This talk:

Several classes of models where the shape function is not known precisely but differentiability holds:

- continuous space directed polymers
  - zero temperature
  - positive temperature
- homogenization in HJB equations with dynamic environments.

# Shear invariant case (Burgers/KPZ type models)

Poissonian points in space-time  $\mathbb{R} \times \mathbb{R}$ :



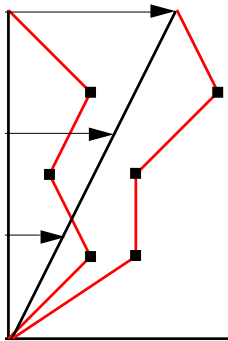
$$A^{0,T}(\gamma) = \int_0^T \dot{\gamma}(t)^2 dt - \#\{\text{Poisson points on } \gamma\}$$

[Bakhtin, Cator, Khanin 2014]

# Shear invariant case (Burgers/KPZ type models)

$$A^{0,T}(\gamma) = \int_0^T \dot{\gamma}(t)^2 dt - \#\{\text{Poisson points on } \gamma\}$$

If  $\gamma(0) = \gamma(T) = 0$  and  $\Xi_v \gamma(t) = \gamma(t) + vt$ , then



$$\begin{aligned} & \int_0^T \left( \frac{d}{dt} \Xi_v \gamma(t) \right)^2 dt \\ &= \int_0^T (\dot{\gamma}(t) + v)^2 dt \\ &= \int_0^T \dot{\gamma}(t)^2 dt + 2v \int_0^T \dot{\gamma}(t) dt + \int_0^T v^2 dt \\ &= \int_0^T \dot{\gamma}(t)^2 dt + Tv^2. \end{aligned}$$

PPP is distributionally invariant under  $\Xi_v$ , so  $\Lambda(v) = \Lambda(0) + v^2$ .

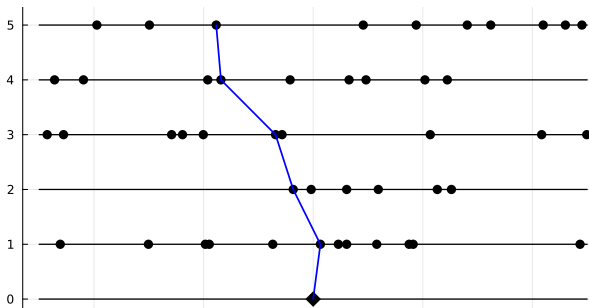
## Several other shear invariant models

$$A(\gamma) = \sum_k (\gamma_{k+1} - \gamma_k)^2 + \sum_k F_k(\gamma_k),$$

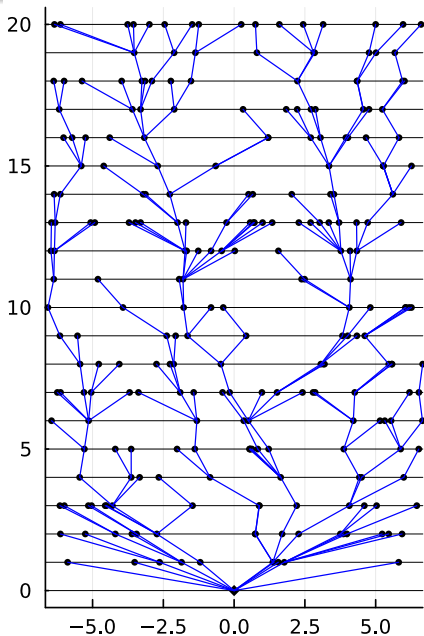
where  $F_k$  are i.i.d. in time and stationary in space [Bakhtin, 2016], or simply

$$A(\gamma) = \sum_k (\gamma_{k+1} - \gamma_k)^2$$

but require  $\gamma_k$  to coincide with one of the Poisson points on  $\{k\} \times \mathbb{R}$ .



