Transport in Unsteady Flows: Deterministic + Stochastic BIRS, Banff, January 2017

### The Atmosphere and Oceans as Unsteady Flows: Intrinsic Variability and Time-Dependent Forcing

#### **Michael Ghil**

Ecole Normale Supérieure, Paris, and University of California, Los Angeles

Based on joint work with M.D. Chekroun, F.-F. Jin, D. Kondrashov, V. Lucarini, J.D. Neelin, S. Pierini, E. Simonnet, L. Sushama, I. Zaliapin, and many others





Please visit these sites for more info. http://www.atmos.ucla.edu/tcd/ http://www.environnement.ens.fr/

## **Dynamical systems and predictability**

- The initial-value problem  $\rightarrow$  numerical weather prediction (NWP)
  - easiest!
- The asymptotic problem  $\rightarrow$  long-term climate
  - a little harder
- The intermediate problem  $\rightarrow$  low-frequency variability (LFV)
  - multiple equilibria, long-periodic oscillations, intermittency, slow transients, "tipping points"
    - hardest!!

Paraphrasing John von Neumann, in R. L. Pfeffer (ed.), *Dynamics of Climate* (Pergamon, 1960) now re-edited as an Elsevier E-book

## Long-term temperature evolution on Earth

Not only do **global temperatures** move up & down on geological time scales, nor do they just switch from one long-term mean to another: They clearly show **changes in dynamic regime** — from high to low variability, from one dominant periodicity to another, from high to low drift, and so on.



Overall, to model this **complex behavior** we do need to consider both **chaotic** & **random** ingredients, both **intrinsic** & **forced** variability. Compiled by Glen Fergus, <u>https://commons.wikimedia.org/wiki/File%3AAII\_palaeotemps.png</u>

N.B. Plot is ~"log-linear": time axis is logarithmic+linear, temperature axis is linear.

#### **Composite spectrum of climate variability**

Standard treatement of frequency bands:

- 1. High frequencies white noise (or "colored")
- 2. Low frequencies slow evolution of parameters



\*\* 27 years – Brier (1968, *Rev. Geophys.*)

### **Climate and Its Sensitivity**

Let's say CO<sub>2</sub> doubles: How will "climate" change?

- Climate is in stable equilibrium (fixed point); if so, mean temperature will just shift gradually to its new equilibrium value.
- 2. Climate is purely periodic; if so, mean temperature will (maybe) shift gradually to its new equilibrium value.
  But how will the period, amplitude and phase of the limit cycle change?
- 3. And how about some "real stuff" now: chaotic + random?

Ghil (in *Encycl. Global Environmental Change*, 2002)



# Motivation

- The *climate system* is highly *nonlinear and* quite *complex*.
- The system's *major components* the atmosphere, oceans, ice sheets *evolve* on many time and space scales.
- Its *predictive understanding* has to rely on the system's physical, chemical and biological modeling, but also on the thorough mathematical analysis of the models thus obtained: *the forest vs. the trees*.
- The *hierarchical modeling* approach allows one to give proper weight to the understanding provided by the models vs. their realism: back-and-forth between *"toy"* (conceptual) and *detailed* ("realistic") *models*, and between *models* and *data*.
- How do we disentangle *natural variability* from *the anthropogenic forcing*: *can we & should we, or not?*

# The gyres and the eddies

Many scales of motion, dominated in the mid-latitudes by (i) *the double-gyre circulation*; and (ii) *the rings and eddies*.

Much of the focus of physical oceanography over the '70s to '90s has been with the "*meso-scale*": the meanders, rings & eddies, and the associated two-dimensional and quasi-geostrophic *turbulence*.



