

# On non-topological solutions for planar Liouville Systems of Toda-type

Arkady Poliakovsky

Ben-Gurion University of the Negev

# A class of systems arising from the study of vortex configurations in self-dual gauge field theories

# A class of systems arising from the study of vortex configurations in self-dual gauge field theories

- Consider the problem:

$$\begin{cases} -\Delta w_i = \sum_{j=1}^m a_{ij} e^{w_j} - 4\pi N_i \delta_0 & \text{in } \mathbb{R}^2 \\ \int_{\mathbb{R}^2} e^{w_i} < \infty. \end{cases} \quad (1)$$

# A class of systems arising from the study of vortex configurations in self-dual gauge field theories

- Consider the problem:

$$\begin{cases} -\Delta w_i = \sum_{j=1}^m a_{ij} e^{w_j} - 4\pi N_i \delta_0 & \text{in } \mathbb{R}^2 \\ \int_{\mathbb{R}^2} e^{w_i} < \infty. \end{cases} \quad (1)$$

- Here  $N_j > -1$ ,  $j = 1, \dots, m$  and  $A = \{a_{ij}\}$  is a symmetric matrix.

# A class of systems arising from the study of vortex configurations in self-dual gauge field theories

- Consider the problem:

$$\begin{cases} -\Delta w_i = \sum_{j=1}^m a_{ij} e^{w_j} - 4\pi N_i \delta_0 & \text{in } \mathbb{R}^2 \\ \int_{\mathbb{R}^2} e^{w_i} < \infty. \end{cases} \quad (1)$$

- Here  $N_j > -1$ ,  $j = 1, \dots, m$  and  $A = \{a_{ij}\}$  is a symmetric matrix.
- By setting:  $w_i(x) = u_i(x) + 2N_i \ln |x|$ , we reduce (1) to:

$$\begin{cases} -\Delta u_i = \sum_{j=1}^m a_{ij} |x|^{2N_j} e^{u_j} & \text{in } \mathbb{R}^2 \\ \frac{1}{2\pi} \int_{\mathbb{R}^2} |x|^{2N_i} e^{u_i} dx < +\infty \end{cases} \quad (2)$$

# A class of systems arising from the study of vortex configurations in self-dual gauge field theories

- Consider the problem:

$$\begin{cases} -\Delta w_i = \sum_{j=1}^m a_{ij} e^{w_j} - 4\pi N_i \delta_0 & \text{in } \mathbb{R}^2 \\ \int_{\mathbb{R}^2} e^{w_i} < \infty. \end{cases} \quad (1)$$

- Here  $N_j > -1$ ,  $j = 1, \dots, m$  and  $A = \{a_{ij}\}$  is a symmetric matrix.
- By setting:  $w_i(x) = u_i(x) + 2N_i \ln |x|$ , we reduce (1) to:

$$\begin{cases} -\Delta u_i = \sum_{j=1}^m a_{ij} |x|^{2N_j} e^{u_j} & \text{in } \mathbb{R}^2 \\ \frac{1}{2\pi} \int_{\mathbb{R}^2} |x|^{2N_i} e^{u_i} dx < +\infty \end{cases} \quad (2)$$

- Clearly If  $u_i$  solves (2) and for every  $R > 0$  we define

# A class of systems arising from the study of vortex configurations in self-dual gauge field theories

- Consider the problem:

$$\begin{cases} -\Delta w_i = \sum_{j=1}^m a_{ij} e^{w_j} - 4\pi N_i \delta_0 & \text{in } \mathbb{R}^2 \\ \int_{\mathbb{R}^2} e^{w_i} < \infty. \end{cases} \quad (1)$$

- Here  $N_j > -1$ ,  $j = 1, \dots, m$  and  $A = \{a_{ij}\}$  is a symmetric matrix.
- By setting:  $w_i(x) = u_i(x) + 2N_i \ln |x|$ , we reduce (1) to:

$$\begin{cases} -\Delta u_i = \sum_{j=1}^m a_{ij} |x|^{2N_j} e^{u_j} & \text{in } \mathbb{R}^2 \\ \frac{1}{2\pi} \int_{\mathbb{R}^2} |x|^{2N_i} e^{u_i} dx < +\infty \end{cases} \quad (2)$$

- Clearly If  $u_i$  solves (2) and for every  $R > 0$  we define

$$u_i^{(R)}(x) := u_i(x/R) - 2(N_i + 1) \ln(R) \quad \forall i = 1, \dots, m, \quad (3)$$

# A class of systems arising from the study of vortex configurations in self-dual gauge field theories

- Consider the problem:

$$\begin{cases} -\Delta w_i = \sum_{j=1}^m a_{ij} e^{w_j} - 4\pi N_i \delta_0 & \text{in } \mathbb{R}^2 \\ \int_{\mathbb{R}^2} e^{w_i} < \infty. \end{cases} \quad (1)$$

- Here  $N_j > -1$ ,  $j = 1, \dots, m$  and  $A = \{a_{ij}\}$  is a symmetric matrix.
- By setting:  $w_i(x) = u_i(x) + 2N_i \ln |x|$ , we reduce (1) to:

$$\begin{cases} -\Delta u_i = \sum_{j=1}^m a_{ij} |x|^{2N_j} e^{u_j} & \text{in } \mathbb{R}^2 \\ \frac{1}{2\pi} \int_{\mathbb{R}^2} |x|^{2N_i} e^{u_i} dx < +\infty \end{cases} \quad (2)$$

- Clearly If  $u_i$  solves (2) and for every  $R > 0$  we define

$$u_i^{(R)}(x) := u_i(x/R) - 2(N_i + 1) \ln(R) \quad \forall i = 1, \dots, m, \quad (3)$$

then  $u_i^{(R)}$  also solves (2) and moreover



# A class of systems arising from the study of vortex configurations in self-dual gauge field theories

- Consider the problem:

$$\begin{cases} -\Delta w_i = \sum_{j=1}^m a_{ij} e^{w_j} - 4\pi N_i \delta_0 & \text{in } \mathbb{R}^2 \\ \int_{\mathbb{R}^2} e^{w_i} < \infty. \end{cases} \quad (1)$$

- Here  $N_j > -1$ ,  $j = 1, \dots, m$  and  $A = \{a_{ij}\}$  is a symmetric matrix.
- By setting:  $w_i(x) = u_i(x) + 2N_i \ln |x|$ , we reduce (1) to:

$$\begin{cases} -\Delta u_i = \sum_{j=1}^m a_{ij} |x|^{2N_j} e^{u_j} & \text{in } \mathbb{R}^2 \\ \frac{1}{2\pi} \int_{\mathbb{R}^2} |x|^{2N_i} e^{u_i} dx < +\infty \end{cases} \quad (2)$$

- Clearly If  $u_i$  solves (2) and for every  $R > 0$  we define

$$u_i^{(R)}(x) := u_i(x/R) - 2(N_i + 1) \ln(R) \quad \forall i = 1, \dots, m, \quad (3)$$

then  $u_i^{(R)}$  also solves (2) and moreover

$$\frac{1}{2\pi} \int_{\mathbb{R}^2} |x|^{2N_i} e^{u_i^{(R)}} dx = \frac{1}{2\pi} \int_{\mathbb{R}^2} |x|^{2N_i} e^{u_i} dx. \quad (4)$$

