

Dissipative Solitons (DSs) and the FitzHugh-Nagumo (FN) Equation

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Workshop

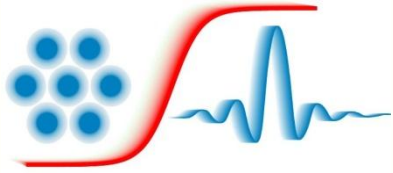
Localized Multi-Dimensional Patterns in Dissipative Systems:

Theory, Modelling, and Experiments

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Banff 25.07. – 29.07.2011

Extended Version



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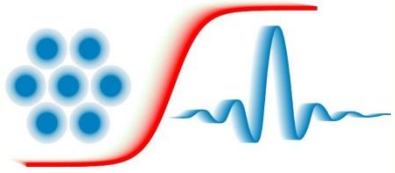
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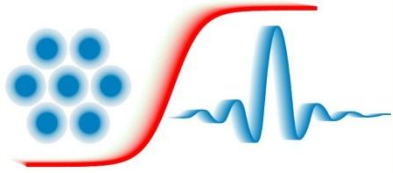
Key reference:

H.-G. Purwins, H. U. Bödeker and Sh. Amiranashvili,
***Dissipative solitons*, Advances in Physics, vol. 59, pp. 485-701 (2010)**

<http://www.uni-muenster.de/Physik.AP/Purwins/Research-Summary>.



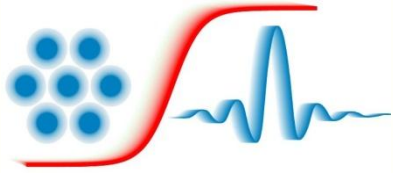
1. What is a Dissipative Soliton (DS)?



The Particle Concept

the particle concept

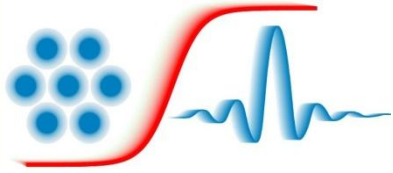
- **division of objects of natural perception into subunits**
- **homogeneous space**
- **small number, simple (short range) interaction**
- **most successful example: concept of the atom**
- **usually particles exist in closed systems**
- **particle description by ordinary differential equations**
- **other point of view: localized solitary deviation of some field variable on an otherwise homogeneous background**
- **relatively new aspect: localized structures in certain classes of dissipative systems behave like particles in many respect**



Dissipative Solitons (DSs) I:

The Ideal Object

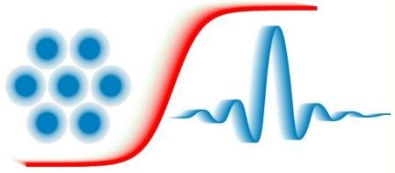
- DSs are localized deviations of one or more state variable from an otherwise stationary homogeneous background.
- DSs or ensembles of them are attractors of a stationary or periodically driven spatially extended dissipative systems that are homogeneous by construction.
- Provided they are sufficiently far away from each other and from the boundary, for a given set of parameters, DSs can show up in any number in one or more distinct classes with same size, shape and dynamical properties within a given class.
- DSs coexist with some stable stationary homogeneous state that coincides with the background state far away from any DS.
- DSs undergo interaction, such that their individuality is retained to large extent, alternatively they are annihilated or generated as a whole.



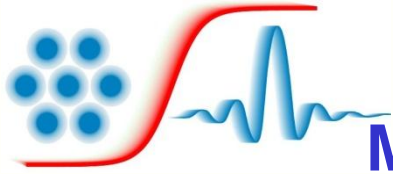
Dissipative Solitons (DSs) II:

The Real Object

- Usually DSs exist only in some 1- or 2-dimensional subspace.
- DSs interact (weakly) with the boundary and with each other.
- Perturbations are present in the form of temporal and spatial noise.
- Well defined inhomogeneities may be present e.g. in the form of gradients as well as local or periodic, distortions.
- The background may be weakly modulated periodically by self-organization.
- One may deal with a discrete system.
- In the presence of global restrictions the background may depend weakly on the number of LSs.
- Possibly, in the case of different classes of DSs interaction may give rise to a change from one to another class.

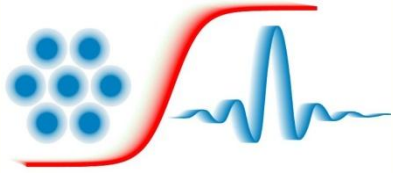


2. The Emergence of Dissipative Solitons (DSs)



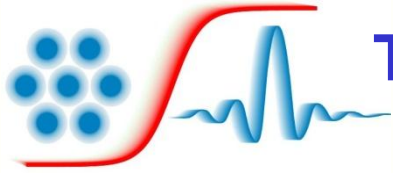
Detection of Nerve Fibres and Measurement of the Speed of Pulse Propagation

- **Hermann Ludwig Ferdinand von Helmholtz (1812 - 1894)**
medical doctor and physicist
- perhaps first quantitative experimental characterization of a DS
- **1842 : detection of nerve fibres connecting ganglia**
(clusters of cells)
- **1850: first measurement of speed of pulse propagation**
on a nerve fibre for a frog and human being: $v \approx 30\text{m/s}$
- because of the slow propagation Helmholtz thought, that
some substance should propagate



Nerve Pulse Propagation and Reaction-Diffusion Systems

- **Robert Thomas Dietrich Luther (1868 – 1945)**
theoretical chemist, scholar of Oswald
- **1906: meeting of the Bunsengesellschaft**
- **presentation of various experiments demonstrating waves in chemical reaction systems being homogeneous by construction**
- **Presentation of the result of the theoretical analysis of a reaction-diffusion equation (which equation he analysed is not known)**
c = speed of propagation
D = diffusion constant
 τ = time constant of the chemical reaction
$$c = \text{const.} \times \sqrt{D/\tau}$$
- **in the discussion with Nernst, Luther speculated that nerve pulses may be the result of the propagation of chemical reactions**



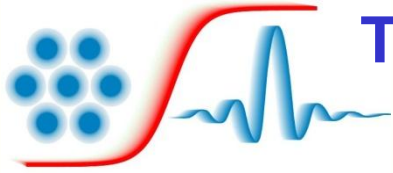
The Zeldovich-Frank-Kamenetsky (ZF) Equation I: The Equation

$$u_t = Du_{xx} + f(u),$$

- showed up in the 1930s in relation to genetic diffusion (Fischer (1937); Kolmogorov, Petrovsky, Piskunoff (1937))
- Zeldovich and Frank-Kamenetsky (1938) investigated the equation with

$$f(u) = \lambda u - u^3 + \kappa_1, \quad \lambda = \frac{1}{\tau}, \quad \lambda \geq 0$$

- equation is referred to as the ZF equation
- 1940: discussion of nerve pulse propagation in relation to the Zeldovich-Frank-Kamenetsky equation



The Zeldovich-Frank-Kamenetsky (ZFK) Equation II: Analytical Expression for Kink Solutions in \mathcal{R}^1

condition: two stable stationary solution of the system

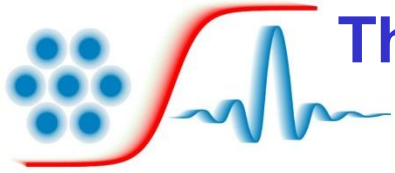
$$u_t = Du_{xx} + \lambda u - u^3 + \kappa_1$$

$$u(x, t) = \frac{1}{2} \left[(u_+ + u_-) + (u_+ - u_-) \tanh k(x - ct) \right]$$

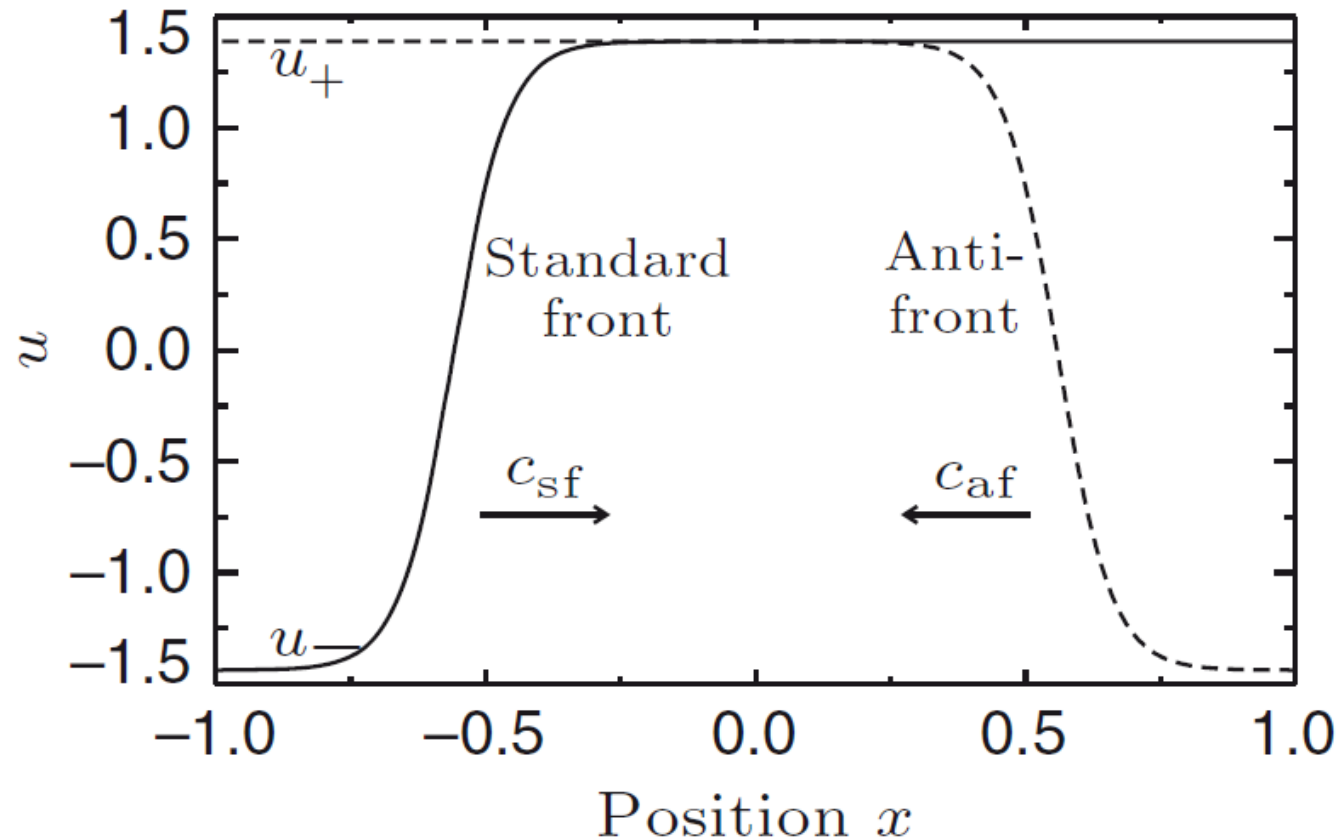
$$c = -(3/2)\sqrt{2D}(u_+ + u_-), \quad k = (u_+ - u_-) / \sqrt{8D}$$

