

Large-time behavior in hypocoercive BGK-models

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Decay for nonsymmetric ODEs: find Lyapunov functionals!

$$\dot{x} = -\mathbf{C}x, \quad t \geq 0, \quad x(t) \in \mathbb{R}^n \quad (1)$$

Definition: \mathbf{C} is *coercive* if $x^T \mathbf{C}x \geq \kappa \|x\|^2 \quad \forall x$ (for some $\kappa > 0$).

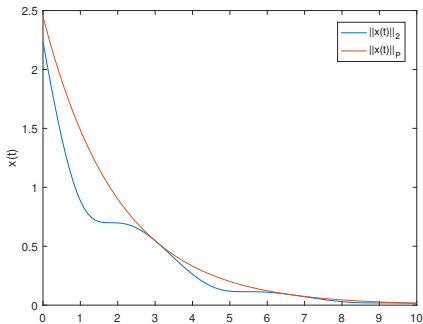
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ex: $\mathbf{C} = \begin{pmatrix} 1 & -1 \\ 1 & 0 \end{pmatrix}$, $\lambda_{\mathbf{C}} = \frac{1}{2} \pm i\frac{\sqrt{3}}{2} \Rightarrow$ decay rate = $\frac{1}{2}$ for (1).

- \mathbf{C} not coercive \Rightarrow no decay of $\|x(t)\|_2$ by trivial energy method!
- But decay of **modified norm** $\|x(t)\|_{\mathbf{P}} := \sqrt{x^T \mathbf{P} x}$; $\mathbf{P} := \begin{bmatrix} 2 & -1 \\ -1 & 2 \end{bmatrix}$



Decay for nonsymmetric ODEs: find Lyapunov functionals!

$$\dot{x} = -\mathbf{C}x, \quad t \geq 0, \quad x(t) \in \mathbb{R}^n$$

Definition: \mathbf{C} is *hypocoercive* if $\exists \mu > 0$ such that:

$$\Re(\lambda_j) \geq \mu, \quad j = 1, \dots, n.$$

If all eigenvalues of \mathbf{C} are non-defective:

$$\exists c \geq 1 : \quad \|x(t)\|_2 \leq c \|x(0)\|_2 e^{-\mu t}, \quad t \geq 0.$$

- always: $\kappa \leq \mu$

Choice of \mathbf{P} / Lyapunov's direct method

Lemma 1

Let $\mathbf{C} \in \mathbb{R}^{n \times n}$ be positive stable, i.e. $\mu := \min\{\Re \lambda_{\mathbf{C}}\} > 0$.

- ① If all $\lambda_{\mathbf{C}}^{\min} \in \{\lambda \in \sigma(\mathbf{C}) \mid \Re \lambda = \mu\}$ are *non-defective* (i.e. geometric = algebraic multiplicity)

$$\Rightarrow \exists \mathbf{P} \in \mathbb{R}^{n \times n}, \mathbf{P} > 0 : \quad \mathbf{P}\mathbf{C} + \mathbf{C}^{\top}\mathbf{P} \geq 2\mu\mathbf{P}.$$

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- ② If (at least) one $\lambda_{\mathbf{C}}^{\min}$ is *defective* \Rightarrow

$$\forall \varepsilon > 0 \quad \exists \mathbf{P} = \mathbf{P}(\varepsilon) > 0 : \mathbf{P}\mathbf{C} + \mathbf{C}^{\top}\mathbf{P} \geq 2(\mu - \varepsilon)\mathbf{P}.$$

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Proof: \mathbf{P} can be constructed explicitly; e.g. for \mathbf{C} non-defective / diagonalizable:

$$\mathbf{P} := \sum_{j=1}^n z_j \otimes \bar{z}_j^{\top}; \quad z_j \dots \text{eigenvectors of } \mathbf{C}^{\top}$$

- \mathbf{P} not unique; but the decay rates μ (or $\mu - \varepsilon$) are independent of \mathbf{P} . □
- For complex \mathbf{C} : $\mathbf{P} > 0$ Hermitian with $\mathbf{P}\mathbf{C} + \mathbf{C}^*\mathbf{P} \geq 2\mu\mathbf{P}$.