- We focus on the radial solvability of the problem

$$\begin{cases} -\Delta u_i = \sum_{j=1}^m a_{ij} |x|^{2N_j} e^{u_j} & \text{in } \mathbb{R}^2 \\ \frac{1}{2\pi} \int_{\mathbb{R}^2} |x|^{2N_i} e^{u_i} dx = \beta_i. \end{cases} \quad (5)$$

- We focus on the radial solvability of the problem

$$\begin{cases} -\Delta u_i = \sum_{j=1}^m a_{ij} |x|^{2N_j} e^{u_j} & \text{in } \mathbb{R}^2 \\ \frac{1}{2\pi} \int_{\mathbb{R}^2} |x|^{2N_i} e^{u_i} dx = \beta_i. \end{cases} \quad (5)$$

- It is well known that any solution to (5) satisfies Pohozaev identity:

- We focus on the radial solvability of the problem

$$\begin{cases} -\Delta u_i = \sum_{j=1}^m a_{ij} |x|^{2N_j} e^{u_j} & \text{in } \mathbb{R}^2 \\ \frac{1}{2\pi} \int_{\mathbb{R}^2} |x|^{2N_i} e^{u_i} dx = \beta_i. \end{cases} \quad (5)$$

- It is well known that any solution to (5) satisfies Pohozaev identity:

$$\sum_{i=1}^m \sum_{j=1}^m \frac{1}{2} a_{ij} \beta_i \beta_j - \sum_{k=1}^m 2(N_k + 1) \beta_k = 0. \quad (6)$$

- We focus on the radial solvability of the problem

$$\begin{cases} -\Delta u_i = \sum_{j=1}^m a_{ij} |x|^{2N_j} e^{u_j} & \text{in } \mathbb{R}^2 \\ \frac{1}{2\pi} \int_{\mathbb{R}^2} |x|^{2N_i} e^{u_i} dx = \beta_i. \end{cases} \quad (5)$$

- It is well known that any solution to (5) satisfies Pohozaev identity:

$$\sum_{i=1}^m \sum_{j=1}^m \frac{1}{2} a_{ij} \beta_i \beta_j - \sum_{k=1}^m 2(N_k + 1) \beta_k = 0. \quad (6)$$

So (6) is one of the necessary conditions of solvability of problem (5).

- We focus on the radial solvability of the problem

$$\begin{cases} -\Delta u_i = \sum_{j=1}^m a_{ij} |x|^{2N_j} e^{u_j} & \text{in } \mathbb{R}^2 \\ \frac{1}{2\pi} \int_{\mathbb{R}^2} |x|^{2N_i} e^{u_i} dx = \beta_i. \end{cases} \quad (5)$$

- It is well known that any solution to (5) satisfies Pohozaev identity:

$$\sum_{i=1}^m \sum_{j=1}^m \frac{1}{2} a_{ij} \beta_i \beta_j - \sum_{k=1}^m 2(N_k + 1) \beta_k = 0. \quad (6)$$

So (6) is one of the necessary conditions of solvability of problem (5).

- Moreover, the following condition, arising from the finiteness of integrals in (5):

- We focus on the radial solvability of the problem

$$\begin{cases} -\Delta u_i = \sum_{j=1}^m a_{ij} |x|^{2N_j} e^{u_j} & \text{in } \mathbb{R}^2 \\ \frac{1}{2\pi} \int_{\mathbb{R}^2} |x|^{2N_i} e^{u_i} dx = \beta_i. \end{cases} \quad (5)$$

- It is well known that any solution to (5) satisfies Pohozaev identity:

$$\sum_{i=1}^m \sum_{j=1}^m \frac{1}{2} a_{ij} \beta_i \beta_j - \sum_{k=1}^m 2(N_k + 1) \beta_k = 0. \quad (6)$$

So (6) is one of the necessary conditions of solvability of problem (5).

- Moreover, the following condition, arising from the finiteness of integrals in (5):

$$\sum_{j=1}^m a_{ij} \beta_j > 2(N_i + 1) \quad \forall i = 1, \dots, m, \quad (7)$$

- We focus on the radial solvability of the problem

$$\begin{cases} -\Delta u_i = \sum_{j=1}^m a_{ij} |x|^{2N_j} e^{u_j} & \text{in } \mathbb{R}^2 \\ \frac{1}{2\pi} \int_{\mathbb{R}^2} |x|^{2N_i} e^{u_i} dx = \beta_i. \end{cases} \quad (5)$$

- It is well known that any solution to (5) satisfies Pohozaev identity:

$$\sum_{i=1}^m \sum_{j=1}^m \frac{1}{2} a_{ij} \beta_i \beta_j - \sum_{k=1}^m 2(N_k + 1) \beta_k = 0. \quad (6)$$

So (6) is one of the necessary conditions of solvability of problem (5).

- Moreover, the following condition, arising from the finiteness of integrals in (5):

$$\sum_{j=1}^m a_{ij} \beta_j > 2(N_i + 1) \quad \forall i = 1, \dots, m, \quad (7)$$

is also one of the necessary conditions of radial solvability of problem (5).



# The "classical" Liouville equation

# The "classical" Liouville equation

- In the case of  $m = 1$ , and  $N = 0$  problem (2) reduces to:

$$\begin{cases} -\Delta v = \lambda e^v & \text{in } \mathbb{R}^2 \\ \int_{\mathbb{R}^2} e^v dx < +\infty. \end{cases} \quad (8)$$

# The "classical" Liouville equation

- In the case of  $m = 1$ , and  $N = 0$  problem (2) reduces to:

$$\begin{cases} -\Delta v = \lambda e^v & \text{in } \mathbb{R}^2 \\ \int_{\mathbb{R}^2} e^v dx < +\infty. \end{cases} \quad (8)$$

## Theorem (Chen-Li)

$$\text{If } \lambda > 0 \text{ then Eq.(8)} \Rightarrow \frac{\lambda}{2\pi} \int_{\mathbb{R}^2} e^v dx = 4, \quad (9)$$

# The "classical" Liouville equation

- In the case of  $m = 1$ , and  $N = 0$  problem (2) reduces to:

$$\begin{cases} -\Delta v = \lambda e^v & \text{in } \mathbb{R}^2 \\ \int_{\mathbb{R}^2} e^v dx < +\infty. \end{cases} \quad (8)$$

## Theorem (Chen-Li)

$$\text{If } \lambda > 0 \text{ then Eq.(8)} \Rightarrow \frac{\lambda}{2\pi} \int_{\mathbb{R}^2} e^v dx = 4, \quad (9)$$