Based on SSTs, from satellite IR data

# Kuroshio Extension (KE) Path Changes

 Monthly
 36°N

 paths from
 28°N

 paths from
 36°N

 altimeter:
 36°N

 32°N
 36°N

 32°N
 36°N

 Stable vs.
 36°N

 unstable
 36°N

 periods
 36°N

Qiu & Chen (*Deep-Sea Res.*, 2009)



#### Transitions Between Blocked and Zonal Flows in a Barotropic Rotating Annulus with Topography

**Zonal Flow** 13–22 Dec. 1978

#### **Blocked Flow** 10–19 Jan. 1963



Fig. 1. Atmospheric pictures of (A) zonal and (B) blocked flow, showing contour plots of the height (m) of the 700-hPa (700 mbar) surface, with a contour interval of 60 m for both panels. The plots were obtained by averaging 10 days of twice-daily data for (A) 13 to 22 December 1978 and (B) 10 to 19 January 1963; the data are from the National Oceanic and Atmospheric

Administration's Climate Analysis Center. The nearly zonal flow of (A) includes quasi-stationary, small-amplitude waves (32). Blocked flow advects cold Arctic air southward over eastern North America or Europe, while decreasing precipitation in the continent's western part (26).

#### Weeks, Tian, Urbach, Ide, Swinney, & Ghil (Science, 1997)

"Limited-contour" analysis for atmospheric low-frequency variability

10-day sequences of subtropical jet paths: blocked vs. zonal flow regimes

Kimoto & Ghil, JAS, 1993a



FIG. 1. Limited contour analysis of Northern Hemisphere (NH) flows. Daily contours of a prescribed height (2940 m in this case—roughly corresponding to the jet axis) are superimposed for successive 10-day intervals during NH winter 1978/79. Persistence is illustrated by some of the punels (see text for details).

#### **Outline – Unsteady Flows & Climate**

- Atmospheric & oceanic flows
  - scales of motion, in time & space
  - one person's signal ("deterministic") is another one's noise ("stochastic")
- Time-dependent forcing
  - intrinsic vs. forced variability
  - pullback and random attractors
- An illustrative example
  - the Lorenz convection model with time-dependent forcing
- A "grand unification"
  - a mathematical definition of climate sensitivity
- Conclusions and references
  - what do we & don't we know?
  - selected bibliography

#### Atmospheric CO<sub>2</sub> at Mauna Loa Observatory



## **Temperatures and GHGs**

# Greenhouse gases (GHGs) go up, temperatures go up:

It's gotta do with us, at least a bit, doesn't it?



Wikicommons, from Hansen *et al.* (*PNAS*, 2006); see also http://data.giss.nasa.gov/ gistemp/graphs/

# Global warming and its socio-economic impacts

Temperatures rise:

- What about impacts?
- How to adapt?

The answer, my friend, is blowing in the wind, *i.e.*, it depends on the accuracy and reliability of the forecast ...



Source : IPCC (2007), AR4, WGI, SPM Figure SPM.5. Solid lines are multi-model global averages of surface warming (relative to 1980–1999) for the scenarios A2, A1B and B1, shown as continuations of the 20th century simulations. Shading denotes the ±1 standard deviation range of individual model annual averages. The orange line is for the experiment where concentrations were held constant at year 2000 values. The grey bars at right indicate the best estimate (solid line within each bar) and the likely range assessed for the six SRES marker scenarios. The assessment of the best estimate and likely ranges in the grey bars includes the AOGCMs in the left part of the figure, as well as results from a hierarchy of independent models and observational constraints. [Figures 10.4 and 10.29]

#### Time-Dependent Forcing -> Pullback Attractors

Consider the scalar, linear ordinary differential equation (ODE)

 $\dot{x} = -\alpha x + \sigma t \,, \ \alpha > 0 \,, \ \sigma > 0 \,.$ 

When there's no forcing,  $\sigma = 0$ , the ODE is purely dissipative  $\dot{x} = -\alpha x$ , and all solutions converge to the fixed point x = 0 as  $t \to +\infty$ .