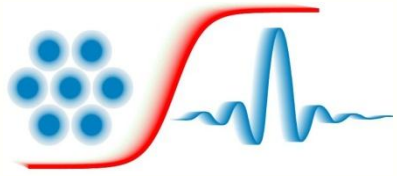
$$|c|_{\max} = \sqrt{\frac{3}{2}D\lambda} = \text{const.} \times \sqrt{\frac{D}{\tau}}, \quad \tau = \frac{1}{\lambda}$$

$u_+, u_- =$ largest and smallest zeros of $f(u)$ in the case of a total of three zeros of $f(u)$
zeros depend on λ (> 1) and κ_1

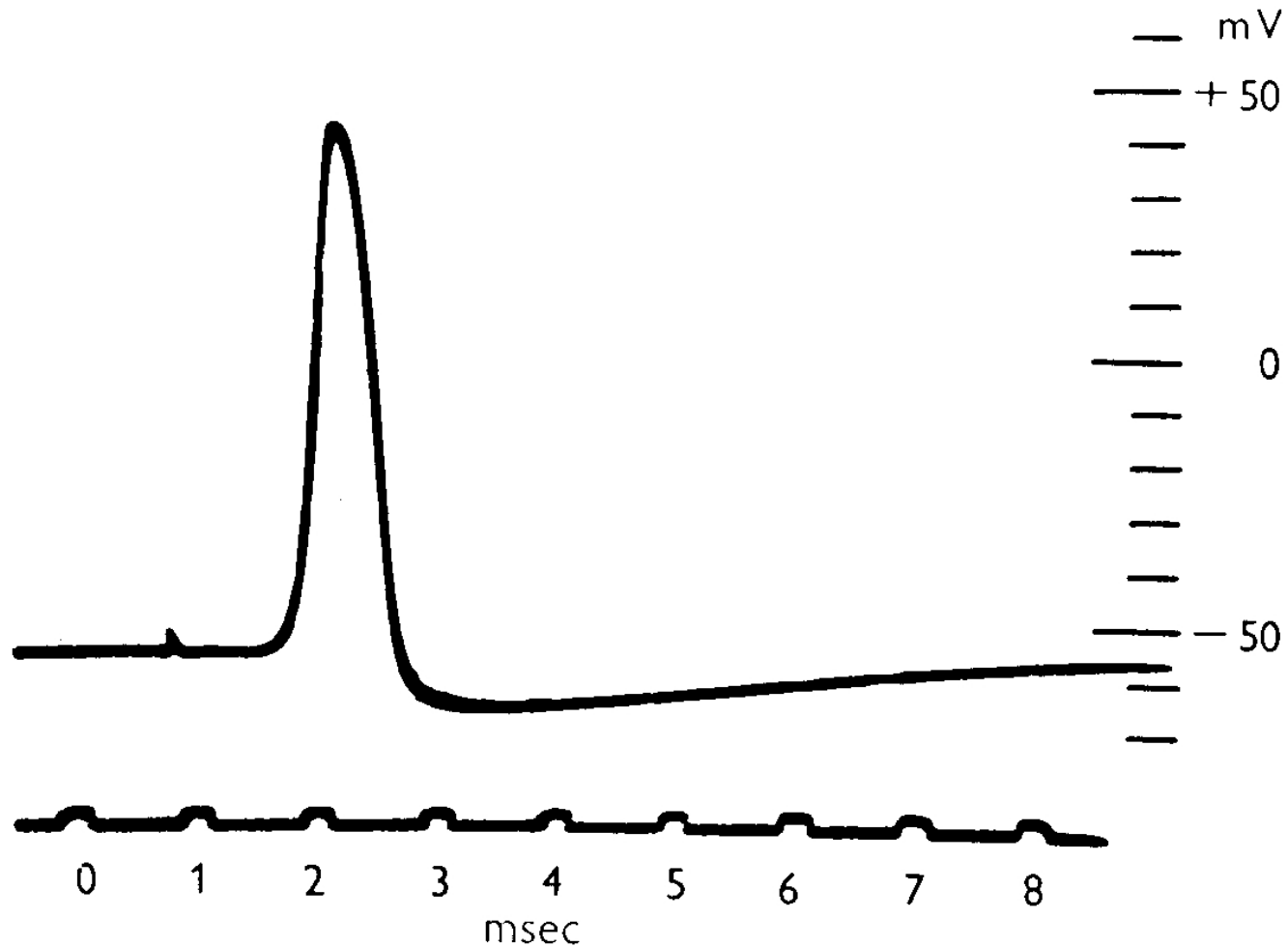


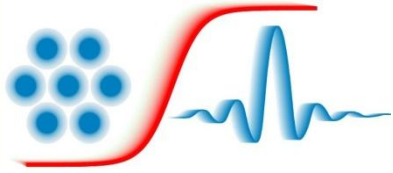
The Zeldovich-Frank-Kamenetsky (ZFK) Equation IV: Illustration of Kinks and Localized Solutions in \mathcal{R}^1





Nerve Pulse of an Octopus: Electrical Potential at a Given Position of the Fibre as a Function of Time





Properties of Nerve Fibers

- **Nerve fiber (axon)**

connected to neuron via the axon hillock for outgoing signals

diameter = 0.5 – 10 μm

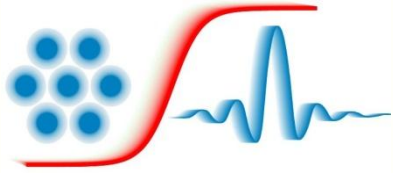
length = 1 μm – 1 m and more

speed of pulses propagation $v \approx 1 – 100 \text{ m/s}$

- **Dendrites**

connecting nerve fibers via synapses for incoming signals

up to 100 000 to 200 000 per neuron



The FitzHugh-Nagumo (FN) Equation I: The Equation

- **1952: Hodgkin and Huxley set up a system of 5 equations to describe nerve pulse excitation in the fibre**
- **1962: FitzHugh and Nagumo reduced the system to two variables**

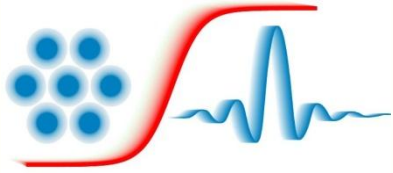
$$u_t = d_u^2 u_{xx} + \lambda u - u^3 - v + \kappa_1,$$