*and all the solutions are fully classified.*

# The "singular" Liouville equation

# The "singular" Liouville equation

- In the case of  $m = 1$ , and  $N \neq 0$  problem (2) reduces to:

$$\begin{cases} -\Delta v = |x|^{2N} e^v & \text{in } \mathbb{R}^2 \\ \int_{\mathbb{R}^2} |x|^{2N} e^v dx < +\infty. \end{cases} \quad (10)$$

# The "singular" Liouville equation

- In the case of  $m = 1$ , and  $N \neq 0$  problem (2) reduces to:

$$\begin{cases} -\Delta v = |x|^{2N} e^v & \text{in } \mathbb{R}^2 \\ \int_{\mathbb{R}^2} |x|^{2N} e^v dx < +\infty. \end{cases} \quad (10)$$

## Theorem (Prajapat-Tarantello)

$$Eq.(10) \Rightarrow \frac{1}{2\pi} \int_{\mathbb{R}^2} |x|^{2N} e^v dx = 4(N + 1), \quad (11)$$

# The "singular" Liouville equation

- In the case of  $m = 1$ , and  $N \neq 0$  problem (2) reduces to:

$$\begin{cases} -\Delta v = |x|^{2N} e^v & \text{in } \mathbb{R}^2 \\ \int_{\mathbb{R}^2} |x|^{2N} e^v dx < +\infty. \end{cases} \quad (10)$$

## Theorem (Prajapat-Tarantello)

$$Eq.(10) \Rightarrow \frac{1}{2\pi} \int_{\mathbb{R}^2} |x|^{2N} e^v dx = 4(N + 1), \quad (11)$$

*and all the solutions are fully classified.*



The system (2) when  $a_{ij} > 0$ .

# The system (2) when $a_{ij} > 0$ .

Consider the conditions:

$$\begin{cases} \beta_i > 0 & i \in \{1, \dots, m\} \\ \left( \sum_{i=1}^m \sum_{j=1}^m \frac{1}{2} a_{ij} \beta_i \beta_j \right) - \sum_{i=1}^m 2(N_i + 1) \beta_i = 0 \\ \left( \sum_{i \in J} \sum_{j \in J} \frac{1}{2} a_{ij} \beta_i \beta_j \right) - \sum_{i \in J} 2(N_i + 1) \beta_i < 0 \quad \forall J, 1 \leq |J| < m \end{cases} \quad (12)$$

## The system (2) when $a_{ij} > 0$ .

Consider the conditions:

$$\begin{cases} \beta_i > 0 & i \in \{1, \dots, m\} \\ \left( \sum_{i=1}^m \sum_{j=1}^m \frac{1}{2} a_{ij} \beta_i \beta_j \right) - \sum_{i=1}^m 2(N_i + 1) \beta_i = 0 \\ \left( \sum_{i \in J} \sum_{j \in J} \frac{1}{2} a_{ij} \beta_i \beta_j \right) - \sum_{i \in J} 2(N_i + 1) \beta_i < 0 \quad \forall J, 1 \leq |J| < m \end{cases} \quad (12)$$

### Theorem (Chipot-Shafrir-Wolanski)

*If  $a_{ij} > 0$ ,  $\det A \neq 0$  and  $N_i = 0$  then (12) are necessary and sufficient condition for radial solvability of (5).*

## The system (2) when $a_{ij} > 0$ .

Consider the conditions:

$$\begin{cases} \beta_i > 0 & i \in \{1, \dots, m\} \\ \left( \sum_{i=1}^m \sum_{j=1}^m \frac{1}{2} a_{ij} \beta_i \beta_j \right) - \sum_{i=1}^m 2(N_i + 1) \beta_i = 0 \\ \left( \sum_{i \in J} \sum_{j \in J} \frac{1}{2} a_{ij} \beta_i \beta_j \right) - \sum_{i \in J} 2(N_i + 1) \beta_i < 0 \quad \forall J, 1 \leq |J| < m \end{cases} \quad (12)$$

### Theorem (Chipot-Shafrir-Wolanski)

*If  $a_{ij} > 0$ ,  $\det A \neq 0$  and  $N_i = 0$  then (12) are necessary and sufficient condition for radial solvability of (5).*

### Theorem (C.S.Lin-Zhang)

*In the settings of previous theorem a radial solution to (5) is unique (up to scaling (3)).*

## Theorem (P-Tarantello)

*If  $a_{ij} > 0$  and  $N_i > -1$  then (12) are necessary and sufficient condition for radial solvability of (5). Moreover, such a solution is unique (up to scaling (3)).*

## Theorem (P-Tarantello)

*If  $a_{ij} > 0$  and  $N_i > -1$  then (12) are necessary and sufficient condition for radial solvability of (5). Moreover, such a solution is unique (up to scaling (3)).*

Here  $A$  can be degenerate. The particular case when  $\det A \neq 0$  was treated independently by C.S.Lin and Zhang.

# A degenerate system arisen in the study of selfgravitating strings

# A degenerate system arisen in the study of selfgravitating strings

- For  $b > 0$  and  $N > -1$  consider the problem:

$$\begin{cases} -\Delta v = e^{bv} + |x|^{2N} e^v & \text{in } \mathbb{R}^2 \\ \frac{1}{2\pi} \int_{\mathbb{R}^2} (e^{bv} + |x|^{2N} e^v) dx = \alpha. \end{cases} \quad (13)$$



# A degenerate system arisen in the study of selfgravitating strings

- For  $b > 0$  and  $N > -1$  consider the problem:

$$\begin{cases} -\Delta v = e^{bv} + |x|^{2N} e^v & \text{in } \mathbb{R}^2 \\ \frac{1}{2\pi} \int_{\mathbb{R}^2} (e^{bv} + |x|^{2N} e^v) dx = \alpha. \end{cases} \quad (13)$$

- It can be easily verified that if  $b = 1/(N + 1)$  then  $\alpha = 4(N + 1)$ .

# A degenerate system arisen in the study of selfgravitating strings

- For  $b > 0$  and  $N > -1$  consider the problem:

$$\begin{cases} -\Delta v = e^{bv} + |x|^{2N} e^v & \text{in } \mathbb{R}^2 \\ \frac{1}{2\pi} \int_{\mathbb{R}^2} (e^{bv} + |x|^{2N} e^v) dx = \alpha. \end{cases} \quad (13)$$

- It can be easily verified that if  $b = 1/(N + 1)$  then  $\alpha = 4(N + 1)$ .
- Then by setting  $u_1 = bv - \ln b$ ,  $u_2 = v$  we reduce (13) to the degenerate system of the form (5):

# A degenerate system arisen in the study of selfgravitating strings

- For  $b > 0$  and  $N > -1$  consider the problem:

$$\begin{cases} -\Delta v = e^{bv} + |x|^{2N} e^v & \text{in } \mathbb{R}^2 \\ \frac{1}{2\pi} \int_{\mathbb{R}^2} (e^{bv} + |x|^{2N} e^v) dx = \alpha. \end{cases} \quad (13)$$