Now what about when we do have forcing,  $\sigma \neq 0$ ? At each time  $t = t_1$ , say, we have to "pull back" and start at some time  $s = s_1 \ll t_1$ , say, to see where the flow takes us at  $t = t_1$ . As  $s \to -\infty$ , we get the pullback attractor a = a(t) in the figure,

$$a(t) = \frac{\sigma}{\alpha}(t - \frac{1}{\alpha}).$$



#### Random attractor of the stochastic Lorenz system

#### Snapshot of the random attractor (RA)



- A snapshot of the RA, A(ω), computed at a fixed time t and for the same realization ω; it is made up of points transported by the stochastic flow, from the remote past t T, T >> 1.
- We use multiplicative noise in the deterministic Lorenz model, with the classical parameter values b = 8/3,  $\sigma = 10$ , and r = 28.
- Even computed pathwise, this object supports meaningful statistics.



A day in the life of the Lorenz (1963) model's random attractor, or LORA for short; see SI in Chekroun, Simonnet & Ghil (2011, *Physica D*)

### **Classical Strange Attractor**

Physically closed system, modeled mathematically as autonomous system: neither deterministic (anthropogenic) nor random (natural) forcing.

The attractor is strange, but still fixed in time ~ "irrational" number.

Climate sensitivity ~ change in the average value (first moment) of the coordinates (*x*, *y*, *z*) as a parameter  $\lambda$  changes.



## **Random Attractor**

Physically open system, modeled mathematically as non-autonomous system: allows for deterministic (anthropogenic) as well as random (natural) forcing.

The attractor is "pullback" and evolves in time ~ "imaginary" or "complex" number.

Climate sensitivity ~ change in the statistical properties (first and higher-order moments) of the attractor as one or more parameters ( $\lambda$ ,  $\mu$ , ...) change.

Ghil (*Encyclopedia of Atmospheric Sciences*, 2<sup>nd</sup> ed., 2012)



#### How to define climate sensitivity or, What happens when there's natural variability?

This definition allows us to watch how "the earth moves," as it is affected by both natural and anthropogenic forcing, in the presence of natural variability, which includes both chaotic & random behavior: chaotic + random behavior:



Clearly the invariant measure  $\nu(t;\mu)$  changes in its position (i.e., its support), as well as in its probability density — with time *t*, as shown here — but also with respect to an arbitrary parameter  $\mu$ , where  $\mu = \tau$  in the present case. Hence, in general,  $\gamma = \partial d_W / \partial \mu$ .

## Yet another (grand?) unification

Lorenz (JAS, 1963)

Climate is deterministic and autonomous, but highly nonlinear.

Trajectories diverge exponentially,

forward asymptotic PDF is multimodal.

Hasselmann (*Tellus*, 1976) Climate is stochastic and noise-driven, but quite linear. Trajectories decay back to the mean,

forward asymptotic PDF is unimodal.

#### **Grand unification (?)**

Climate is deterministic + stochastic, as well as highly nonlinear. Internal variability and forcing interact strongly, change and sensitivity refer to both mean and higher moments.



#### **Concluding remarks –**

#### What do we & don't we know?

#### What do we know?

- It's getting warmer.
- We do contribute to it.
- So we should act as best we know and can!

#### What do we know less well?

- By how much?
  - Is it getting warmer ...
  - Do we contribute to it ...
  - How does the climate system (atmosphere, ocean, ice, etc.) really work?
  - How does natural variability interact with anthropogenic forcing?

#### What to do?

- Better understand the system and its forcings.
- Explore the models', and the system's
  - robustness and sensitivity
  - pullback & random attractors