# Trees of minimizers



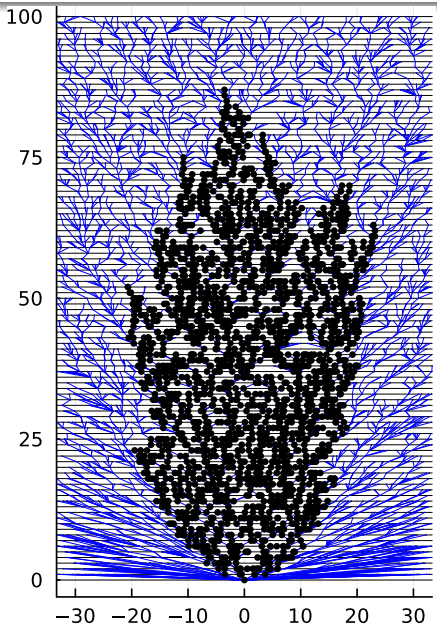
paths minimizing

$$A = \sum (\gamma_{k+1} - \gamma_k)^2$$

for various endpoints.



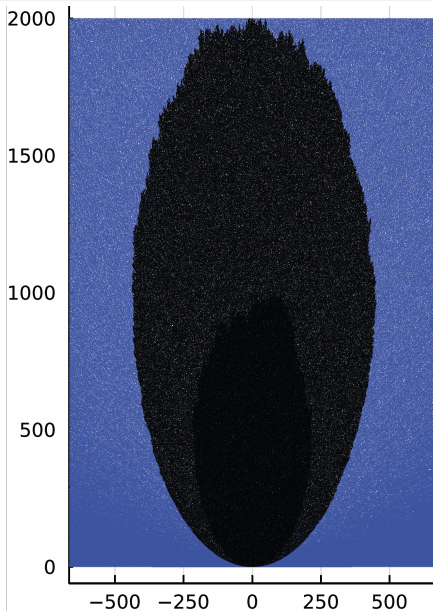
# Limit shape



A point is shown in black if there is a path to that point with

$$A = \sum_{i=1}^n (\gamma_i - \gamma_{i-1})^2 < \mathbf{20}$$

# Limit shape



A point is shown in black if there is a path to that point with

$$A = \sum_{i=1}^n (\gamma_i - \gamma_{i-1})^2 < \mathbf{400}$$

The boundary of the limit shape is an ellipse

$$x^2 + C \left( t - \frac{1}{2C} \right)^2 = \frac{1}{4C}$$

with  $C = \Lambda(0)$ .

# General (nonquadratic) action

$$A(\gamma) = \sum_k L(\Delta_k \gamma) + \sum_k F_k(\gamma_k) \quad (\Delta_k \gamma = \gamma_{k+1} - \gamma_k)$$

**Theorem** [with Douglas Dow] Assume that

- $F$  is i.i.d. in time, stationary in space, continuous, bounded from below,  $\mathbb{E} F(0) < \infty$ ;
- $L \in C^2$ ,  $\lim_{|v| \rightarrow \infty} L(v) = +\infty$ ,  $\limsup_{|v| \rightarrow \infty} \frac{L''(v)}{L(v)} < \infty$ .

(Doesn't have to be convex, e.g.  $L(v) = v^{2p} + \sum_{k=0}^{2p-1} a_k v^k$ )

Then there is a deterministic, convex, and **differentiable** shape function  $\Lambda$ : for each  $v \in \mathbb{R}$ , with probability 1,

$$\Lambda(v) = \lim_{n \rightarrow \infty} \frac{1}{n} A^n(0, nv),$$

$$\Lambda'(v) = \lim_{n \rightarrow \infty} \frac{1}{n} \sum_{k=0}^{n-1} L'(\Delta_k \gamma^n(v)),$$

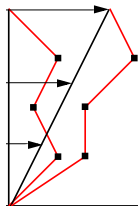
where  $\gamma^n(v)$  realizes  $A^n(0, nv) = \inf\{A(\gamma) \mid \gamma_0 = 0, \gamma_n = nv\}$ .