- It can be easily verified that if  $b = 1/(N+1)$  then  $\alpha = 4(N+1)$ .
- Then by setting  $u_1 = bv - \ln b$ ,  $u_2 = v$  we reduce (13) to the degenerate system of the form (5):

$$\begin{cases} -\Delta u_1 = b^2 e^{u_1} + b|x|^{2N} e^{u_2} & \text{in } \mathbb{R}^2 \\ -\Delta u_2 = b e^{u_1} + |x|^{2N} e^{u_2} & \text{in } \mathbb{R}^2 \\ \frac{1}{2\pi} \int_{\mathbb{R}^2} e^{u_1} dx = \beta_1 \\ \frac{1}{2\pi} \int_{\mathbb{R}^2} |x|^{2N} e^{u_2} dx = \beta_2 \\ b\beta_1 + \beta_2 = \alpha. \end{cases} \quad (14)$$

## Theorem (P-Tarantello)

Assume that  $v$  is a radial solution of (13).

$$\text{if } b > \frac{1}{N+1} \text{ then } \max \left\{ \frac{4}{b}, 4(N+1) - \frac{4}{b} \right\} < \alpha < 4(N+1),$$

$$\text{if } 0 < b < \frac{1}{N+1} \text{ then } \max \left\{ 4(N+1), \frac{4}{b} - 4(N+1) \right\} < \alpha < \frac{4}{b}.$$

Moreover, in the later cases there exist the unique radial solution to (13).

# A system (5) with positively defined matrix $A$

# A system (5) with positively defined matrix $A$

- What can we say about solvability of (5) if  $A$  is positively defined but can contain negative entries.

# A system (5) with positively defined matrix $A$

- What can we say about solvability of (5) if  $A$  is positively defined but can contain negative entries.
- We focus on the case  $m = 2$ .

# A system (5) with positively defined matrix $A$

- What can we say about solvability of (5) if  $A$  is positively defined but can contain negative entries.
- We focus on the case  $m = 2$ .
- Then (5) reads as:

$$\begin{cases} -\Delta\psi = a_{11}|x|^{2N_1}e^\psi + a_{12}|x|^{2N_2}e^\varphi & \text{in } \mathbb{R}^2 \\ -\Delta\varphi = a_{22}|x|^{2N_2}e^\varphi + a_{12}|x|^{2N_1}e^\psi & \text{in } \mathbb{R}^2 \\ \frac{1}{2\pi} \int_{\mathbb{R}^2} |x|^{2N_1}e^\psi dx = \beta \\ \frac{1}{2\pi} \int_{\mathbb{R}^2} |x|^{2N_2}e^\varphi dx = \alpha, \end{cases} \quad (15)$$



# A system (5) with positively defined matrix $A$

- What can we say about solvability of (5) if  $A$  is positively defined but can contain negative entries.
- We focus on the case  $m = 2$ .
- Then (5) reads as:

$$\begin{cases} -\Delta\psi = a_{11}|x|^{2N_1}e^\psi + a_{12}|x|^{2N_2}e^\varphi & \text{in } \mathbb{R}^2 \\ -\Delta\varphi = a_{22}|x|^{2N_2}e^\varphi + a_{12}|x|^{2N_1}e^\psi & \text{in } \mathbb{R}^2 \\ \frac{1}{2\pi} \int_{\mathbb{R}^2} |x|^{2N_1}e^\psi dx = \beta \\ \frac{1}{2\pi} \int_{\mathbb{R}^2} |x|^{2N_2}e^\varphi dx = \alpha, \end{cases} \quad (15)$$

and positive definiteness reads as:

# A system (5) with positively defined matrix $A$

- What can we say about solvability of (5) if  $A$  is positively defined but can contain negative entries.
- We focus on the case  $m = 2$ .
- Then (5) reads as:

$$\begin{cases} -\Delta\psi = a_{11}|x|^{2N_1}e^\psi + a_{12}|x|^{2N_2}e^\varphi & \text{in } \mathbb{R}^2 \\ -\Delta\varphi = a_{22}|x|^{2N_2}e^\varphi + a_{12}|x|^{2N_1}e^\psi & \text{in } \mathbb{R}^2 \\ \frac{1}{2\pi} \int_{\mathbb{R}^2} |x|^{2N_1}e^\psi dx = \beta \\ \frac{1}{2\pi} \int_{\mathbb{R}^2} |x|^{2N_2}e^\varphi dx = \alpha, \end{cases} \quad (15)$$

and positive definiteness reads as:

$$a_{11} > 0, \quad a_{22} > 0, \quad \text{and} \quad a_{12}^2 < a_{11}a_{22}. \quad (16)$$

- Defining in (15):

$$u_1(x) = \psi(x) - \ln(a_{11}) \quad \text{and} \quad u_2(x) = \varphi(x) - \ln(a_{22}) \quad (17)$$

- Defining in (15):

$$u_1(x) = \psi(x) - \ln(a_{11}) \quad \text{and} \quad u_2(x) = \varphi(x) - \ln(a_{22}) \quad (17)$$

we rewrite (15) as:

- Defining in (15):

$$u_1(x) = \psi(x) - \ln(a_{11}) \quad \text{and} \quad u_2(x) = \varphi(x) - \ln(a_{22}) \quad (17)$$

we rewrite (15) as:

$$\begin{cases} -\Delta u_1 = |x|^{2N_1} e^{u_1} - \tau_1 |x|^{2N_2} e^{u_2} & \text{in } \mathbb{R}^2 \\ -\Delta u_2 = |x|^{2N_2} e^{u_2} - \tau_2 |x|^{2N_1} e^{u_1} & \text{in } \mathbb{R}^2 \\ \frac{1}{2\pi} \int_{\mathbb{R}^2} |x|^{2N_1} e^{u_1} dx = \beta_1, \\ \frac{1}{2\pi} \int_{\mathbb{R}^2} |x|^{2N_2} e^{u_2} dx = \beta_2, \end{cases} \quad (18)$$

- Defining in (15):

$$u_1(x) = \psi(x) - \ln(a_{11}) \quad \text{and} \quad u_2(x) = \varphi(x) - \ln(a_{22}) \quad (17)$$

we rewrite (15) as:

$$\begin{cases} -\Delta u_1 = |x|^{2N_1} e^{u_1} - \tau_1 |x|^{2N_2} e^{u_2} & \text{in } \mathbb{R}^2 \\ -\Delta u_2 = |x|^{2N_2} e^{u_2} - \tau_2 |x|^{2N_1} e^{u_1} & \text{in } \mathbb{R}^2 \\ \frac{1}{2\pi} \int_{\mathbb{R}^2} |x|^{2N_1} e^{u_1} dx = \beta_1, \\ \frac{1}{2\pi} \int_{\mathbb{R}^2} |x|^{2N_2} e^{u_2} dx = \beta_2, \end{cases} \quad (18)$$

where

$$\tau_1 := -\frac{a_{12}}{a_{22}}, \quad \tau_2 := -\frac{a_{12}}{a_{11}} \quad \text{and} \quad \beta_1 = \frac{\beta}{a_{11}}, \quad \beta_2 = \frac{\alpha}{a_{22}}. \quad (19)$$