Decay of \mathbf{P} -norm

- Sharp decay estimate for $\dot{x} = -\mathbf{C}x$ (non-defective case):

Let $\|x\|_{\mathbf{P}}^2 := x^T \mathbf{P} x$.

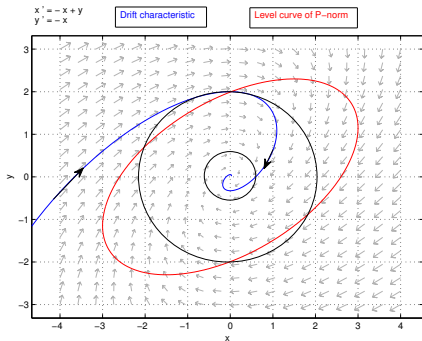
$$\frac{d}{dt} \|x\|_{\mathbf{P}}^2 = -x^T \underbrace{(\mathbf{P}\mathbf{C} + \mathbf{C}^T \mathbf{P})}_{\geq 2\mu \mathbf{P}} x \leq -2\mu \|x\|_{\mathbf{P}}^2$$

$$\Rightarrow \|x(t)\|_{\mathbf{P}} \leq \|x(0)\|_{\mathbf{P}} e^{-\mu t}, \quad t \geq 0.$$

Decay of \mathbf{P} -norm (cont'd)

ex: $\dot{x} = -\mathbf{C}x$ with $\mathbf{C} = \begin{pmatrix} 1 & -1 \\ 1 & 0 \end{pmatrix}$

- At x_2 -axis: trajectory $x(t)$ tangent to level curve of $|x|$:



- level curve of “distorted” vector norm $\sqrt{x^T \mathbf{P} x}$; $\mathbf{P} = \begin{pmatrix} 2 & -1 \\ -1 & 2 \end{pmatrix}$
→ uniform decay with sharp rate $\frac{1}{2}$

Wavy decay of energy/entropy in degenerate kinetic eq.

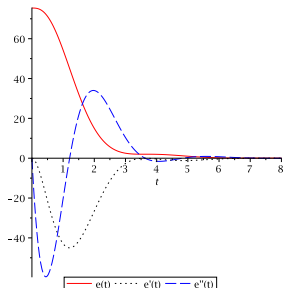
kinetic Fokker-Planck equation from plasma physics, for density

$f(x, v, t)$, $x, v \in \mathbb{R}^d$:

$$f_t + \underbrace{v \cdot \nabla_x f}_{\text{free transport}} - \underbrace{\nabla_x V \cdot \nabla_v f}_{\text{influence of potential } V(x)} = \underbrace{\Delta_v f}_{\text{diffusion}} + \underbrace{\text{div}_v(vf)}_{\text{friction}}$$

in 1D, $V(x) = x^2/2$:

$$e(t) = \|f(t) - f_\infty\|_{L^2(f_\infty^{-1})}^2$$



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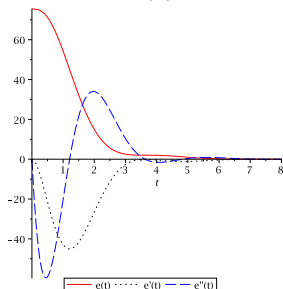
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- **local-in-x equilibration** of $f(x, v, t)$ towards $\rho(x) \frac{1}{\sqrt{2\pi}} e^{-|v|^2/2}$
- (Hamiltonian) **rotation in phase space** “mixes” positions
⇒ global equilibration

Space-inhomogen. Bhatnagar-Gross-Krook model (1954)

- application: gas dynamics, computational fluid dynamics
- simplified kinetic model: **relaxation operator** (instead of $Q(f, f)$):

$$f_t + v \cdot \nabla_x f = M_f(x, v, t) - f(x, v, t), \quad t \geq 0, \quad x \in \mathbb{T}^d, \quad v \in \mathbb{R}^d$$