Using shear  $\Xi_v(t, x) = (t, x + tv)$

For a path  $\gamma$  with  $\gamma_0 = \gamma_n = 0$

$$B(v, \gamma) = \sum_k L(\Delta_k \gamma + v) + \sum_k F_k(\gamma_k)$$

$$B^n(v) = \inf\{B(v, \gamma) \mid \gamma_0 = \gamma_n = 0\}$$



Since  $\Xi_v F = \Xi_{-v} F = F$  in distribution,

$$(B^n(v))_{n \in \mathbb{N}} \stackrel{d}{=} (A^n(0, nv))_{n \in \mathbb{N}},$$

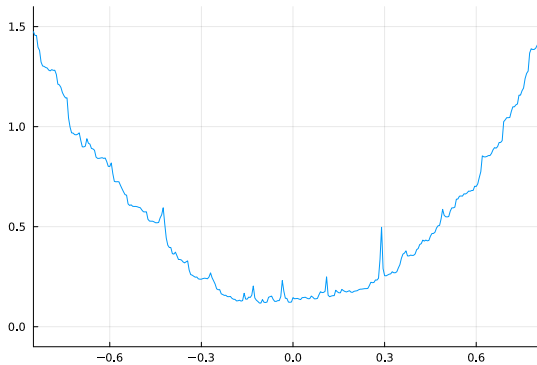
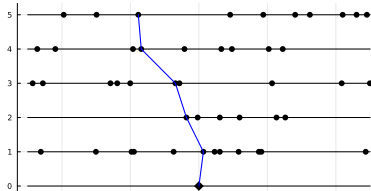
so

$$\Lambda(v) = \lim_{n \rightarrow \infty} \frac{1}{n} A^n(0, nv) = \lim_{n \rightarrow \infty} \frac{1}{n} B^n(v).$$

In addition,  $(B^n(v))_{v \in \mathbb{R}}$  is “nicer” than  $(A^n(0, nv))_{v \in \mathbb{R}}$

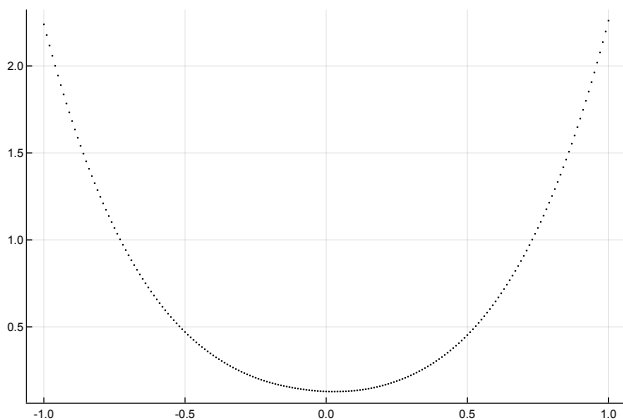
# Poisson points model with $A(\gamma) = \sum_k (\Delta_k \gamma)^4$

$$\frac{1}{n} A^n(0, nv) \text{ for } n = 200$$



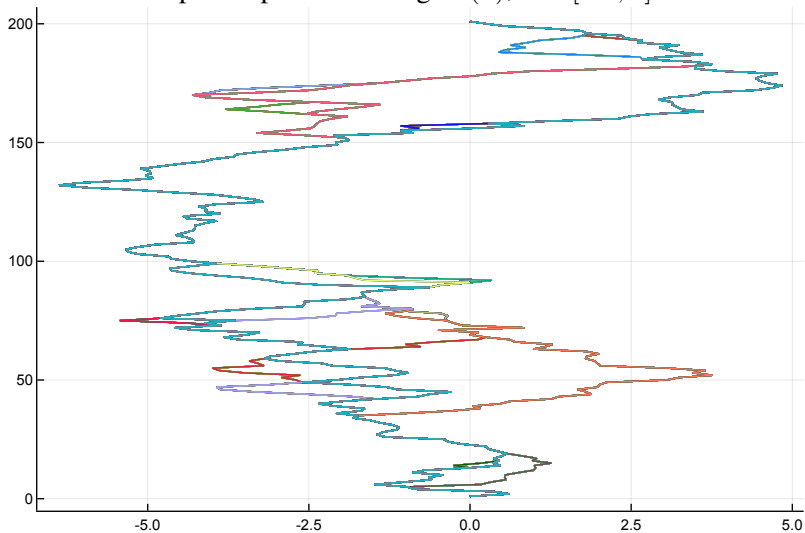
# Poisson points model with $A(\gamma) = \sum_k (\Delta_k \gamma)^4$

$$\frac{1}{n} B^n(v) \text{ for } n = 200$$



# Poisson points model with $A(\gamma) = \sum_k (\Delta_k \gamma)^4$

Optimal paths realizing  $B^n(v)$ ,  $v \in [-1, 1]$ .