- Defining in (15):

$$u_1(x) = \psi(x) - \ln(a_{11}) \quad \text{and} \quad u_2(x) = \varphi(x) - \ln(a_{22}) \quad (17)$$

we rewrite (15) as:

$$\begin{cases} -\Delta u_1 = |x|^{2N_1} e^{u_1} - \tau_1 |x|^{2N_2} e^{u_2} & \text{in } \mathbb{R}^2 \\ -\Delta u_2 = |x|^{2N_2} e^{u_2} - \tau_2 |x|^{2N_1} e^{u_1} & \text{in } \mathbb{R}^2 \\ \frac{1}{2\pi} \int_{\mathbb{R}^2} |x|^{2N_1} e^{u_1} dx = \beta_1, \\ \frac{1}{2\pi} \int_{\mathbb{R}^2} |x|^{2N_2} e^{u_2} dx = \beta_2, \end{cases} \quad (18)$$

where

$$\tau_1 := -\frac{a_{12}}{a_{22}}, \quad \tau_2 := -\frac{a_{12}}{a_{11}} \quad \text{and} \quad \beta_1 = \frac{\beta}{a_{11}}, \quad \beta_2 = \frac{\alpha}{a_{22}}. \quad (19)$$

Moreover, in the case  $a_{12} \neq 0$  (16) reads as:

- Defining in (15):

$$u_1(x) = \psi(x) - \ln(a_{11}) \quad \text{and} \quad u_2(x) = \varphi(x) - \ln(a_{22}) \quad (17)$$

we rewrite (15) as:

$$\begin{cases} -\Delta u_1 = |x|^{2N_1} e^{u_1} - \tau_1 |x|^{2N_2} e^{u_2} & \text{in } \mathbb{R}^2 \\ -\Delta u_2 = |x|^{2N_2} e^{u_2} - \tau_2 |x|^{2N_1} e^{u_1} & \text{in } \mathbb{R}^2 \\ \frac{1}{2\pi} \int_{\mathbb{R}^2} |x|^{2N_1} e^{u_1} dx = \beta_1, \\ \frac{1}{2\pi} \int_{\mathbb{R}^2} |x|^{2N_2} e^{u_2} dx = \beta_2, \end{cases} \quad (18)$$

where

$$\tau_1 := -\frac{a_{12}}{a_{22}}, \quad \tau_2 := -\frac{a_{12}}{a_{11}} \quad \text{and} \quad \beta_1 = \frac{\beta}{a_{11}}, \quad \beta_2 = \frac{\alpha}{a_{22}}. \quad (19)$$

Moreover, in the case  $a_{12} \neq 0$  (16) reads as:

$$0 < \tau_1 \tau_2 < 1. \quad (20)$$



- In the case  $a_{12} < 0$  we have  $\tau_1 > 0$ ,  $\tau_2 > 0$  and  $\tau_1\tau_2 < 1$ .

- In the case  $a_{12} < 0$  we have  $\tau_1 > 0$ ,  $\tau_2 > 0$  and  $\tau_1\tau_2 < 1$ .  
Then, Pohozaev identity (6) reads as:

- In the case  $a_{12} < 0$  we have  $\tau_1 > 0$ ,  $\tau_2 > 0$  and  $\tau_1\tau_2 < 1$ .

Then, Pohozaev identity (6) reads as:

$$\tau_2\beta_1^2 - 4\tau_2(N_1 + 1)\beta_1 + \tau_1\beta_2^2 - 4\tau_1(N_2 + 1)\beta_2 - 2\tau_1\tau_2\beta_1\beta_2 = 0. \quad (21)$$

- In the case  $a_{12} < 0$  we have  $\tau_1 > 0$ ,  $\tau_2 > 0$  and  $\tau_1\tau_2 < 1$ .

Then, Pohozaev identity (6) reads as:

$$\tau_2\beta_1^2 - 4\tau_2(N_1 + 1)\beta_1 + \tau_1\beta_2^2 - 4\tau_1(N_2 + 1)\beta_2 - 2\tau_1\tau_2\beta_1\beta_2 = 0. \quad (21)$$

Moreover, (6) and (7) together read as:

- In the case  $a_{12} < 0$  we have  $\tau_1 > 0$ ,  $\tau_2 > 0$  and  $\tau_1\tau_2 < 1$ .

Then, Pohozaev identity (6) reads as:

$$\tau_2\beta_1^2 - 4\tau_2(N_1 + 1)\beta_1 + \tau_1\beta_2^2 - 4\tau_1(N_2 + 1)\beta_2 - 2\tau_1\tau_2\beta_1\beta_2 = 0. \quad (21)$$

Moreover, (6) and (7) together read as:

$$\left\{ \begin{array}{l} \beta_2 < \frac{2}{1-\tau_1\tau_2} \left( (N_2 + 1) + \tau_2(N_1 + 1) \right. \\ \left. + \sqrt{(N_2 + 1)^2 + 2\tau_2(N_2 + 1)(N_1 + 1) + \frac{\tau_2}{\tau_1}(N_1 + 1)^2} \right), \\ \\ \beta_2 > \frac{2}{1-\tau_1\tau_2} \left( (N_2 + 1) + \tau_2(N_1 + 1) \right. \\ \left. + (\sqrt{\tau_1\tau_2}) \sqrt{(N_2 + 1)^2 + 2\tau_2(N_2 + 1)(N_1 + 1) + \frac{\tau_2}{\tau_1}(N_1 + 1)^2} \right), \\ \\ \beta_1 = \frac{(2(N_1 + 1) + \tau_1\beta_2)}{\tau_1} \\ \left. + \sqrt{(2(N_1 + 1) + \tau_1\beta_2)^2 - \frac{\tau_1}{\tau_2}\beta_2(\beta_2 - 4(N_2 + 1))} \right\} \end{array} \right.$$

Moreover, similarly as it was done in the case  $a_{12} > 0$ , in the case  $a_{12} < 0$  we also can find that the following condition

Moreover, similarly as it was done in the case  $a_{12} > 0$ , in the case  $a_{12} < 0$  we also can find that the following condition

$$\beta_1 > 4(N_1 + 1) \quad \text{and} \quad \beta_2 > 4(N_2 + 1), \quad (23)$$

Moreover, similarly as it was done in the case  $a_{12} > 0$ , in the case  $a_{12} < 0$  we also can find that the following condition

$$\beta_1 > 4(N_1 + 1) \quad \text{and} \quad \beta_2 > 4(N_2 + 1), \quad (23)$$

is also one of the necessary conditions of radial solvability of problem (18).



The following cases of system (18) are special:

The following cases of system (18) are special:

- $\tau_1 = \tau_2 = \frac{1}{2}$ ,

The following cases of system (18) are special:

- $\tau_1 = \tau_2 = \frac{1}{2}$ ,
- Either  $\tau_1 = \frac{1}{2}$  and  $\tau_2 = 1$ , or  $\tau_1 = 1$  and  $\tau_2 = \frac{1}{2}$ .