$$\tau v_t = d_v^2 v_{xx} + u - v$$

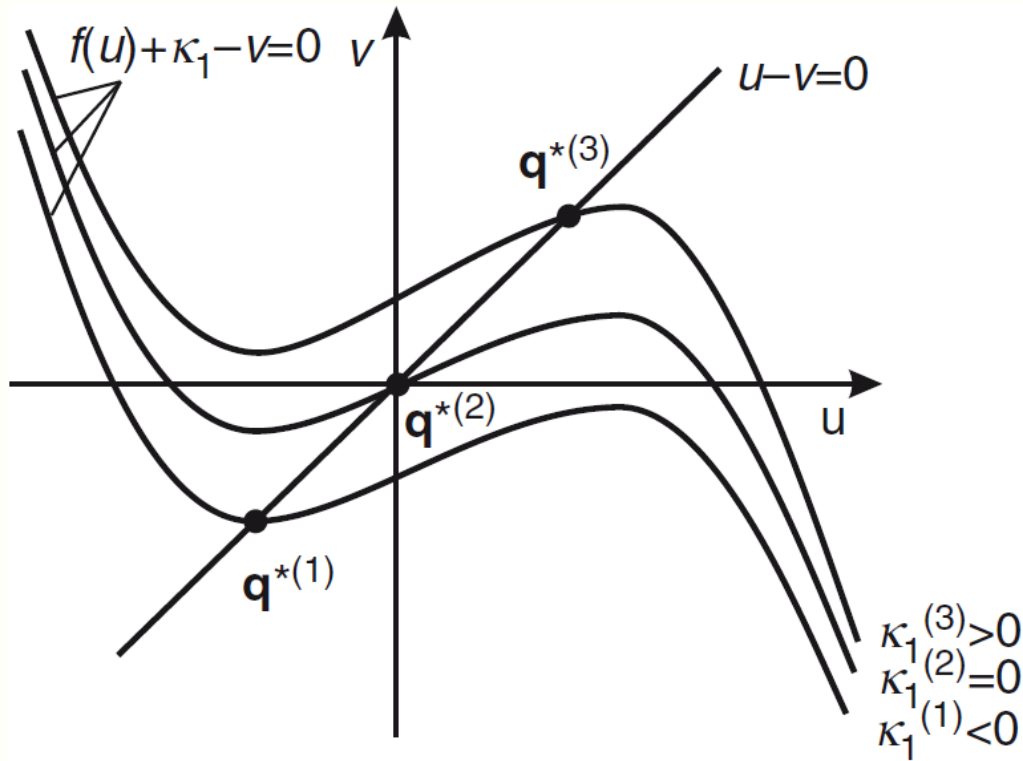
$$\lambda \geq 0, \tau \gg 1$$

u = transmembrane voltage

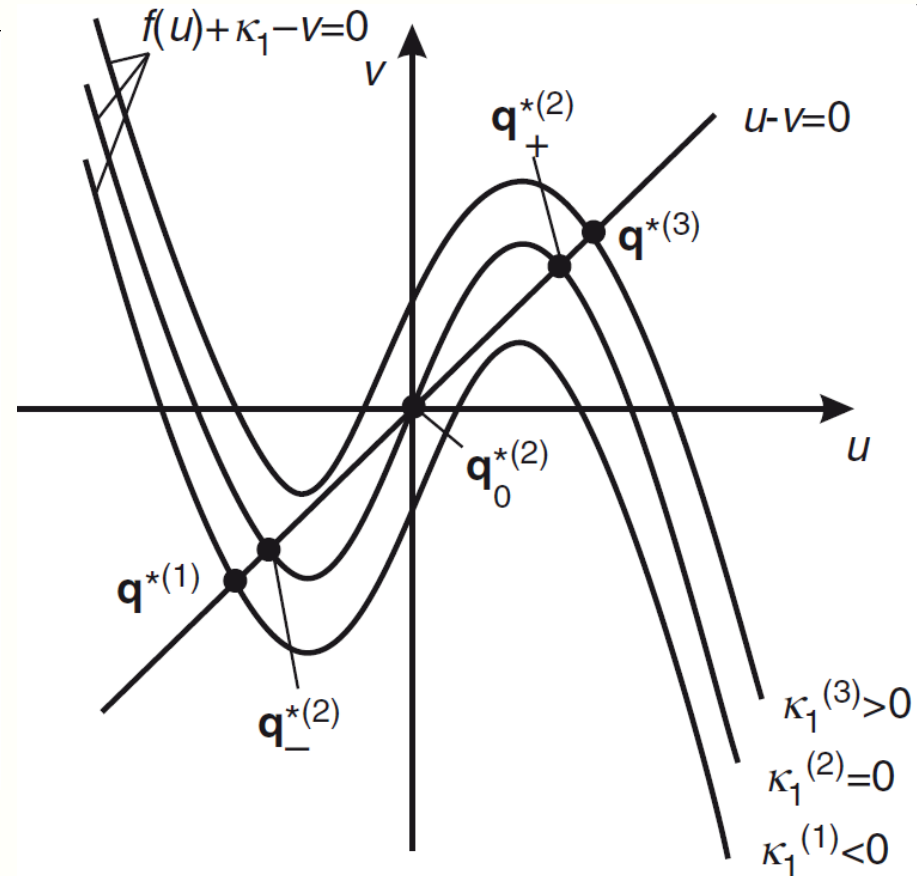
v = recovery variable (related to the potassium current)



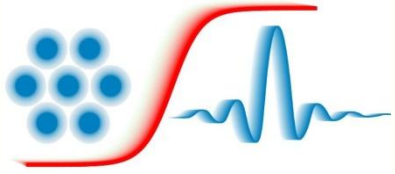
The FitzHugh-Nagumo (FN) Equation II: The Space Clamped System



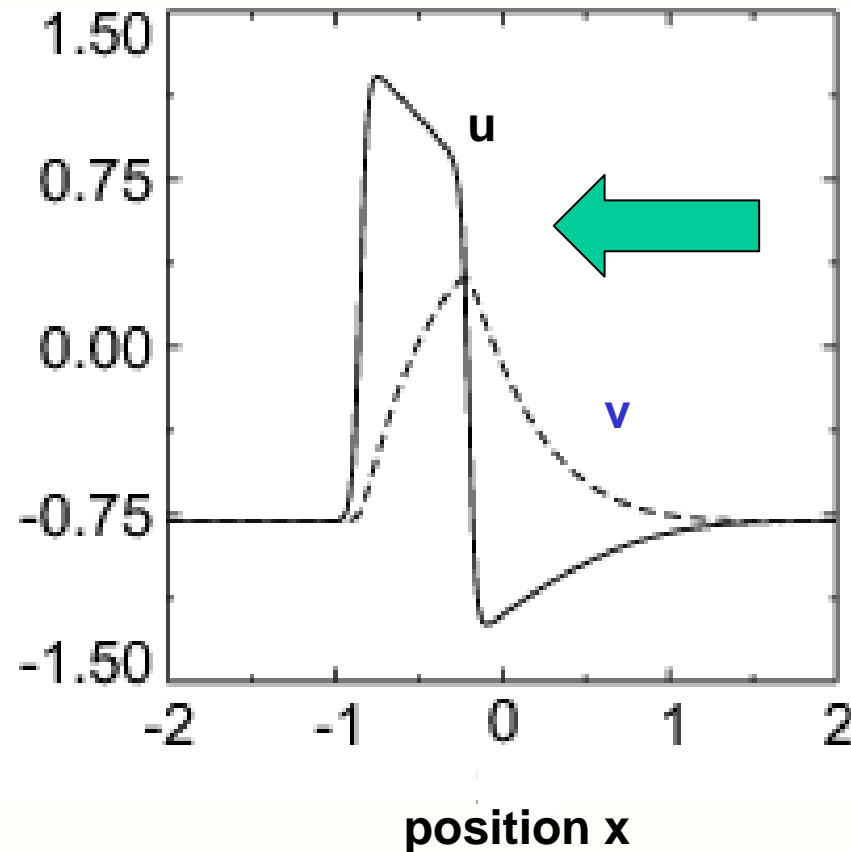
$\lambda < 1$

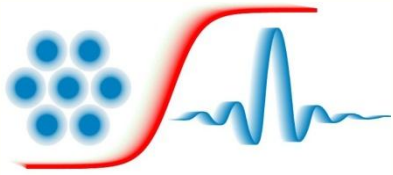


$\lambda > 1$



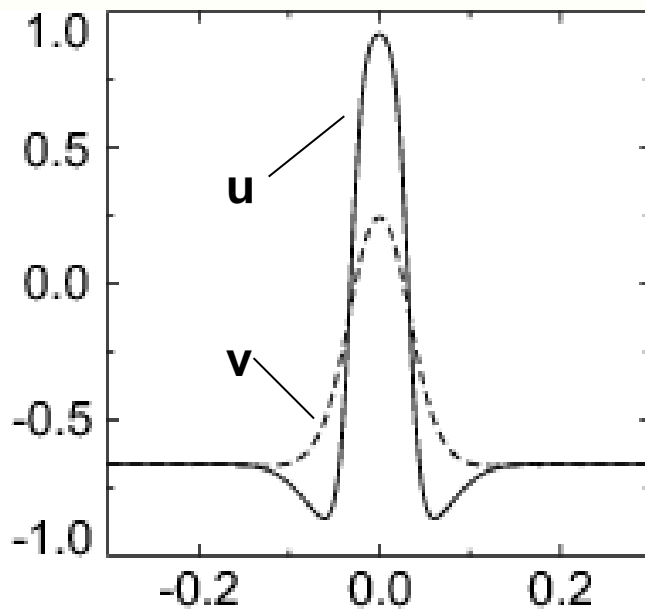
The FitzHugh-Nagumo (FN) Equation III: Example for a Numerical Solution for a Travelling DS in \mathcal{R}^1 (Nerve Pulse)





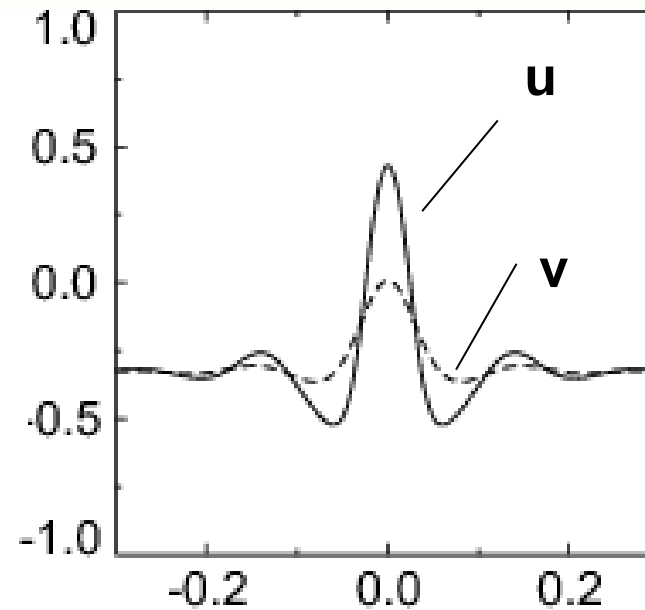
The FitzHugh-Nagumo (FN) Equation IV: Examples for Numerical Solutions for Stationary DSs in \mathcal{R}^1