Each path serves a range of values of  $v$ .

# Proof of differentiability

Use  $\gamma^n(v)$ , the optimal path for  $B^n(v)$  to estimate  $B^n(w)$  for  $w \approx v$ .

$$B^n(w) \leq B^n(w, \gamma^n(v))$$

$$\begin{aligned} &\leq B^n(v) + (w - v) \sum_{k=0}^{n-1} L'(\Delta_k \gamma(v) + v) \\ &\quad + \frac{(w - v)^2}{2} \sum_{k=0}^{n-1} L''(\Delta_k \gamma(v) + v + s(w - v)) \end{aligned}$$

For  $w - v > 0$ ,

$$\frac{B^n(w) - B^n(v)}{n(w - v)} \leq \frac{1}{n} \sum_{k=0}^{n-1} L'(\Delta_k \gamma(v) + v) + C(w - v)$$

$$\frac{\Lambda(w) - \Lambda(v)}{w - v} \leq \liminf_{n \rightarrow \infty} \frac{1}{n} \sum_{k=0}^{n-1} L'(\Delta_k \gamma(v) + v) + C(w - v)$$

$$\partial^+ \Lambda(v) \leq \liminf_{n \rightarrow \infty} \frac{1}{n} \sum_{k=0}^{n-1} L'(\Delta_k \gamma(v) + v).$$



# Proof of differentiability

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For  $w - v < 0$ ,

$$\frac{B^n(w) - B^n(v)}{n(w - v)} \geq \frac{1}{n} \sum_{k=0}^{n-1} L'(\Delta_k \gamma(v) + v) + C|w - v|$$

$$\frac{\Lambda(w) - \Lambda(v)}{w - v} \geq \limsup_{n \rightarrow \infty} \frac{1}{n} \sum_{k=0}^{n-1} L'(\Delta_k \gamma(v) + v) + C|w - v|$$

$$\partial^- \Lambda(v) \geq \limsup_{n \rightarrow \infty} \frac{1}{n} \sum_{k=0}^{n-1} L'(\Delta_k \gamma(v) + v).$$

# Proof of differentiability

So

$$\begin{aligned}\partial^+ \Lambda(v) &\leq \liminf_{n \rightarrow \infty} \frac{1}{n} \sum_{k=0}^{n-1} L'(\Delta_k \gamma(v) + v) \\ &\leq \limsup_{n \rightarrow \infty} \frac{1}{n} \sum_{k=0}^{n-1} L'(\Delta_k \gamma(v) + v) \leq \partial^- \Lambda(v).\end{aligned}$$

But  $\Lambda$  is convex, so

$$\partial^- \Lambda(v) \leq \partial^+ \Lambda(v).$$

Therefore,

$$\partial^- \Lambda(v) = \partial^+ \Lambda(v) = \Lambda'(v) = \lim_{n \rightarrow \infty} \frac{1}{n} \sum_{k=0}^{n-1} L'(\Delta_k \gamma(v) + v)$$

In terms of the optimal path from  $(0, 0)$  to  $(n, nv)$ :

$$\Lambda'(v) = \lim_{n \rightarrow \infty} \frac{1}{n} \sum_{k=0}^{n-1} L'(\Delta_k \gamma(v))$$

## Doesn't quite imply strict convexity

$$\Lambda'(v) = \lim_{n \rightarrow \infty} \frac{1}{n} \sum_{k=0}^{n-1} L'(\Delta_k \gamma(v)) = \langle L'(\Delta \gamma(v)) \rangle$$

For example, if  $L(v) = v^4$ , then  $L'(v) = 4v^3$ ,

$$\Lambda'(v) = 4 \langle (\Delta \gamma(v))^3 \rangle.$$

It is natural to expect this to strictly grow in  $v$  (strict convexity) because

$$\langle \Delta \gamma(v) \rangle = v.$$

But the first moment does not control the third moment, so this is not obvious.