The following cases of system (18) are special:

- $\tau_1 = \tau_2 = \frac{1}{2}$ ,
- Either  $\tau_1 = \frac{1}{2}$  and  $\tau_2 = 1$ , or  $\tau_1 = 1$  and  $\tau_2 = \frac{1}{2}$ .
- Either  $\tau_1 = \frac{1}{2}$  and  $\tau_2 = \frac{3}{2}$ , or  $\tau_1 = \frac{3}{2}$  and  $\tau_2 = \frac{1}{2}$ .

The following cases of system (18) are special:

- $\tau_1 = \tau_2 = \frac{1}{2}$ ,
- Either  $\tau_1 = \frac{1}{2}$  and  $\tau_2 = 1$ , or  $\tau_1 = 1$  and  $\tau_2 = \frac{1}{2}$ .
- Either  $\tau_1 = \frac{1}{2}$  and  $\tau_2 = \frac{3}{2}$ , or  $\tau_1 = \frac{3}{2}$  and  $\tau_2 = \frac{1}{2}$ .

In these cases it was proved (C.S.Lin,Wei and their coauthors) that the set of  $(\beta_1, \beta_2)$  for which we have a radial solvability of (18) reduces to a single point.

The following cases of system (18) are special:

- $\tau_1 = \tau_2 = \frac{1}{2}$ ,
- Either  $\tau_1 = \frac{1}{2}$  and  $\tau_2 = 1$ , or  $\tau_1 = 1$  and  $\tau_2 = \frac{1}{2}$ .
- Either  $\tau_1 = \frac{1}{2}$  and  $\tau_2 = \frac{3}{2}$ , or  $\tau_1 = \frac{3}{2}$  and  $\tau_2 = \frac{1}{2}$ .

In these cases it was proved (C.S.Lin,Wei and their coauthors) that the set of  $(\beta_1, \beta_2)$  for which we have a radial solvability of (18) reduces to a single point.

For example: if  $\tau_1 = \tau_2 = \frac{1}{2}$  then necessarily  
$$\beta_1 = \beta_2 = 4(N_1 + 1) + 4(N_2 + 1)$$

The following cases of system (18) are special:

- $\tau_1 = \tau_2 = \frac{1}{2}$ ,
- Either  $\tau_1 = \frac{1}{2}$  and  $\tau_2 = 1$ , or  $\tau_1 = 1$  and  $\tau_2 = \frac{1}{2}$ .
- Either  $\tau_1 = \frac{1}{2}$  and  $\tau_2 = \frac{3}{2}$ , or  $\tau_1 = \frac{3}{2}$  and  $\tau_2 = \frac{1}{2}$ .

In these cases it was proved (C.S.Lin, Wei and their coauthors) that the set of  $(\beta_1, \beta_2)$  for which we have a radial solvability of (18) reduces to a single point.

For example: if  $\tau_1 = \tau_2 = \frac{1}{2}$  then necessarily

$$\beta_1 = \beta_2 = 4(N_1 + 1) + 4(N_2 + 1)$$

and if  $\tau_1 = \frac{1}{2}$ ,  $\tau_2 = 1$  then necessarily  $\beta_1 = 8(N_1 + 1) + 4(N_2 + 1)$

and  $\beta_2 = 8(N_1 + 1) + 8(N_2 + 1)$ .

The system (18) in the case  $\tau_1 = \tau_2$



# The system (18) in the case $\tau_1 = \tau_2$

For  $\tau \in (0, 1)$  consider the system:

$$\begin{cases} -\Delta u_1 = |x|^{2N_1} e^{u_1} - \tau |x|^{2N_2} e^{u_2} & \text{in } \mathbb{R}^2 \\ -\Delta u_2 = |x|^{2N_2} e^{u_2} - \tau |x|^{2N_1} e^{u_1} & \text{in } \mathbb{R}^2 \\ \frac{1}{2\pi} \int_{\mathbb{R}^2} |x|^{2N_1} e^{u_1} dx = \beta_1, \\ \frac{1}{2\pi} \int_{\mathbb{R}^2} |x|^{2N_2} e^{u_2} dx = \beta_2, \end{cases} \quad (24)$$

# The system (18) in the case $\tau_1 = \tau_2$

For  $\tau \in (0, 1)$  consider the system:

$$\begin{cases} -\Delta u_1 = |x|^{2N_1} e^{u_1} - \tau |x|^{2N_2} e^{u_2} & \text{in } \mathbb{R}^2 \\ -\Delta u_2 = |x|^{2N_2} e^{u_2} - \tau |x|^{2N_1} e^{u_1} & \text{in } \mathbb{R}^2 \\ \frac{1}{2\pi} \int_{\mathbb{R}^2} |x|^{2N_1} e^{u_1} dx = \beta_1, \\ \frac{1}{2\pi} \int_{\mathbb{R}^2} |x|^{2N_2} e^{u_2} dx = \beta_2, \end{cases} \quad (24)$$

- If  $\tau = 1/2$  then it is well known Toda system, the radial solution exists if and only if  $\beta_1 = \beta_2 = 4(N_1 + N_2 + 2)$  and they are completely classified (C.S.Lin-Wei-Ye).

## Theorem (P-Tarantello)

For every  $\tau \in (0, 1) \setminus \{1/2\}$  the necessary and sufficient conditions on  $(\beta_1, \beta_2)$  for the existence of a radial solution to (24) are the following:

$$\begin{cases} \frac{1}{2}\beta_1^2 - 2(N_1 + 1)\beta_1 + \frac{1}{2}\beta_2^2 - 2(N_2 + 1)\beta_2 - \tau\beta_1\beta_2 = 0, \\ \underline{\beta}_1(\tau) < \beta_1 < \bar{\beta}_1(\tau) \\ \underline{\beta}_2(\tau) < \beta_2 < \bar{\beta}_2(\tau). \end{cases} \quad (25)$$

where  $\underline{\beta}_1(\tau), \bar{\beta}_1(\tau), \underline{\beta}_2(\tau), \bar{\beta}_2(\tau)$  are given by some formulas.

# Definitions of $\underline{\beta}_1(\tau), \overline{\beta}_1(\tau), \underline{\beta}_2(\tau), \overline{\beta}_2(\tau)$

# Definitions of $\underline{\beta}_1(\tau), \overline{\beta}_1(\tau), \underline{\beta}_2(\tau), \overline{\beta}_2(\tau)$

- There exists unique  $\delta_1 \in (0, 1/2)$  such that

$$4(N_2 + 1) = 2\delta_1 \left( 4(N_1 + 1) + 8\delta_1(N_2 + 1) \right) \quad (26)$$

# Definitions of $\underline{\beta}_1(\tau), \overline{\beta}_1(\tau), \underline{\beta}_2(\tau), \overline{\beta}_2(\tau)$

- There exists unique  $\delta_1 \in (0, 1/2)$  such that

$$4(N_2 + 1) = 2\delta_1 \left( 4(N_1 + 1) + 8\delta_1(N_2 + 1) \right) \quad (26)$$

- there exists unique  $\delta_2 \in (1/2, 1/\sqrt{2})$  such that

$$8(N_2 + 1) + \frac{2}{\delta_2}(N_1 + 1) = 2\delta_2 \left( 8\delta_2(N_2 + 1) + 4(N_1 + 1) \right) \quad (27)$$