non-oscillatory tails

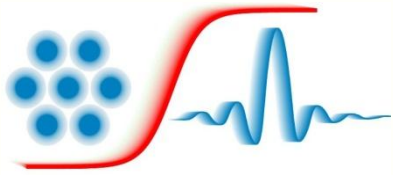


position x

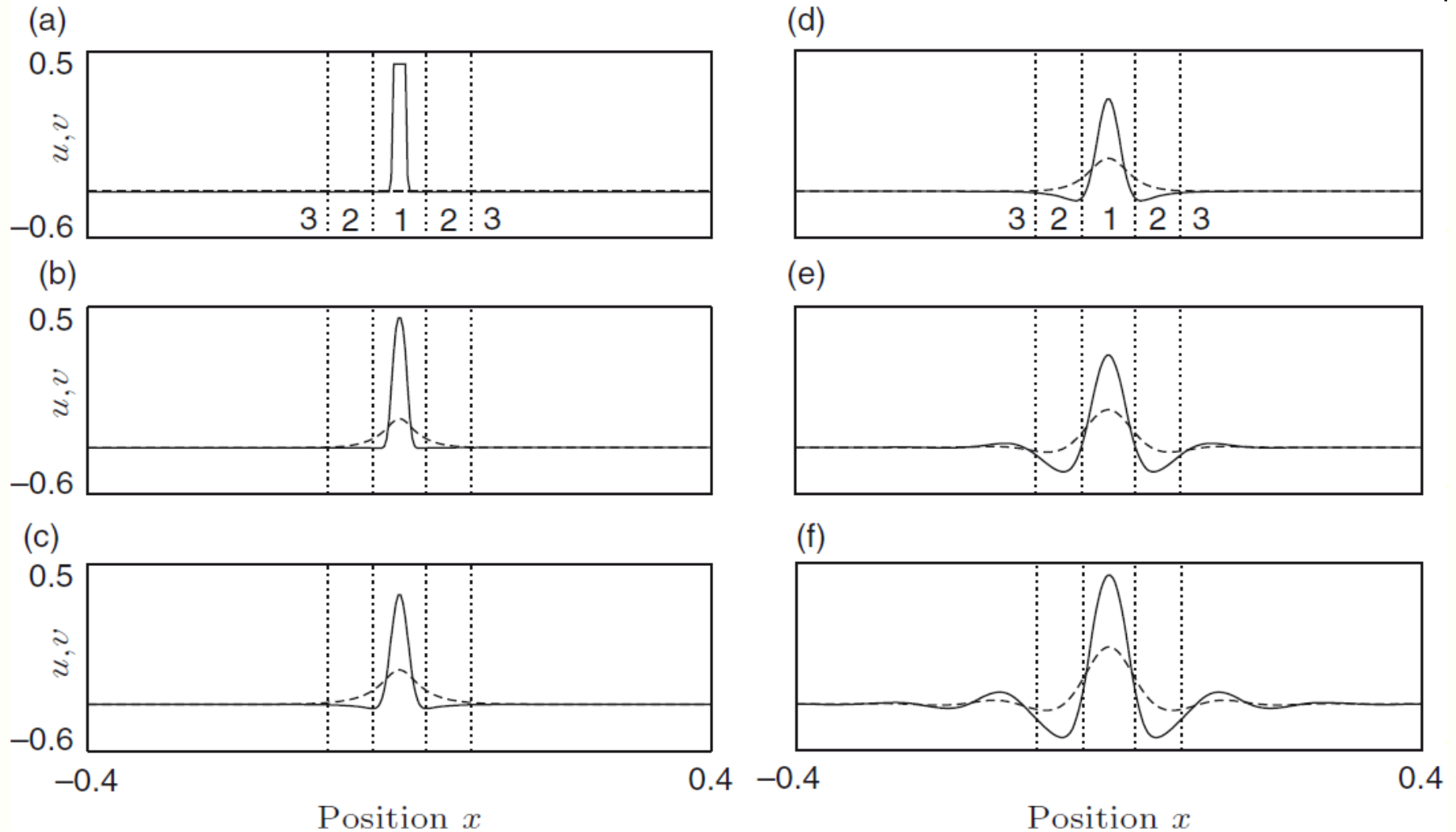
oscillatory tails

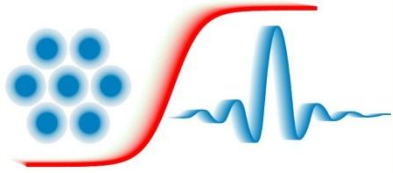


position x

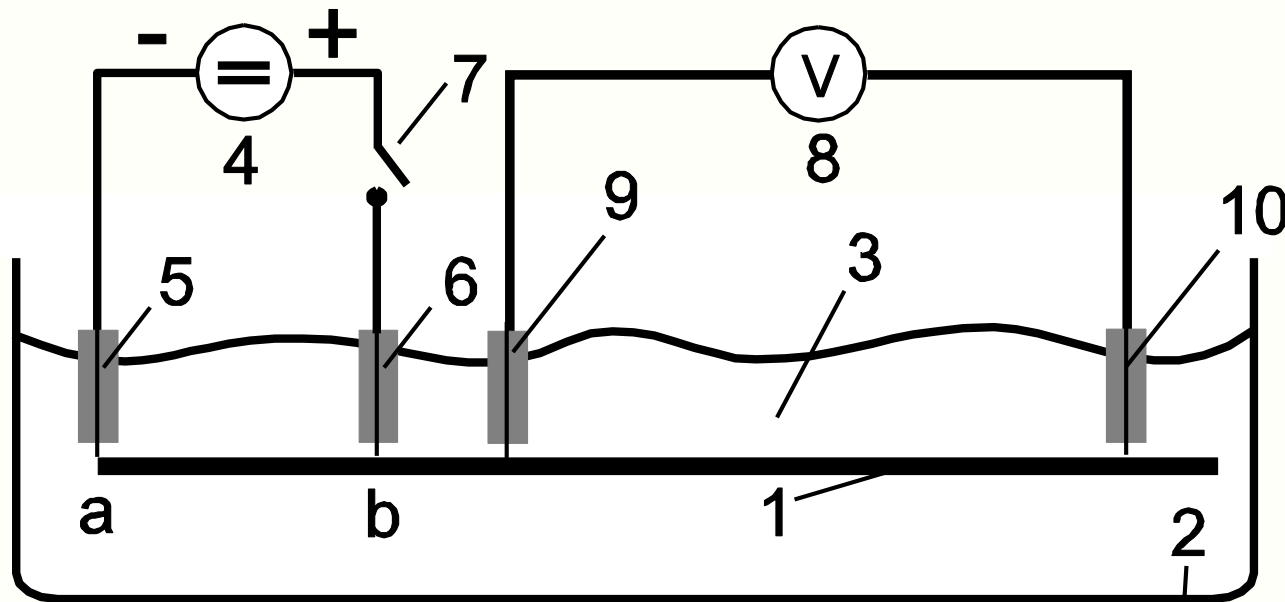


The FitzHugh-Nagumo (FN) Equation V: Illustration of the Formation of a Stationary DS in \mathcal{R}^1

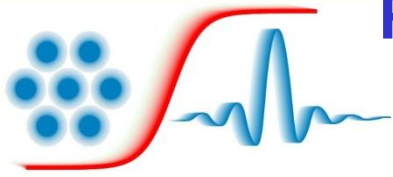




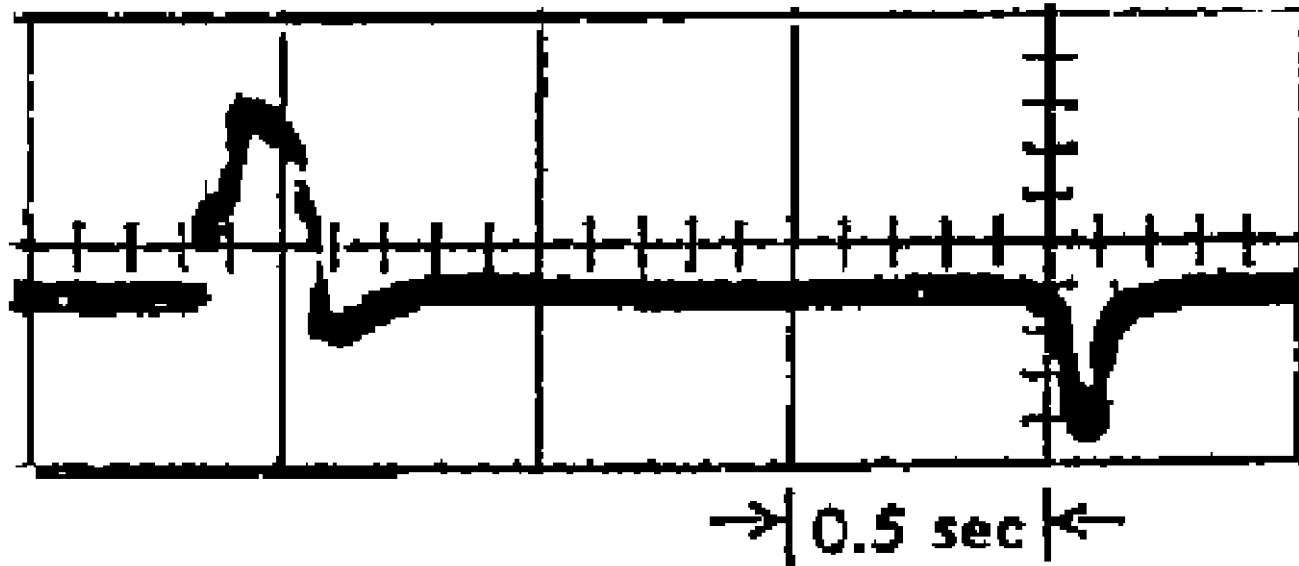
Pulse Propagation on an Iron Wire in Nitric Acid I: Experimental Set-Up



1. iron wire 2. dish 3. nitric acid 4. voltage source 5., 6. Pt electrodes
7. switch 8. potential meter 9., 10. fixed potential probes

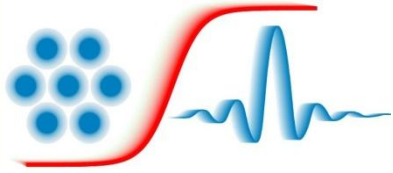


Pulse Propagation on an Iron Wire in Nitric Acid II: Electrical Potential in Dependence of Time at Fixed Position of the Wire



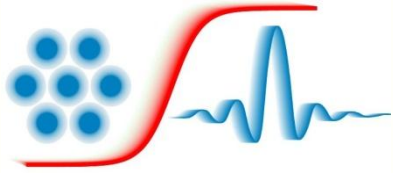
potential difference between the probes 9. and 10. as a
function of time

propagation of a well defined section of naked iron on an otherwise
oxidised wire after an electrical perturbation initially deoxidising the
wire locally