## Positive temperature case. Average free energy.

$$\begin{aligned} Z^n(y) &= \int \exp \left[ -A(\gamma) \right] \delta_0(d\gamma_0) d\gamma_1 \dots d\gamma_{n-1} \delta_y(d\gamma_n) \\ &= \int \exp \left[ -\sum_{k=0}^{n-1} L(\Delta_k \gamma) - \sum_{k=0}^{n-1} F_k(\gamma_k) \right] \delta_0(d\gamma_0) d\gamma_1 \dots d\gamma_{n-1} \delta_y(d\gamma_n) \end{aligned}$$

Need more requirements:  $\mathbb{E} \sup_{|x| \leq 1/2} F_k(x) < \infty$ ,

$$\liminf_{|v| \rightarrow \infty} |L'(v)| > 0, \quad \limsup_{|v| \rightarrow \infty} \frac{|L'(v)|}{|L(v)|^\theta} < \infty \quad \text{for some } \theta \in (0, 1).$$

**Theorem** [with Douglas Dow] There is a deterministic, convex, differentiable  $\Lambda : \mathbb{R} \rightarrow \mathbb{R}$  s.t. for every  $v \in \mathbb{R}$ , with prob. 1,

$$\begin{aligned} \Lambda(v) &= -\lim_{n \rightarrow \infty} \frac{1}{n} \log Z^n(nv), \\ \Lambda'(v) &= \lim_{n \rightarrow \infty} \mu_{nv}^n \left( \frac{1}{n} \sum_{k=0}^{n-1} L'(\Delta_k \gamma) \right) \end{aligned}$$

(  $\mu_{nv}^n$  is the polymer measure on paths connecting  $(0, 0)$  to  $(n, nv)$  )

## The main estimate (after applying the shear)

All paths connect 0 to 0 (not 0 to  $Tv$ )

$$\begin{aligned} \log \tilde{Z}^n(w) &= \log \int \exp \left( - \sum_k \left[ F_k(\gamma_k) + L(\Delta_k \gamma + w) \right] \right) d\gamma \\ &\geq \log \int e^{-\sum_k \left[ F_k(\gamma_k) + L(\Delta_k \gamma + v) + (w-v)L'(\Delta_k \gamma + v) + \frac{1}{2}(w-v)^2 L''(\dots) \right]} d\gamma \\ &= \log \tilde{Z}^n(v) \\ &\quad + \log \frac{1}{\tilde{Z}^n(v)} \int e^{-\sum_k \left[ F_k(\gamma_k) + L(\Delta_k \gamma + v) + (w-v)L'(\Delta_k \gamma + v) + \frac{1}{2}(w-v)^2 L''(\dots) \right]} d\gamma \\ &= \log \tilde{Z}^n(v) + \log \tilde{\mu}_v^n \left( e^{-\sum_k \left[ (w-v)L'(\Delta_k \gamma + v) + \frac{1}{2}(w-v)^2 L''(\dots) \right]} \right) \\ &\geq \log \tilde{Z}^n(v) - (w-v) \tilde{\mu}_v^n \left( \sum_{k=0}^{n-1} L'(\Delta_k \gamma + v) \right) - \frac{1}{2} (w-v)^2 \tilde{\mu}_v^n \left( \dots \right) \end{aligned}$$

# Continuous time, HJB eqns, non-white noise, in $\mathbb{R}^d$ , $d \geq 1$

$$A^t(\gamma) = \int_0^t L(\dot{\gamma}_s) ds + \int_0^t F(s, \gamma_s) ds$$

- We no longer assume  $F$  is white in time:

$$F(t, x) = \sum_i \varphi_i(t-t_i, x-x_i) = \int_{\mathbb{R} \times \mathbb{R}^d \times \mathcal{C}} \varphi(t-s, x-y) \mathbf{N}(ds, dy, d\varphi)$$

$(t_i, x_i)$  are Poisson points in  $\mathbb{R} \times \mathbb{R}^d$ , each convolved with its own random  $\varphi_i$  (i.i.d.,  $C^2$ , uniformly bounded support, an exp-moment)