# Definitions of $\underline{\beta}_1(\tau), \overline{\beta}_1(\tau), \underline{\beta}_2(\tau), \overline{\beta}_2(\tau)$

- There exists unique  $\delta_1 \in (0, 1/2)$  such that

$$4(N_2 + 1) = 2\delta_1 \left( 4(N_1 + 1) + 8\delta_1(N_2 + 1) \right) \quad (26)$$

- there exists unique  $\delta_2 \in (1/2, 1/\sqrt{2})$  such that

$$8(N_2 + 1) + \frac{2}{\delta_2}(N_1 + 1) = 2\delta_2 \left( 8\delta_2(N_2 + 1) + 4(N_1 + 1) \right) \quad (27)$$

- there exists unique  $\sigma_1 \in (0, 1/2)$  such that

$$4(N_1 + 1) = 2\sigma_1 \left( 4(N_2 + 1) + 8\sigma_1(N_1 + 1) \right) \quad (28)$$

# Definitions of $\underline{\beta}_1(\tau), \overline{\beta}_1(\tau), \underline{\beta}_2(\tau), \overline{\beta}_2(\tau)$

- There exists unique  $\delta_1 \in (0, 1/2)$  such that

$$4(N_2 + 1) = 2\delta_1 \left( 4(N_1 + 1) + 8\delta_1(N_2 + 1) \right) \quad (26)$$

- there exists unique  $\delta_2 \in (1/2, 1/\sqrt{2})$  such that

$$8(N_2 + 1) + \frac{2}{\delta_2}(N_1 + 1) = 2\delta_2 \left( 8\delta_2(N_2 + 1) + 4(N_1 + 1) \right) \quad (27)$$

- there exists unique  $\sigma_1 \in (0, 1/2)$  such that

$$4(N_1 + 1) = 2\sigma_1 \left( 4(N_2 + 1) + 8\sigma_1(N_1 + 1) \right) \quad (28)$$

- there exists unique  $\sigma_2 \in (1/2, 1/\sqrt{2})$  such that

$$8(N_1 + 1) + \frac{2}{\sigma_2}(N_2 + 1) = 2\sigma_2 \left( 8\sigma_2(N_1 + 1) + 4(N_2 + 1) \right)$$



$$\underline{\beta}_{-1}(\tau) = \begin{cases} 4(N_1 + 1) & \forall \tau \in (0, \sigma_1) \\ 2\tau(4(N_2 + 1) + 8\tau(N_1 + 1)) & \forall \tau \in [\sigma_1, 1/2) \\ (4(N_1 + 1) + 8\tau(N_2 + 1)) & \forall \tau \in [1/2, \delta_2) \\ \frac{2((N_1+1)+\tau(N_2+1)+\tau\sqrt{(N_1+1)^2+(N_2+1)^2+2\tau(N_1+1)(N_2+1)})}{1-\tau^2} & \forall \tau \geq \delta_2 \end{cases} \quad (30)$$

$$\underline{\beta}_1(\tau) = \begin{cases} 4(N_1 + 1) & \forall \tau \in (0, \sigma_1) \\ 2\tau(4(N_2 + 1) + 8\tau(N_1 + 1)) & \forall \tau \in [\sigma_1, 1/2) \\ (4(N_1 + 1) + 8\tau(N_2 + 1)) & \forall \tau \in [1/2, \delta_2) \\ \frac{2((N_1+1)+\tau(N_2+1)+\tau\sqrt{(N_1+1)^2+(N_2+1)^2+2\tau(N_1+1)(N_2+1)})}{1-\tau^2} & \forall \tau \geq \delta_2 \end{cases} \quad (30)$$

$$\overline{\beta}_1(\tau) = \begin{cases} (4(N_1 + 1) + 8\tau(N_2 + 1)) & \forall \tau \in (0, 1/2) \\ 2\tau(4(N_2 + 1) + 8\tau(N_1 + 1)) & \forall \tau \in [1/2, \sigma_2) \\ \frac{2((N_1+1)+\tau(N_2+1)+\sqrt{(N_1+1)^2+(N_2+1)^2+2\tau(N_1+1)(N_2+1)})}{1-\tau^2} & \forall \tau \geq \sigma_2. \end{cases} \quad (31)$$

$$\underline{\beta}_2(\tau) = \begin{cases} 4(N_2 + 1) & \forall \tau \in (0, \delta_1) \\ 2\tau(4(N_1 + 1) + 8\tau(N_2 + 1)) & \forall \tau \in [\delta_1, 1/2) \\ (4(N_2 + 1) + 8\tau(N_1 + 1)) & \forall \tau \in [1/2, \sigma_2) \\ \frac{2((N_2+1)+\tau(N_1+1)+\tau\sqrt{(N_2+1)^2+(N_1+1)^2+2\tau(N_2+1)(N_1+1)})}{1-\tau^2} & \forall \tau \geq \sigma_2 \end{cases}$$

$$\underline{\beta}_2(\tau) = \begin{cases} 4(N_2 + 1) & \forall \tau \in (0, \delta_1) \\ 2\tau(4(N_1 + 1) + 8\tau(N_2 + 1)) & \forall \tau \in [\delta_1, 1/2) \\ (4(N_2 + 1) + 8\tau(N_1 + 1)) & \forall \tau \in [1/2, \sigma_2) \\ \frac{2((N_2+1)+\tau(N_1+1)+\tau\sqrt{(N_2+1)^2+(N_1+1)^2+2\tau(N_2+1)(N_1+1)})}{1-\tau^2} & \forall \tau \geq \sigma_2 \end{cases} \quad (32)$$

$$\overline{\beta}_2(\tau) = \begin{cases} (4(N_2 + 1) + 8\tau(N_1 + 1)) & \forall \tau \in (0, 1/2) \\ 2\tau(4(N_1 + 1) + 8\tau(N_2 + 1)) & \forall \tau \in [1/2, \delta_2) \\ \frac{2((N_2+1)+\tau(N_1+1)+\sqrt{(N_2+1)^2+(N_1+1)^2+2\tau(N_2+1)(N_1+1)})}{1-\tau^2} & \forall \tau \geq \delta_2 \end{cases} \quad (33)$$

## Lemma

For every  $\tau \in (0, 1)$   $\theta \in \mathbb{R}$  consider  $(v_1^{(\theta)}, v_2^{(\theta)})$  be radial solution of

## Lemma

For every  $\tau \in (0, 1)$   $\theta \in \mathbb{R}$  consider  $(v_1^{(\theta)}, v_2^{(\theta)})$  be radial solution of

$$\begin{cases} -\Delta v_1^{(\theta)} = |x|^{2N_1} e^{v_1^{(\theta)}} - \tau |x|^{2N_2} e^{v_2^{(\theta)}} & \text{in } \mathbb{R}^2 \\ -\Delta v_2^{(\theta)} = |x|^{2N_2} e^{v_2^{(\theta)}} - \tau |x|^{2N_1} e^{v_1^{(\theta)}} & \text{in } \mathbb{R}^2 \\ \psi^{(\theta)}(0) = \theta \\ \varphi^{(\theta)}(0) = 0, \end{cases}$$