#### **Some general references**

- Arnold, L., 1998: Random Dynamical Systems, Springer Monographs in Math., Springer, 625 pp.
- Arnol'd, V. I., 1983: *Geometrical Methods in the Theory of Ordinary Differential Equations*, Springer-Verlag, New York/Heidelberg/Berlin, 334 pp.
- Charney, J.G., *et al.*, 1979: *Carbon Dioxide and Climate: A Scientific Assesment*. National Academic Press, Washington, D.C.
- Chekroun, M. D., E. Simonnet, and M. Ghil, 2011: Stochastic climate dynamics: Random attractors and time-dependent invariant measures, *Physica D*, **240**, 1685–1700.
- Ghil, M., and S. Childress, 1987: *Topics in Geophysical Fluid Dynamics: Atmospheric Dynamics, Dynamo Theory and Climate Dynamics*, Ch. 5, Springer-Verlag, New York, 485 pp.
- Ghil, M., 2001: Hilbert problems for the geosciences in the 21<sup>st</sup> century, *Nonlin. Processes Geophys.*, **8**, 211–222.
- Ghil, M., 2017: The wind-driven ocean circulation: Applying dynamical systems theory to a climate problem, *Discr. Cont. Dyn. Syst. A*, **37(1)**, 189–228, doi:<u>10.3934/dcds.2017008</u>.
- Ghil, M., M.D. Chekroun, and E. Simonnet, 2008: Climate dynamics and fluid mechanics: Natural variability and related uncertainties, *Physica D*, **237**, 2111–2126.
- Hasselmann, K., 1976: Stochastic climate models. Part I. Theory, Tellus, 28, 474-485.
- Houghton, J.T., G.J. Jenkins, and J.J. Ephraums (Eds.), 1991: *Climate Change, The IPCC Scientific Assessment*, Cambridge Univ. Press, Cambridge, MA, 365 pp.
- Lorenz, E.N., 1963: Deterministic nonperiodic flow. J. Atmos. Sci., 20, 130-141.
- Lucarini, V., and F. Ragone, 2011: Energetics of climate models: Net energy balance and meridional enthalpy transport, *Rev. Geophys.*, **49**, RG1001, doi: <u>10.1029/2009RG000323</u>
- Ruelle, D., 1997: Application of hyperbolic dynamics to physics: Some problems and conjectures, *Bull. Amer. Math. Soc.*, **41**, 275–278.

# **Reserve slides**

# Unfortunately, things aren't all that easy!

What to do?

Try to achieve better interpretation of, and agreement between, models ...

**Ghil, M**., 2002: Natural climate variability, in *Encyclopedia of Global Environmental Change*, T. Munn (Ed.), Vol. 1, Wiley Natural variability introduces additional complexity into the anthropogenic climate change problem

The most common interpretation of observations and GCM simulations of climate change is still in terms of a scalar, linear Ordinary Differential Equation (ODE)



Linear response to CO<sub>2</sub> vs. observed change in T

Hence, we need to consider instead a <u>system of nonlinear</u> <u>Partial Differential Equations (PDEs)</u>, with parameters and multiplicative, as well as additive forcing (deterministic + stochastic)

 $\frac{dX}{dt} = N(X, t, \mu, \beta)$ 

## So what's it gonna be like, by 2100?

Table SPM.2. Recent trends, assessment of human influence on the trend and projections for extreme weather events for which there is an observed late-20th century trend. {Tables 3.7, 3.8, 9.4; Sections 3.8, 5.5, 9.7, 11.2–11.9}

Phenomenon <sup>a</sup> and direction of trend	Likelihood that trend occurred in late 20th century (typically post 1960)	Likelihood of a human contribution to observed trend <sup>b</sup>	Likelihood of future trends based on projections for 21st century using SRES scenarios
Warmer and fewer cold days and nights over most land areas	Very likely <sup>o</sup>	Likely <sup>d</sup>	Virtually certain <sup>d</sup>
Warmer and more frequent hot days and nights over most land areas	Very likely <sup>e</sup>	Likely (nights)⁴	Virtually certain <sup>d</sup>
Warm spells/heat waves. Frequency increases over most land areas	Likely	More likely than not <sup>r</sup>	Very likely
Heavy precipitation events. Frequency (or proportion of total rainfall from heavy falls) increases over most areas	Likely	More likely than not	Very likely
Area affected by droughts increases	Likely in many regions since 1970s	More likely than not	Likely
Intense tropical cyclone activity increases	Likely in some regions since 1970	More likely than not	Likely
Increased incidence of extreme high sea level (excludes tsunamis)9	Likely	More likely than not <sup>th</sup>	Likely <sup>i</sup>

## F. Bretherton's "horrendogram" of Earth System Science



Earth System Science Overview, NASA Advisory Council, 1986

## Global warming and its socio-economic impacts- II

#### Temperatures rise:

- What about impacts?
- How to adapt?