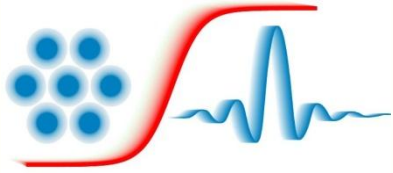


Remark on Russian Contributions

- at the end of the 80th s of the last century interesting contributions
- Rosanov in optics and here in relation with wave equation
- Osipov and Kerner in the field of electrical transport systems
- in particular the work of the Osipov and Kerner is in close relation to the 2-component FN equation



3. A Special Class of Electrical Transport Systems and the Generalize FitzHugh-Nagumo (FN) Equation



FN Equation with Additional Global Coupling

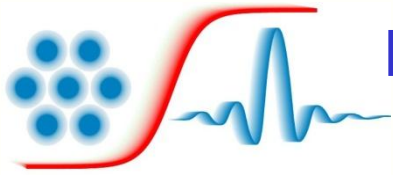
$$u_t = d_u^2 u_{xx} + f(u) - v + \kappa_1 \left[-\kappa_2 \int_{\Omega} u dx \right],$$

$$\tau v_t = d_v^2 v_{xx} + u - v,$$

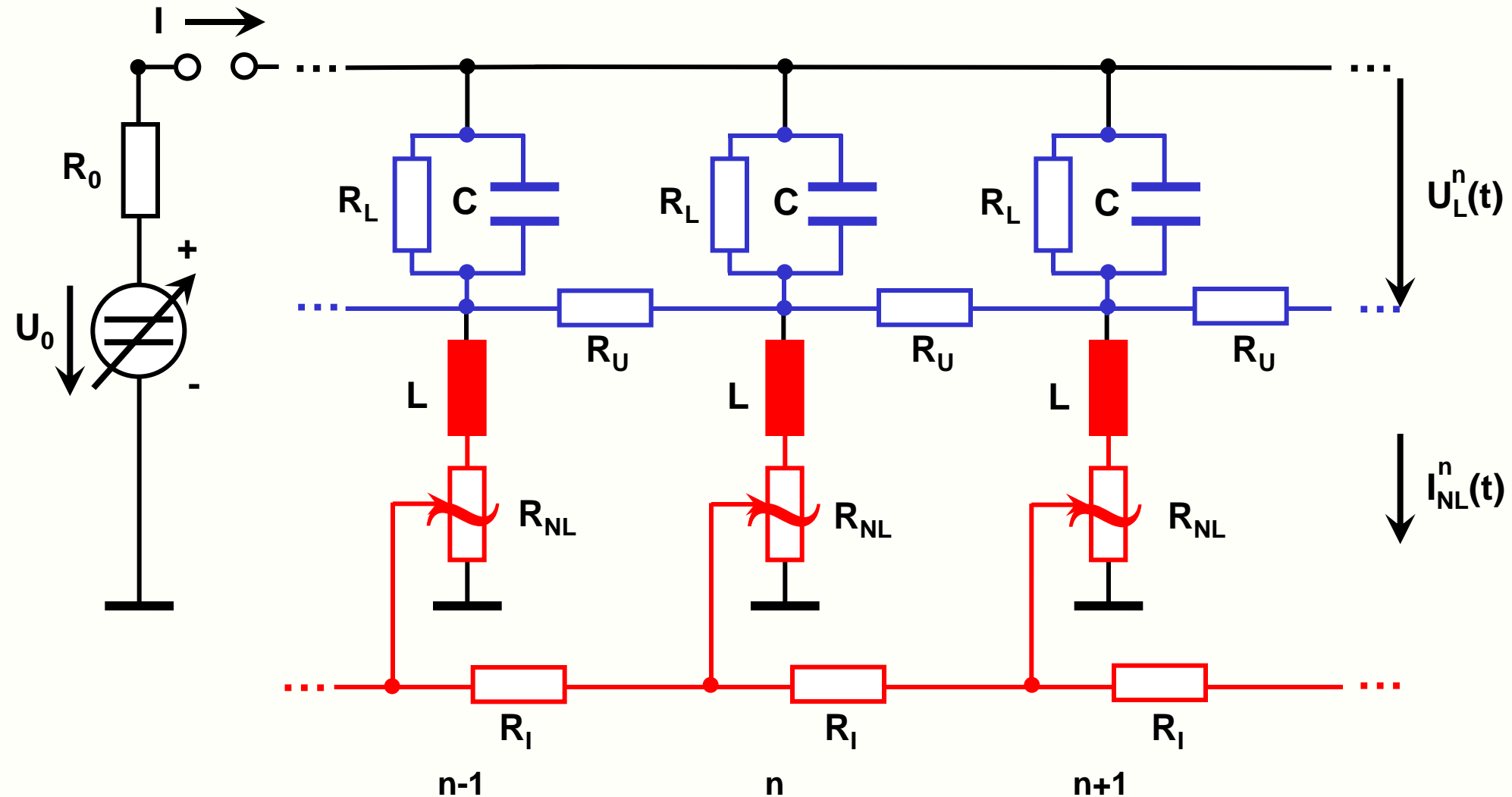
$$u = u(x, t), \quad v = v(x, t),$$

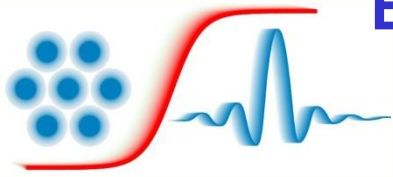
$$f(u) = \lambda u - u^3,$$

$$d_u, d_v, \tau, \kappa_2, \lambda \geq 0$$



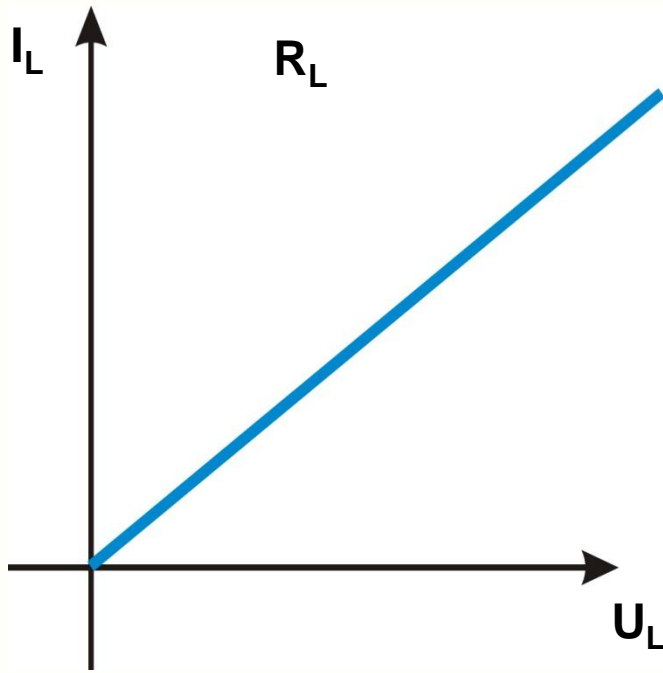
Electrical Equivalent Circuit for the FN Equation I: The 1-Dimensional Circuit



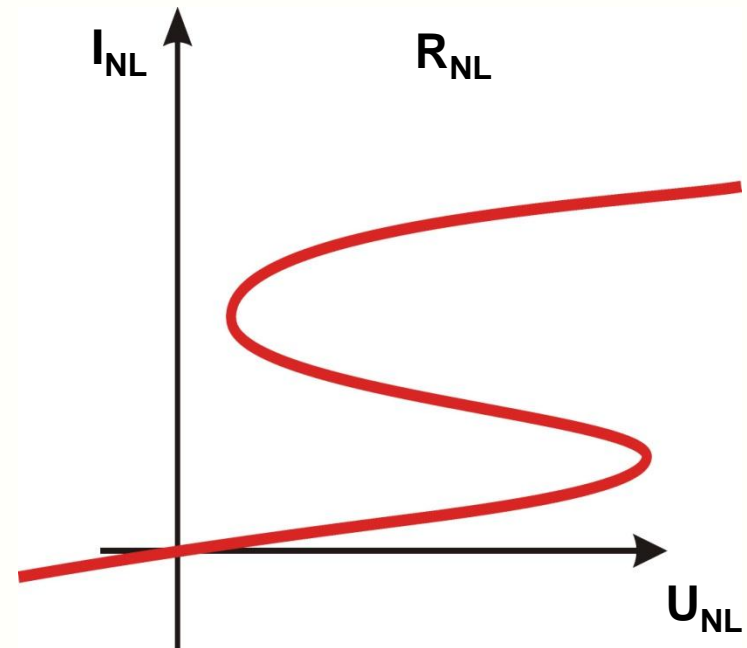


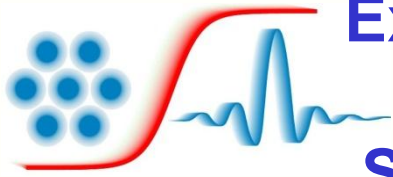
Electrical Equivalent Circuit for the FN Equation II: (Current)-(Voltage) Characteristics of the Linear and the Nonlinear Resistance

monotonic
(approximative linear)