- $L : \mathbb{R}^d \rightarrow \mathbb{R}$  convex, twice differentiable,

$$\lim_{|v| \rightarrow \infty} \frac{L(v)}{|v|} = +\infty$$

$$\limsup_{|v| \rightarrow \infty} \sup_{|r| \leq \delta} \frac{\|\nabla^2 L(v+r)\|}{L(v)} < \infty \quad \text{for some } \delta > 0.$$

# Continuous time, HJB eqns, non-white noise, in $\mathbb{R}^d$ , $d \geq 1$

$$A^t(\gamma) = \int_0^t L(\dot{\gamma}_s) ds + \int_0^t F(s, \gamma_s) ds$$

$$A(t, x) = \inf \{A^t(\gamma) : \gamma(0) = 0, \gamma(t) = x\}$$

$$\partial_t A(t, x) + H(\nabla A(t, x)) = F(t, x), \quad t \in (0, \infty), x \in \mathbb{R}^d$$

The Hamiltonian  $H$  is the Legendre–Fenchel transform of  $L$

$$H(p) = \sup_{v \in \mathbb{R}^d} \{\langle p, v \rangle - L(v)\}, \quad p \in \mathbb{R}^d$$

$$\lim_{t \searrow 0} A(t, x) = \begin{cases} 0, & x = 0, \\ +\infty, & x \neq 0. \end{cases}$$

# Shape Theorem

**Theorem** [with Douglas Dow]

Under these conditions, there is a convex deterministic and **differentiable** function  $\Lambda : \mathbb{R}^d \rightarrow \mathbb{R}$  such that for every  $v \in \mathbb{R}^d$ , with probability 1,

$$\Lambda(v) = \lim_{t \rightarrow \infty} \frac{1}{T} A(T, Tv)$$

$$\nabla \Lambda(v) = \lim_{T \rightarrow \infty} \frac{1}{T} \int_0^T \left[ \nabla L(\dot{\gamma}_t^T(v)) + \Theta(t, \gamma_t^T(v)) \right] dt,$$

where

$$\begin{aligned} \Theta(t, x) &= \int (t-s) \nabla \varphi(t-s, x-y) \mathbf{N}(ds, dy, d\varphi) \\ &= \partial_v \int \varphi(t-s, x-y+v(t-s)) \mathbf{N}(ds, dy, d\varphi) \Big|_{v=0} \end{aligned}$$



# Shape Theorem. HJB Homogenization Version.

## Corollary

For every  $t \in (0, \infty)$  and  $x \in \mathbb{R}^d$ , with probability 1,

$$\lim_{\epsilon \searrow 0} \epsilon A(t/\epsilon, x/\epsilon) = t\Lambda(x/t)$$

The nonrandom function  $\bar{U}(t, x) = t\Lambda(x/t)$  is the fundamental viscosity solution of the deterministic HJB equation

$$\partial_t \bar{U}(t, x) + \bar{H}(\nabla \bar{U}(t, x)) = 0,$$

where  $\bar{H}$  is the Legendre–Fenchel transform of  $\Lambda$ :

$$\bar{H}(p) = \sup_{v \in \mathbb{R}^d} \{\langle v, p \rangle - \Lambda(v)\}, \quad p \in \mathbb{R}^d.$$

Moreover,  $\bar{H}$  is **strictly convex** (no flat edges), and  $\bar{U}(t, x)$  is a **classical** solution which is  $C^1$  for all  $t > 0$ ,  $x \in \mathbb{R}^d$ .

# Existing homogenization results for dynamic environments

## Zero viscosity HJB:

- Schwab 2009
- Bakhtin, Cator, Khanin 2014, Bakhtin 2016 (for quadratic  $L$  in the context of Burgers equation),
- Seeger 2021

## Positive viscosity HJB:

- Kosygina, Varadhan 2008
- Jing, Souganidis, Tran 2017
- Bakhtin, Li 2019 (Burgers equation, quadratic  $L$ )
- Janjigian, Rassoul-Agha, Seppäläinen (KPZ eqn, quadratic  $L$ )

## Further questions

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- continuous FPP with asymmetries
- (parabolic) HJB with positive viscosity/diffusion
- More general potentials (we still need shear invariance for the background process)?
- Does this say anything about lattice models?
- Strict convexity? Use our formulas for  $\nabla\Lambda$ ?