#### AR5 vs. AR4

A certain air of *déjà vu*: GHG "scenarios" have been replaced by "representative concentration pathways" (RCPs), more dire predictions, but the uncertainties remain.

Source : IPCC (2013), AR5, WGI, SPM



## But deterministic chaos doesn't explain all: there are many other sources of irregularity!

- The energy spectrum of the atmosphere and ocean is "full": all space & time scales are active and they all contribute to forecasting uncertainties.
- Still, one can imagine that the longest & slowest scales contribute most to the longest-term forecasts.
- "One person's signal is another person's noise."



After Nastrom & Gage (JAS, 1985)

#### Climate models (atmospheric & coupled) : A classification



- Partial unidirectional asynchronous, hybrid
- Full

Hierarchy: back-and-forth between the simplest and the most elaborate model, and between the models and the observational data

#### **Multiple scales of motion: Space-time organization**

- The most active scales lie along a diagonal in this space vs. time plot.
- Why this is so is far from clear as of now.
- > We'll deal with weather first, then climate.



**Observations** 

N.B. A high-variability ridge lies close to the diagonal of the plot (cf. also Fraedrich & Böttger, 1978, JAS)

\* LFV ≈ 10–100 days (intraseasonal)

# Can we, nonlinear dynamicists, help?

The uncertainties might be *intrinsic*, rather than mere "tuning problems"

If so, maybe *stochastic structural stability* could help!

Might fit in nicely with recent taste for "stochastic parameterizations"



The DDS dream of structural stability (from Abraham & Marsden, 1978)

#### Random Dynamical Systems (RDS), I - RDS theory

This theory is the counterpart for randomly forced dynamical systems (RDS) of the *geometric theory* of ordinary differential equations (ODEs). It allows one to treat stochastic differential equations (SDEs) — and more general systems driven by noise — as flows in (phase space)×(probability space). SDE~ODE, RDS~DDS, L. Arnold (1998)~V.I. Arnol'd (1983).

#### Setting:

- (i) A phase space *X*. **Example**:  $\mathbb{R}^n$ .
- (ii) A probability space (Ω, F, P). Example: The Wiener space Ω = C<sub>0</sub>(R; R<sup>n</sup>) with Wiener measure P.
- (iii) A model of the noise  $\theta(t) : \Omega \to \Omega$  that preserves the measure  $\mathbb{P}$ , i.e.  $\theta(t)\mathbb{P} = \mathbb{P}; \theta$  is called the driving system. **Example:**  $W(t, \theta(s)\omega) = W(t + s, \omega) - W(s, \omega);$ it starts the noise at *s* instead of t = 0.
- (iv) A mapping  $\varphi : \mathbb{R} \times \Omega \times X \to X$  with the cocycle property. **Example**: The solution operator of an SDE.

. . . . . . . . . . . . . .

-

#### Random Dynamical Systems (RDS), I - RDS theory

This theory is the counterpart for randomly forced dynamical systems (RDS) of the *geometric theory* of ordinary differential equations (ODEs). It allows one to treat stochastic differential equations (SDEs) — and more general systems driven by noise — as flows in (phase space) $\times$ (probability space).

SDE~ODE, RDS~DDS, L. Arnold (1998)~V.I. Arnol'd (1983).

#### Setting:

- (i) A phase space X. **Example**:  $\mathbb{R}^n$ .
- (ii) A probability space (Ω, F, ℙ). Example: The Wiener space Ω = C<sub>0</sub>(ℝ; ℝ<sup>n</sup>) with Wiener measure ℙ.
- (iii) A model of the noise  $\theta(t) : \Omega \to \Omega$  that preserves the measure  $\mathbb{P}$ , i.e.  $\theta(t)\mathbb{P} = \mathbb{P}; \theta$  is called the driving system. **Example:**  $W(t, \theta(s)\omega) = W(t + s, \omega) - W(s, \omega);$ it starts the noise at *s* instead of t = 0.
- (iv) A mapping  $\varphi : \mathbb{R} \times \Omega \times X \to X$  with the cocycle property. **Example**: The solution operator of an SDE.