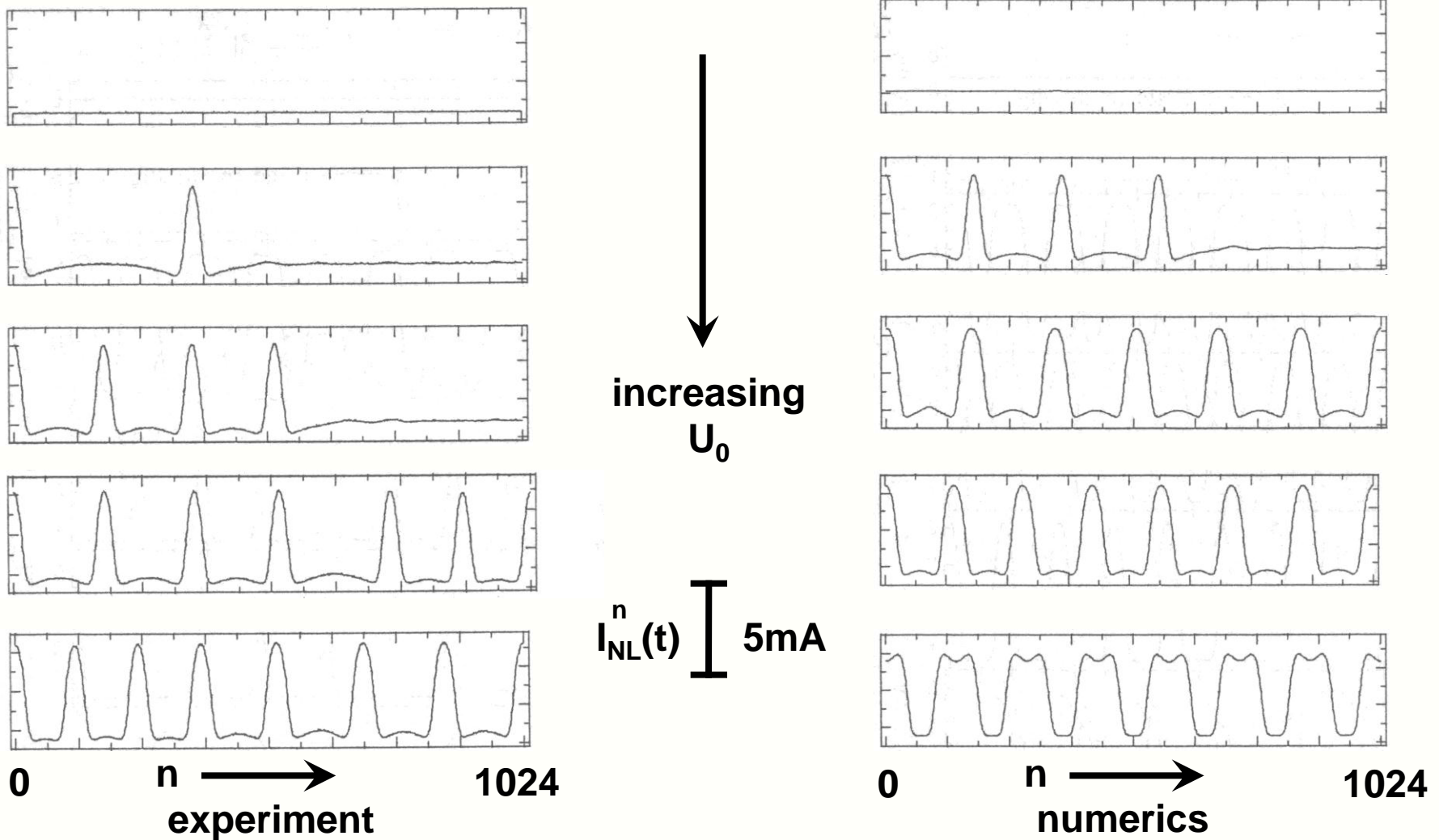


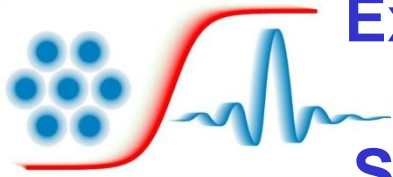
S-shaped
(strongly nonlinear)



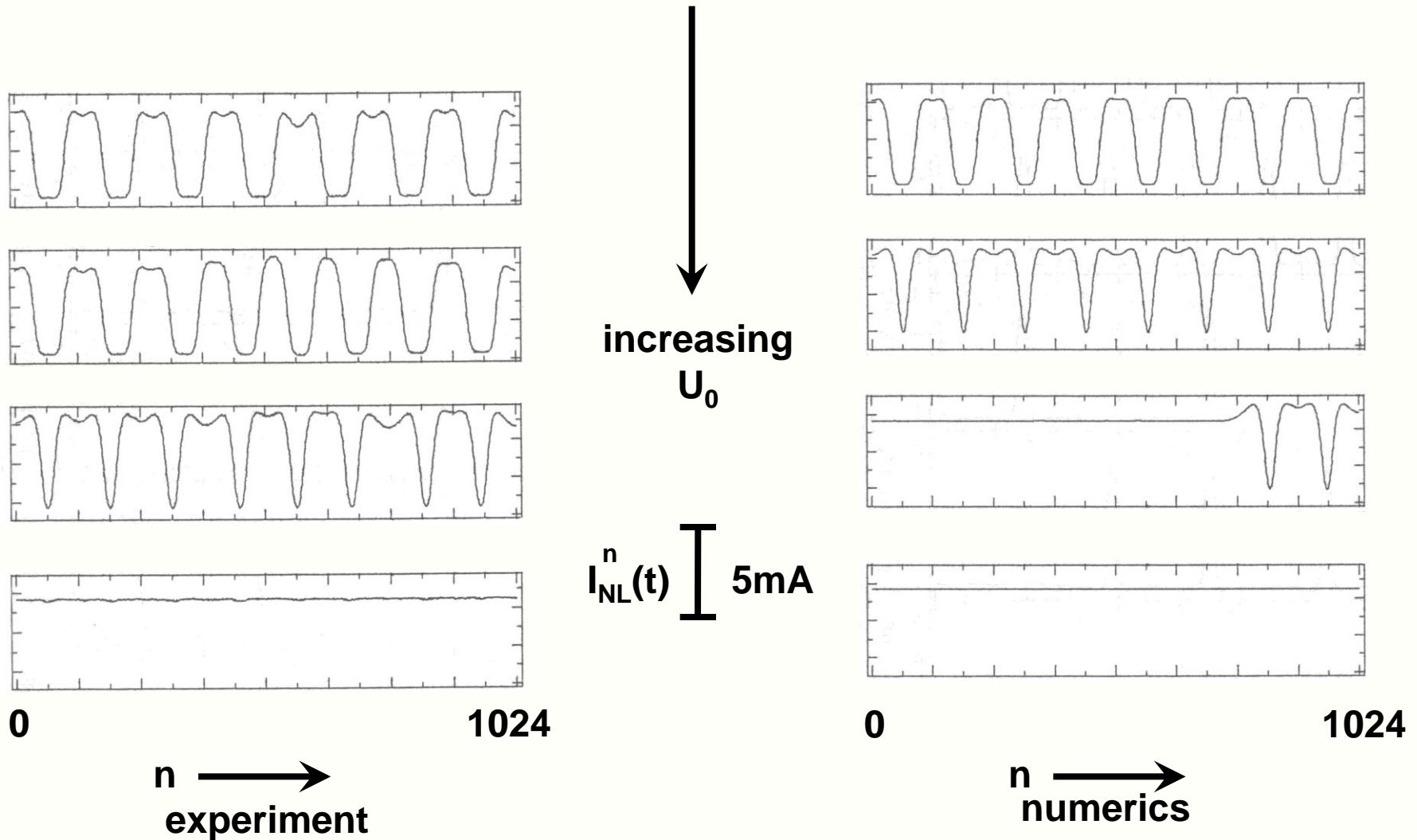


Experimental Result for the 1-Dimensional Electrical Equivalent Circuit and Comparison to Numerical Solutions of the Corresponding Network Equation I



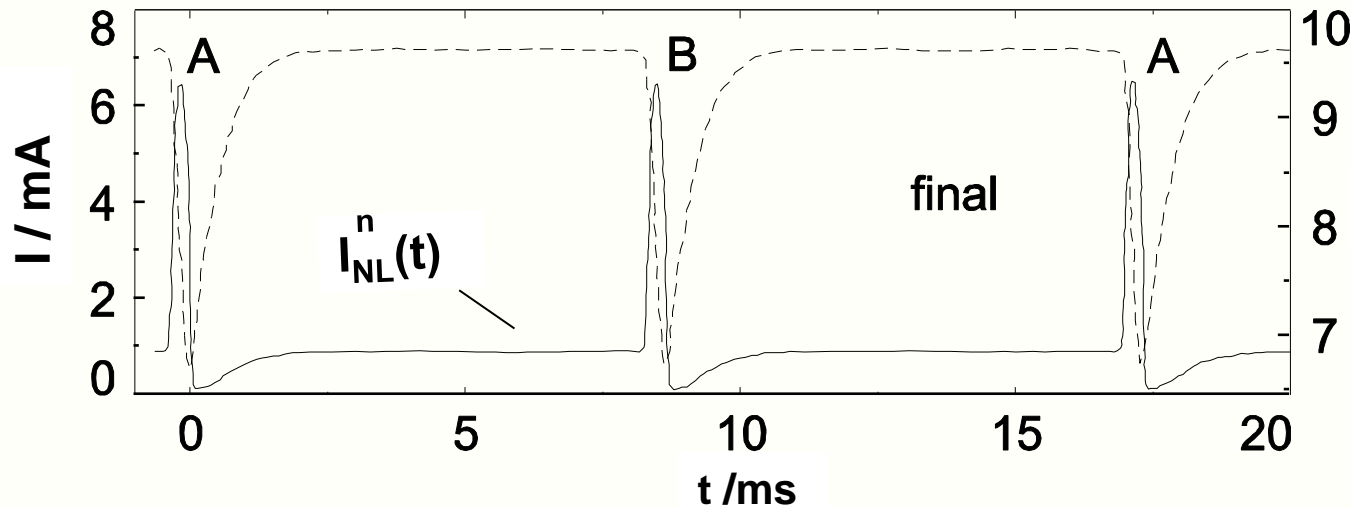
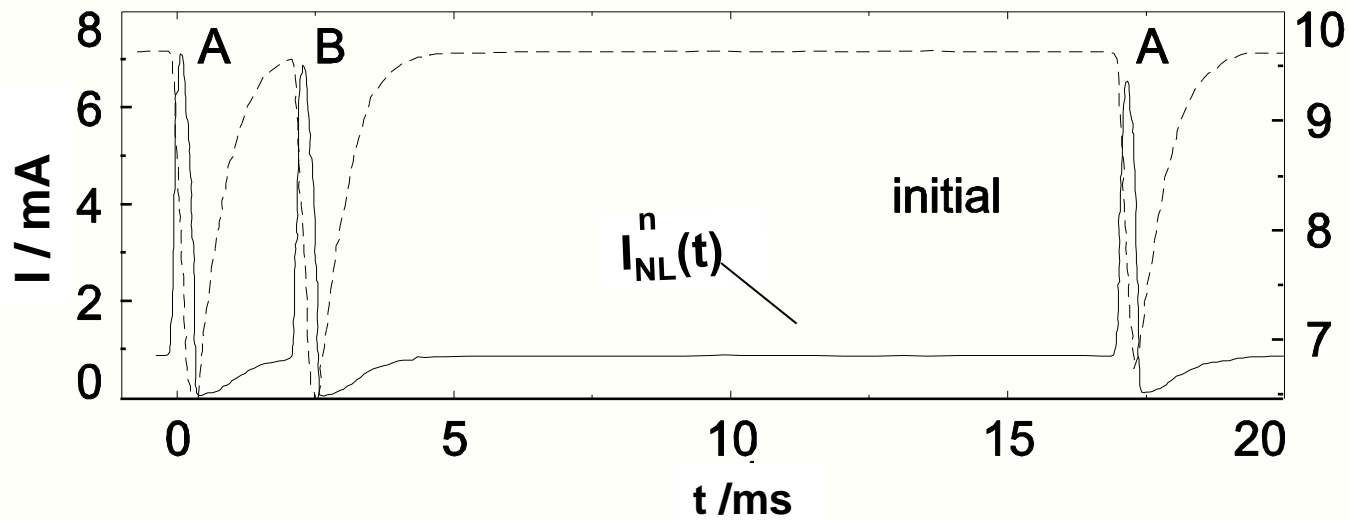