- **relaxation towards local Maxwellian** M_f with same hydrodynamic moments as f :

$$M_f(x, v) = \frac{\rho(x)}{(2\pi T(x))^{\frac{d}{2}}} e^{-\frac{|v-u(x)|^2}{2T(x)}},$$

$$\text{density } \rho(x) := \int_{\mathbb{R}^d} f(x, v) dv,$$

$$\text{mean velocity } u(x) := \frac{1}{\rho(x)} \int_{\mathbb{R}^d} v f(x, v) dv,$$

$$\text{temperature } T(x) := \frac{1}{d\rho(x)} \int_{\mathbb{R}^d} |v - u(x)|^2 f(x, v) dv.$$

nonlinear BGK model

Results:

- Perthame 1989: global existence for $(x, v) \in \mathbb{R}^{2d}$ (compactness of M_f in L^1)
- Perthame-Pulvirenti 1993: uniqueness for $x \in \mathbb{T}^d$, $v \in \mathbb{R}^d$ (contraction, $\|f\|_\infty$ -bounds)

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Goals:

- exponential convergence to equilibrium
- partially with sharp rates

Outline:

- ① hypocoercivity
- ② linear BGK models (discrete & continuous velocity) with 1 conserved quantity
- ③ linearized & nonlinear BGK model with 2 conserved quantities

1D linear BGK model for $f(x, v, t)$: only mass conservation

$$f_t + vf_x = Qf := M_T(v) \int_{\mathbb{R}} f(x, v, t) dv - f(x, v, t), \quad t \geq 0,$$

$$(x, v) \in \mathbb{T}^1 \times \mathbb{R}$$

$$f^\infty(x, v) = M_T(v) := (2\pi T)^{-1/2} e^{-v^2/2T}; \quad \text{some fixed temperature } T,$$

$$Lf := -vf_x + Qf$$

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Q ... selfadjoint on $H := L^2(\mathbb{T}^1 \times \mathbb{R}; M_T^{-1}(v)dv)$

- collision operator Q drives f to a *local* Maxwellian $\rho(x) M_T(v)$.
- transport operator $-v\partial_x$ leads to uniformity in x \rightarrow **hypocoercive**.

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$-L$ is *not* coercive, i.e.:

$$\langle -Lf, f \rangle_H \geq \lambda \|f\|_{L^2(M_T^{-1})}^2 \quad \forall f \in \{f^\infty\}^\perp$$

only holds for $\lambda = 0$; see $f = \rho(x)M_T(v) \Rightarrow$ exponential decay not obvious.

Hypo-coercivity

Definition 1 (Villani 2009)

Consider L on Hilbert space H with $\mathcal{K} = \ker L$; let $\tilde{H} \hookrightarrow \mathcal{K}^\perp$ (densely) (e.g. H ... weighted L^2 , \tilde{H} ... weighted H^1).

– L is called **hypo-coercive** on \tilde{H} if $\exists \lambda > 0, c \geq 1$:

$$\|e^{Lt}f\|_{\tilde{H}} \leq c e^{-\lambda t} \|f\|_{\tilde{H}} \quad \forall f \in \tilde{H}$$

- typically $c > 1$

2 velocity BGK model (= Goldstein-Taylor model)

for $f(x, t) = \begin{pmatrix} f_+(x, t) \\ f_-(x, t) \end{pmatrix}$:

$$\partial_t f_{\pm} \pm \partial_x f_{\pm} = \pm \frac{1}{2}(f_- - f_+), \quad t \geq 0, \quad 2\pi\text{-periodic in } x$$

$$f^{\infty}(x) = \begin{pmatrix} f_+^{\infty} \\ f_-^{\infty} \end{pmatrix} = \begin{pmatrix} \frac{1}{2} \\ \frac{1}{2} \end{pmatrix}$$

x-Fourier modes; discrete velocity basis $\left\{ \begin{pmatrix} 1 \\ 1 \end{pmatrix}, \begin{pmatrix} 1 \\ -1 \end{pmatrix} \right\}$

$$\Rightarrow u_k(t) \in \mathbb{C}^2, \quad k \in \mathbb{Z}$$

Ref's:

[Dolbeault-Mouhot-Schmeiser '15] exp. decay (only 1 conserved quantity)

[Achleitner-AA-Carlen '16] sharp decay rate

2 velocity BGK model: decay of Fourier modes

$$\frac{d}{dt} u_k = -\mathbf{C}_k u_k, \quad \mathbf{C}_k = \begin{pmatrix} 0 & ik \\ ik & 1 \end{pmatrix}, \quad k \in \mathbb{Z}$$

$$\lambda_{\mathbf{C}_0} = 0, \mathbf{1}; \quad \lambda_{\mathbf{C}_k} = \frac{1}{2} \pm i\sqrt{k^2 - \frac{1}{4}}, \quad k \neq 0$$

$$u_0^\infty = \begin{pmatrix} 1 \\ 0 \end{pmatrix}; \quad u_k^\infty = \begin{pmatrix} 0 \\ 0 \end{pmatrix}, \quad k \neq 0$$

Decay of each mode in \mathbf{P}_k -norm:

$$\begin{aligned} \|u_0(t) - u_0^\infty\|_2 &\leq e^{-t} \|u_0(0) - u_0^\infty\|_2 \\ \|u_k(t)\|_{\mathbf{P}_k} &\leq e^{-t/2} \|u_k(0)\|_{\mathbf{P}_k} \quad k \neq 0 \end{aligned}$$

$$\text{Lemma 1} \quad \Rightarrow \quad \mathbf{P}_k \mathbf{C}_k + \mathbf{C}_k^* \mathbf{P}_k \geq 2 \cdot \frac{1}{2} \mathbf{P}_k, \quad k \neq 0$$

2 velocity BGK model: decay of sequence of modes

$$k \neq 0: \quad \mathbf{P}_k := \begin{pmatrix} 1 & \frac{-i}{2k} \\ \frac{i}{2k} & 1 \end{pmatrix} \xrightarrow{|k| \rightarrow \infty} \mathbf{I}$$

\Rightarrow “generic” Lyapunov functional:

$$e(\{u_k\}) := \sqrt{\sum_{k \in \mathbb{Z}} \|u_k\|_{\mathbf{P}_k}^2} \sim \left\| \begin{pmatrix} f_+(x) \\ f_-(x) \end{pmatrix} \right\|_{L^2(0, 2\pi; \mathbb{R}^2)}$$

Theorem 2 (Achleitner-AA-Carlen 2016)

Let $\int_0^{2\pi} [f'_+(x) + f'_-(x)] dx = 2\pi$.

$$\Rightarrow \quad \|f(t) - f^\infty\|_{L^2(0, 2\pi; \mathbb{R}^2)} \leq \sqrt{3} e^{-t/2} \|f^I - f^\infty\|_{L^2(0, 2\pi; \mathbb{R}^2)}, \quad t \geq 0.$$

- Decay rate $\frac{1}{2}$ sharp.
- $\sqrt{3} = \text{cond}(\mathbf{P}_1)$

Continuous velocity BGK model: Fourier modes

$$f_t + vf_x = M_T(v) \int_{\mathbb{R}} f(x, v, t) dv - f(x, v, t), \quad t \geq 0$$

- x -Fourier modes $k \in \mathbb{Z}$ decouple; Hermite function basis in v
 \Rightarrow consider “infinite vector” $\hat{\mathbf{f}}_k(t) \in \ell^2(\mathbb{N}_0)$ for each fixed k :

$$\partial_t \hat{\mathbf{f}}_k + ik\sqrt{T} \mathbf{L}_1 \hat{\mathbf{f}}_k = \mathbf{L}_2 \hat{\mathbf{f}}_k, \quad t \geq 0; k \in \mathbb{Z},$$

$$\mathbf{L}_1 = \begin{pmatrix} 0 & \sqrt{1} & 0 & \cdots \\ \sqrt{1} & 0 & \sqrt{2} & 0 \\ 0 & \sqrt{2} & 0 & \sqrt{3} \\ \vdots & 0 & \sqrt{3} & \ddots \end{pmatrix}, \quad \mathbf{L}_2 = \text{diag}(0, -1, -1, \cdots)$$