#### **RDS, II - A Geometric View of SDEs**



- $\varphi$  is a random dynamical system (RDS)
- $\Theta(t)(x,\omega) = (\theta(t)\omega, \varphi(t,\omega)x)$  is a flow on the bundle

イロン 不良 とくほう 不良 とうしょう

#### RDS, III- Random attractors (RAs)

A random attractor  $A(\omega)$  is both *invariant* and "pullback" *attracting*:

- (a) Invariant:  $\varphi(t,\omega)\mathcal{A}(\omega) = \mathcal{A}(\theta(t)\omega)$ .
- (b) Attracting:  $\forall B \subset X$ ,  $\lim_{t\to\infty} \operatorname{dist}(\varphi(t, \theta(-t)\omega)B, \mathcal{A}(\omega)) = 0$  a.s.

Pullback attraction to  $A(\omega)$ 



#### Sample measures for an NDDE model of ENSO

#### The Galanti-Tziperman (GT) model (JAS, 1999)

$$\frac{dT}{dt} = -\epsilon_T T(t) - M_0(T(t) - T_{sub}(h(t))),$$

$$\begin{split} h(t) &= M_1 e^{-\epsilon_m (\tau_1 + \tau_2)} h(t - \tau_1 - \tau_2) & \text{models for EN} \\ &- M_2 \tau_1 e^{-\epsilon_m (\frac{\tau_1}{2} + \tau_2)} \mu(t - \tau_2 - \frac{\tau_1}{2}) T(t - \tau_2 - \frac{\tau_1}{2}) \\ &+ M_3 \tau_2 e^{-\epsilon_m \frac{\tau_2}{2}} \mu(t - \frac{\tau_2}{2}) T(t - \frac{\tau_2}{2}). \end{split}$$

Seasonal forcing given by  $\mu(t) = 1 + \epsilon \cos(\omega t + \phi).$ The pullback attractor and its invariant measures were computed.

Figure shows the changes in the mean,  $2^{nd} \& 4^{th}$  moment of h(t), along with the Wasserstein distance  $d_W$ , for changes of 0–5% in the delay parameter  $\tau_{\kappa,0}$ 

Note intervals of both smooth & rough dependence!

Neutral delay-differential equation (NDDE), derived from Cane-Zebiak and Jin-Neelin models for ENSO: *T* is East-basin SST and *h* is thermocline depth.



#### letters to nature

Nature 350, 324 - 327 (1991); doi:10.1038/350324a0

#### Interdecadal oscillations and the warming trend in global temperature time series

M. Ghil & R. Vautard

THE ability to distinguish a warming trend from natural variability is critical for an understanding of the climatic response to increasing greenhouse-gas concentrations. Here we use singular spectrum analysis<sup>1</sup> to analyse the time series of global surface air tem-peratures for the past 135 years<sup>2</sup>, allowing a secular warming trend and a small number of oscillatory modes to be separated from the noise. The trend is flat until 1910, with an increase of 0.4 °C since then. The oscillations exhibit interdecadal periods of 21 and 16 years, and interannual periods of 6 and 5 years. The interannual oscillations are probably related to global aspects of the El Niño-Southern Oscillation (ENSO) phenomenon<sup>3</sup>. The interdecadal oscillations could be associated with changes in the extratropical ocean circulation<sup>4</sup>. The oscillatory components have combined (peak-to-peak) amplitudes of 0.2 °C, and therefore limit our ability to predict whether the inferred secular warming trend of 0.005 °Cyr<sup>-1</sup> will continue. This could postpone incontrovertible detection of the greenhouse warming signal for one or two decades.