Experimental Result for the 1-Dimensional Electrical Equivalent Circuit and Comparison to Numerical Solutions of the Corresponding Network Equation II

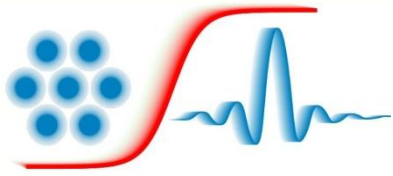




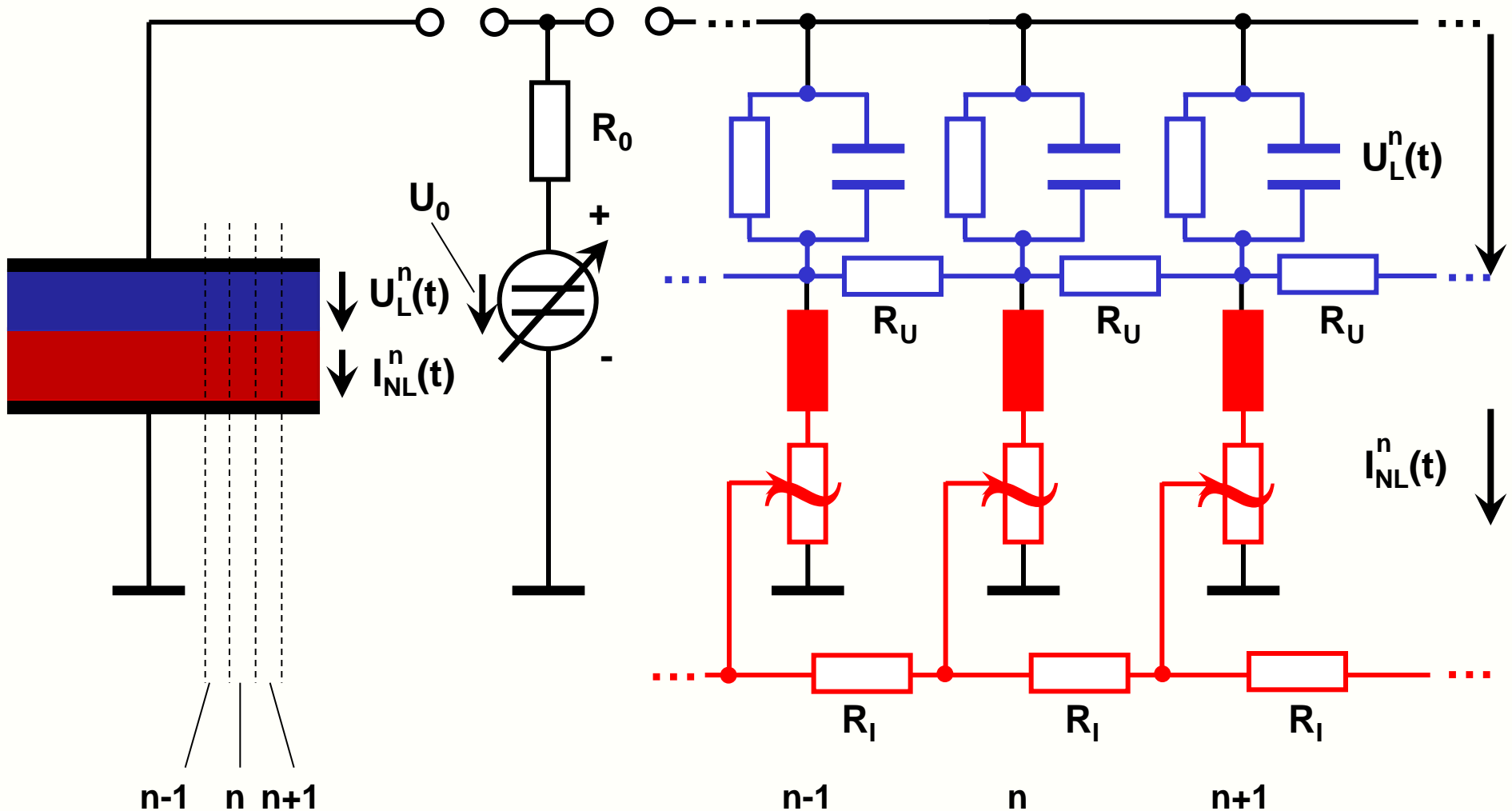
Experimental Results for the 1-Dimensional Electrical Equivalent Circuit Closed to a Ring: Repulsive Interaction of Two Travelling DSs



the current (continuous line) is measured at a given cell; initially two DSs are ignited close to each other and travel in the same direction; in the course of time the interaction leads to separation, finally corresponding to an angle of 180° on the ring



The 1-Dimensional Network Equation as a Model Equation for a Double Layer Continuum System I: The System



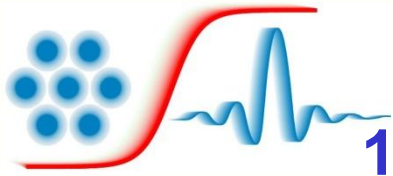
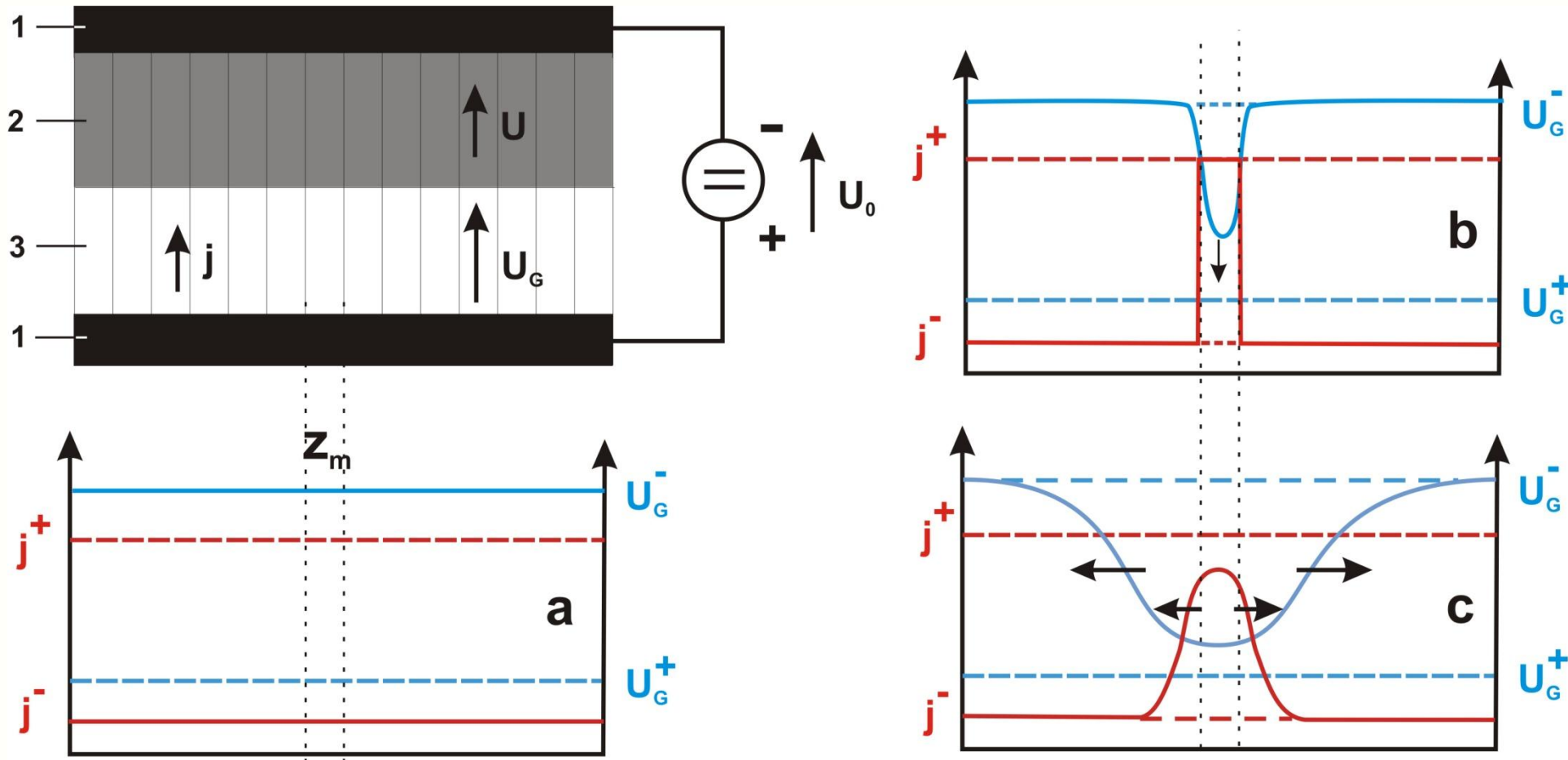