## Lemma

For every  $\tau \in (0, 1)$   $\theta \in \mathbb{R}$  consider  $(v_1^{(\theta)}, v_2^{(\theta)})$  be radial solution of

$$\begin{cases} -\Delta v_1^{(\theta)} = |x|^{2N_1} e^{v_1^{(\theta)}} - \tau |x|^{2N_2} e^{v_2^{(\theta)}} & \text{in } \mathbb{R}^2 \\ -\Delta v_2^{(\theta)} = |x|^{2N_2} e^{v_2^{(\theta)}} - \tau |x|^{2N_1} e^{v_1^{(\theta)}} & \text{in } \mathbb{R}^2 \\ \psi^{(\theta)}(0) = \theta \\ \varphi^{(\theta)}(0) = 0, \end{cases} \quad (34)$$

$$\tilde{\beta}_1(\theta) := \frac{1}{2\pi} \int_{\mathbb{R}^2} |x|^{2N_1} e^{v_1^{(\theta)}} dx, \quad \tilde{\beta}_2(\theta) := \frac{1}{2\pi} \int_{\mathbb{R}^2} |x|^{2N_2} e^{v_2^{(\theta)}} dx.$$

## Lemma

For every  $\tau \in (0, 1)$   $\theta \in \mathbb{R}$  consider  $(v_1^{(\theta)}, v_2^{(\theta)})$  be radial solution of

$$\begin{cases} -\Delta v_1^{(\theta)} = |x|^{2N_1} e^{v_1^{(\theta)}} - \tau |x|^{2N_2} e^{v_2^{(\theta)}} & \text{in } \mathbb{R}^2 \\ -\Delta v_2^{(\theta)} = |x|^{2N_2} e^{v_2^{(\theta)}} - \tau |x|^{2N_1} e^{v_1^{(\theta)}} & \text{in } \mathbb{R}^2 \\ \psi^{(\theta)}(0) = \theta \\ \varphi^{(\theta)}(0) = 0, \end{cases} \quad (34)$$

$$\tilde{\beta}_1(\theta) := \frac{1}{2\pi} \int_{\mathbb{R}^2} |x|^{2N_1} e^{v_1^{(\theta)}} dx, \quad \tilde{\beta}_2(\theta) := \frac{1}{2\pi} \int_{\mathbb{R}^2} |x|^{2N_2} e^{v_2^{(\theta)}} dx.$$

Furthermore, let  $T_\tau^{(1)}$  be the open interval with endpoints  $\lim_{\theta \rightarrow \pm\infty} \tilde{\beta}_1(\theta)$  and  $T_\tau^{(2)}$  be the open interval with endpoints  $\lim_{\theta \rightarrow \pm\infty} \tilde{\beta}_2(\theta)$ . Then



## Lemma

For every  $\tau \in (0, 1)$   $\theta \in \mathbb{R}$  consider  $(v_1^{(\theta)}, v_2^{(\theta)})$  be radial solution of

$$\begin{cases} -\Delta v_1^{(\theta)} = |x|^{2N_1} e^{v_1^{(\theta)}} - \tau |x|^{2N_2} e^{v_2^{(\theta)}} & \text{in } \mathbb{R}^2 \\ -\Delta v_2^{(\theta)} = |x|^{2N_2} e^{v_2^{(\theta)}} - \tau |x|^{2N_1} e^{v_1^{(\theta)}} & \text{in } \mathbb{R}^2 \\ \psi^{(\theta)}(0) = \theta \\ \varphi^{(\theta)}(0) = 0, \end{cases} \quad (34)$$

$$\tilde{\beta}_1(\theta) := \frac{1}{2\pi} \int_{\mathbb{R}^2} |x|^{2N_1} e^{v_1^{(\theta)}} dx, \quad \tilde{\beta}_2(\theta) := \frac{1}{2\pi} \int_{\mathbb{R}^2} |x|^{2N_2} e^{v_2^{(\theta)}} dx.$$

Furthermore, let  $T_\tau^{(1)}$  be the open interval with endpoints  $\lim_{\theta \rightarrow \pm\infty} \tilde{\beta}_1(\theta)$  and  $T_\tau^{(2)}$  be the open interval with endpoints  $\lim_{\theta \rightarrow \pm\infty} \tilde{\beta}_2(\theta)$ . Then

$$T_\tau^{(1)} = \left( \underline{\beta}_1(\tau), \bar{\beta}_1(\tau) \right) \quad \text{and} \quad T_\tau^{(2)} = \left( \underline{\beta}_2(\tau), \bar{\beta}_2(\tau) \right).$$

## Lemma

*Let  $(\tau_1, \tau_2) \neq (1/2, 1/2)$  be such that  $(\tau_1 - 1/2)(\tau_2 - 1/2) \geq 0$  be a radial solution of*

## Lemma

Let  $(\tau_1, \tau_2) \neq (1/2, 1/2)$  be such that  $(\tau_1 - 1/2)(\tau_2 - 1/2) \geq 0$  be a radial solution of

$$\begin{cases} -\Delta u_1 = |x|^{2N_1} e^{u_1} - \tau_1 |x|^{2N_2} e^{u_2} & \text{in } \mathbb{R}^2 \\ -\Delta u_2 = |x|^{2N_2} e^{u_2} - \tau_2 |x|^{2N_1} e^{u_1} & \text{in } \mathbb{R}^2 \\ \frac{1}{2\pi} \int_{\mathbb{R}^2} |x|^{2N_1} e^{u_1} dx = \beta_1, \\ \frac{1}{2\pi} \int_{\mathbb{R}^2} |x|^{2N_2} e^{u_2} dx = \beta_2, \end{cases} \quad (35)$$

## Lemma

Let  $(\tau_1, \tau_2) \neq (1/2, 1/2)$  be such that  $(\tau_1 - 1/2)(\tau_2 - 1/2) \geq 0$  be a radial solution of

$$\begin{cases} -\Delta u_1 = |x|^{2N_1} e^{u_1} - \tau_1 |x|^{2N_2} e^{u_2} & \text{in } \mathbb{R}^2 \\ -\Delta u_2 = |x|^{2N_2} e^{u_2} - \tau_2 |x|^{2N_1} e^{u_1} & \text{in } \mathbb{R}^2 \\ \frac{1}{2\pi} \int_{\mathbb{R}^2} |x|^{2N_1} e^{u_1} dx = \beta_1, \\ \frac{1}{2\pi} \int_{\mathbb{R}^2} |x|^{2N_2} e^{u_2} dx = \beta_2, \end{cases} \quad (35)$$

Then

$$\beta_1 \neq 4(N_1 + 1) + 8\tau_1(N_2 + 1) \quad \text{and} \quad \beta_2 \neq 4(N_2 + 1) + 8\tau_2(N_1 + 1).$$

# Thank You!