# How important are different sources of uncertainty?

Varies, but typically no single source dominates.





Courtesy Tim Palmer, 2009

# Parameter dependence – I

It can be smooth or it can be rough: Niño-3 SSTs from intermediate coupled model for ENSO (Jin, Neelin & Ghil, 1994, 1996)

Skewness & kurtosis of the SSTs: time series of 4000 years,

$$\Delta \delta = 3 \cdot 10^{-4}$$

$$\delta = 0.9557$$





#### M. Chekroun (work in progress)

# Pullback attractor and invariant measure of the GT model

The time-dependent pullback attractor of the GT model supports an invariant measure  $\nu = \nu(t)$ , whose density is plotted in 3-D perspective.

The plot is in delay coordinates h(t+1) vs. h(t) and the density is highly concentrated along 1-D filaments and, furthermore, exhibits sharp, near–0-D peaks on these filaments.

The Wasserstein distance *d*<sub>W</sub> between one such configuration, <sup>0.4</sup> at given parameter values, and <sup>0.2</sup> another one, at a different set of values, is proportional to the work <sup>0</sup> <sup>30</sup> needed to move the total probability mass from one configuration to the other.

Climate sensitivity  $\gamma\,$  can be defined as  $\gamma=\partial d_{\rm W}/\partial \tau$ 



## **Exponential divergence vs. "coarse graining"**

The classical view of dynamical systems theory is: positive Lyapunov exponent → trajectories diverge exponentially

But the presence of multiple regimes implies a much more structured behavior in phase space

Still, the probability distribution function (pdf), when calculated forward in time, is pretty smeared out

L. A. Smith (Encycl. Atmos. Sci., 2003)



## **Global warming and "global weirding"**

#### "CLIMATE STRANGE

FORGET GLOBAL WARMING—AND GET READY for GLOBAL WEIRDING BY BRYAN WALSH"

TIME MAGAZINE, Dec. 29, 2014 - Jan. 5, 2015

"The New Rule: For the next few (?) years, global warming will lead to colder, more brutal winters."



- Oh, thank you for the latest prediction from a science journalist based on interesting but still rather tentative, & hotly debated, suggestions from a few media-loving (& vice-versa) researchers.
- And if this is so certain, why wasn't it predicted by IPCC<sup>(\*)</sup> and other models BEFORE it happened?

<sup>(\*)</sup> Intergovernmental Panel on Climate Change

#### SSA (prefilter) + (low-order) MEM



In good agreement with MTM peaks of **Ghil & Vautard (1991,** *Nature***)** for the Jones *et al.* (1986) temperatures & stack spectra of Vautard *et al.* (1992, *Physica D*) for the IPCC "consensus" record (both global), to wit 26.3, 14.5, 9.6, 7.5 and 5.2 years.

Peaks at 27 & 14 years also in Koch sea-ice index off Iceland (Stocker & Mysak, 1992), etc. Plaut, Ghil & Vautard (1995, Science)

# **Concluding remarks, I – RDS and RAs**

#### **Summary**

- A change of paradigm from closed, autonomous systems to open, non-autonomous ones.
- Random attractors are (i) spectacular, (ii) useful, and (iii) just starting to be explored for climate applications.

#### Work in progress

- Study the effect of specific stochastic parametrizations on model robustness.
- Applications to intermediate models and GCMs.
- Implications for climate sensitivity.
- Implications for predictability?

## Concluding remarks, II – Climate change & climate sensitivity

#### What do we know?

- It's getting warmer.
- We do contribute to it.
- So we should act as best we know and can!

#### What do we know less well?

- By how much?
  - Is it getting warmer ...
  - Do we contribute to it ...
- How does the climate system (atmosphere, ocean, ice, etc.) really work?
- How does natural variability interact with anthropogenic forcing?

#### What to do?

- Better understand the system and its forcings.
- Explore the models', and the system's, robustness and sensitivity
  - stochastic structural and statistical stability!
  - linear response = response function + susceptibility function!!