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- need: spectral gap of the “infinite matrix” $\mathbf{C}_k := ik\sqrt{T} \mathbf{L}_1 - \mathbf{L}_2$, uniform in $k \in \mathbb{Z}$.
- difficult for ∞ -dimensional case

Continuous velocity BGK model: decay of Fourier modes

→ approximate transformation matrices \mathbf{P}_k : ansatz for upper left 2×2 block:

$$\mathbf{P}_k^{2 \times 2} = \begin{pmatrix} 1 & -i\alpha/k \\ i\alpha/k & 1 \end{pmatrix},$$

remaining diagonal = 1.

Motivation: mix 0^{th} , 1^{st} mode of $\mathbf{L}_2 = \text{diag}(0, -1, \dots)$

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- Find $\alpha \in \mathbb{R}$ to maximize $\mu > 0$ in

$$\mathbf{C}_k^* \mathbf{P}_k + \mathbf{P}_k \mathbf{C}_k - 2\mu \mathbf{P}_k \geq \mathbf{0} \quad \forall k \in \mathbb{Z}.$$

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Lyapunov functional:

$$e(f) := \sum_{k \in \mathbb{Z}} \langle (f_k(v) - M_T(v)), \tilde{\mathbf{P}}_k (f_k(v) - M_T(v)) \rangle_{L^2(M_T^{-1})}$$

$\tilde{\mathbf{P}}_k$... bounded operator on $L^2(M_T^{-1})$, represented by matrix \mathbf{P}_k .

Result: Exp. decay for BGK model; estimated rate off by factor 2.5 (compared to numerical result).

Outline:

- 1 hypocoercivity
- 2 linear BGK models (discrete & continuous velocity) with 1 conserved quantity
- 3 linearized & nonlinear BGK model with 2 conserved quantities

1D nonlinear BGK model

$$f_t(x, v, t) + v f_x = M_f(x, v, t) - f ; \quad x \in \mathbb{T}, \quad v \in \mathbb{R}$$

M_f ... local Maxwellian with **same mass & “temperature”** (=2nd moment) as f ; but mean velocity := zero:

$$M_f(v) = \frac{\rho^{3/2}}{\sqrt{2\pi P}} e^{-\frac{v^2 \rho}{2P}} ; \quad \rho = \int f \, dv, \quad P = \int v^2 f \, dv$$

1D linearized BGK model

- **Linearized model** around temperature T ; $f = M_T + h$;
in x Fourier / in v Hermite expansion:

$$\partial_t \hat{\mathbf{h}}_k(t) + ik\sqrt{T} \mathbf{L}_1 \hat{\mathbf{h}}_k(t) = \mathbf{L}_3 \hat{\mathbf{h}}_k(t), \quad k \in \mathbb{Z};$$

$$\mathbf{L}_3 = \text{diag}(0, -1, 0, -1, -1, \dots), \quad \mathbf{L}_1 = \begin{pmatrix} 0 & \sqrt{1} & 0 & \dots \\ \sqrt{1} & 0 & \sqrt{2} & 0 \\ 0 & \sqrt{2} & 0 & \sqrt{3} \\ \vdots & 0 & \sqrt{3} & \ddots \end{pmatrix}.$$

- **2 conserved quantities**
- \mathbf{L}_1 couples mode 0 to mode 1; mode 2 to mode 3 (and 1)

1D linearized BGK model - simplified Lyapunov functional

⇒ **simplified ansatz** for \mathbf{P}_k :

$$\begin{pmatrix} 1 & -i\alpha/k & 0 & 0 \\ i\alpha/k & 1 & 0 & 0 \\ 0 & 0 & 1 & -i\beta/2k \\ 0 & 0 & i\beta/2k & 1 \end{pmatrix} \xrightarrow{k \rightarrow \infty} \mathbf{I}$$

as its upper-left 4×4 block; rest is \mathbf{I} .

- $\alpha = \beta = \frac{1}{3}$ yields for some $\mu_0 > 0$, with $\mathbf{C}_k := ik\mathbf{L}_1 - \mathbf{L}_3$:

$$\mathbf{P}_k \mathbf{C}_k + \mathbf{C}_k^* \mathbf{P}_k \geq 2\mu_0 \mathbf{P}_k, \quad k\text{-uniformly.}$$

This is the essence for the decay result:

1D linearized BGK model - exponential decay

Lyapunov functional (with $h = f - M_T$):

$$e_\gamma(f) := \sum_{k \in \mathbb{Z}} (1 + k^2)^\gamma \langle h_k(v), \mathbf{P}_k h_k(v) \rangle_{L^2(M_T^{-1})} \sim \|f - M_T\|_{\mathcal{H}^\gamma}^2$$

with $\mathcal{H}^\gamma := H^\gamma(0, 2\pi) \otimes L^2(\mathbb{R}; M_T^{-1})$.

Theorem 3 (Achleitner-AA-Carlen 2016)

Let $T = 1$ and $\frac{1}{2\pi} \int_0^{2\pi} \int_{\mathbb{R}} (1, v^2) f' dv dx = (1, 1)$. Then

$$e_\gamma(f(t)) \leq e^{-t/25} e_\gamma(f') \quad t \geq 0; \quad \forall \gamma \geq 0.$$

Remark: Estimated rate off by factor 18 (compared to numerical result).

1D nonlinear BGK model - local exponential stability

Theorem 4 (Achleitner-AA-Carlen 2016)

Additional assumptions (over Theorem 3): Let $\gamma > \frac{1}{2}$, $\|f^I - M_1\|_{\mathcal{H}^\gamma} < \delta_\gamma$ (=explicit constant). Then

$$e_\gamma(f(t)) \leq e^{-t/25} e_\gamma(f^I), \quad t \geq 0.$$

Proof: Rewrite nonlinear BGK like linearized BGK + higher order remainder

True linearized BGK model

Let $T = 1$, $d = 1$:

$$\partial_t \hat{\mathbf{h}}_k + ik \mathbf{L}_1 \hat{\mathbf{h}}_k = \mathbf{L}_4 \hat{\mathbf{h}}_k, \quad k \in \mathbb{Z};$$

$$\mathbf{L}_4 = \text{diag}(0, 0, 0, -1, -1, \dots)$$

- 3 conserved quantities (mass, mean velocity, temperature)
- iterated hypocoercive structure:

Tridiagonal matrix \mathbf{L}_1 couples mass-mode to velocity-mode to energy-mode to the dissipative mode 3.

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- 3 conserved quantities (mass, mean velocity, temperature)
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Tridiagonal matrix \mathbf{L}_1 couples mass-mode to velocity-mode to energy-mode to the dissipative mode 3.

\Rightarrow simplified ansatz for \mathbf{P}_k (with some $\alpha, \beta, \gamma \in \mathbb{R}$):

$$\begin{pmatrix} 1 & -i\alpha/k & 0 & 0 \\ i\alpha/k & 1 & -i\beta/k & 0 \\ 0 & i\beta/k & 1 & -i\gamma/k \\ 0 & 0 & i\gamma/k & 1 \end{pmatrix}$$

as its upper-left 4×4 block; rest is \mathbf{I} .

\Rightarrow exp. decay (also for $d = 2, 3$)

Conclusion

- algebraic lemma $\mathbf{P}\mathbf{C} + \mathbf{C}^*\mathbf{P} \geq 2\mu\mathbf{P}$ yields the generic decay norm in linear ODEs
- exact / approximate generalizations to kinetic BGK-models: exponential decay for discrete / continuous velocities, linearized / nonlinear

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References

- F. Achleitner, A. Arnold, E. Carlen: On linear hypocoercive BGK models, Springer Proc. in Math. & Stat., 2016.
- F. Achleitner, A. Arnold, E. Carlen: On multi-dimensional hypocoercive BGK models, KRM, 2018.