Illustration of the Formation of a Stable Current Filament in the 1-Dimensional Double Layer Continuum System I



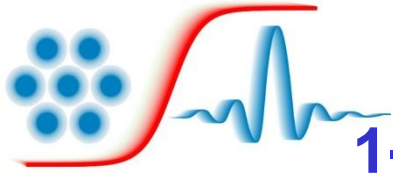


Illustration of the Formation of a Stable Current Filament in the 1-Dimensional Double Layer Continuum System II

presuppositions for the FH equation:

- $u = j$, $v = U$
- **slow relaxation of $u = j$ with respect to $v = U$: $\tau < 1$**
- **slow diffusion of $u = j$ with respect to $v = U$: $d_u < d_v$**
- **$\lambda > 1$ and appropriate κ_1 in order to assure a stable stationary low current state $(u^-, v^-) = (j^-, U^-)$ and a stable stationary high current state $(u^+, v^+) = (j^+, U^+)$ (see Fig.I a)**
- **note: $U^- = U_o - U_G^-$ and $U^+ = U_o - U_G^+$**
- **for simplicity : $\kappa_2 = 0$**

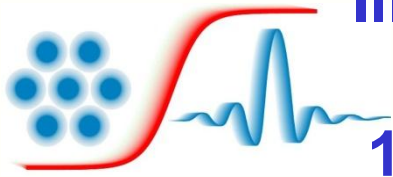


Illustration of the Mechanism for the Formation and Stabilization of a Stable Current Filament in the 1-Dimensional Double Layer Continuum System III

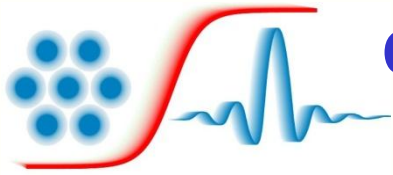
Fig. I b Initial condition. The current density $j = u$ is set to j^+ in the cell Z_m and to j^- elsewhere, the voltage drop $U = v$ is set to U^- everywhere.

Evolution of the system. As a result of the large current density j^+ at Z_m the voltage drop U increases (this can also be seen from the second FH equation), consequently the corresponding voltage drop U_G decreases; this is a relatively, fast process due to the fast relaxation of U and U_G

Fig. I c. Some later time, due to moderate diffusion of j , the initial distribution broadenes; however, somewhat faster also the voltage drop distribution broadens due to large pseudo diffusion.

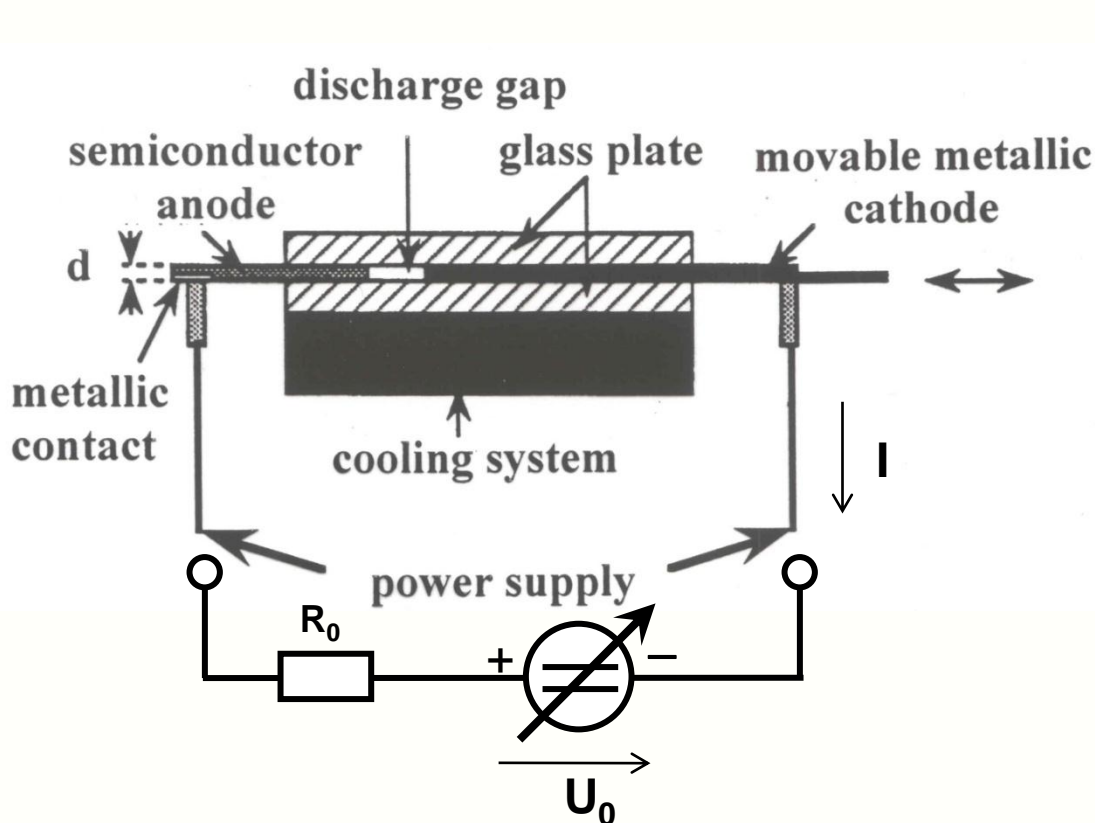
One may now look upon the dynamics as a competition of two phenomena: Two current density fronts travel in opposite direction, such, that the region of breakdown extends – however, two faster U_G voltage fronts, propagating also in opposite direction, lead to an obstruction of further extension of the region of breakdown.

Eventually the two pairs of fronts come to a rest in the final end.

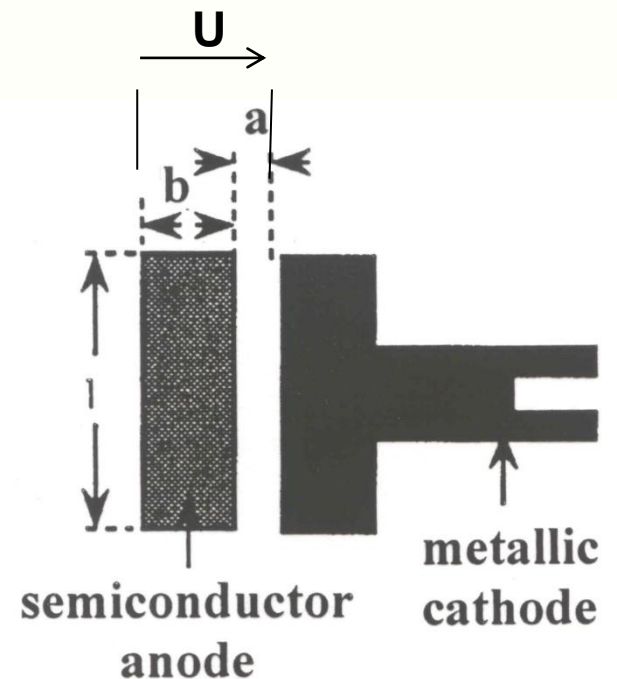


Quasi 1-Dimensional DC Gas-Discharge Systems: Experimental Set-Up

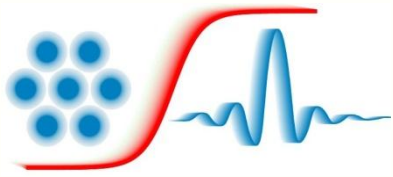
cross section



top view

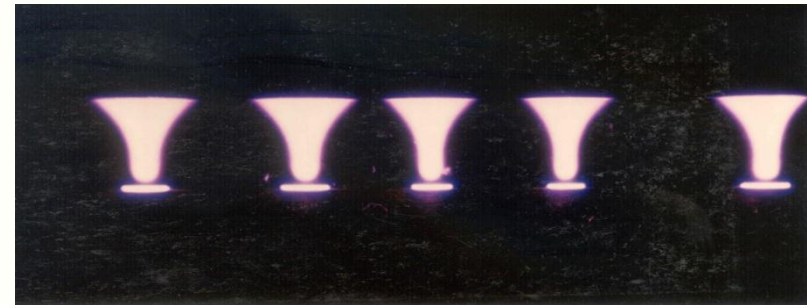
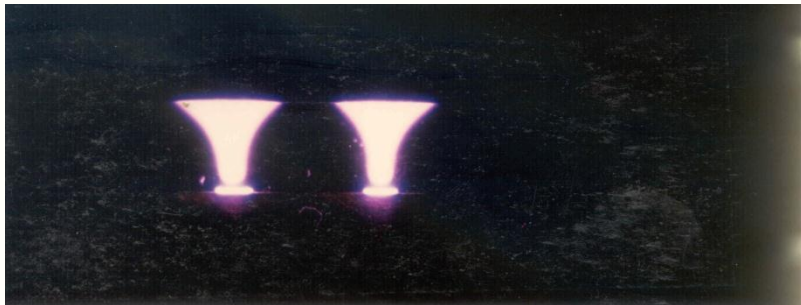
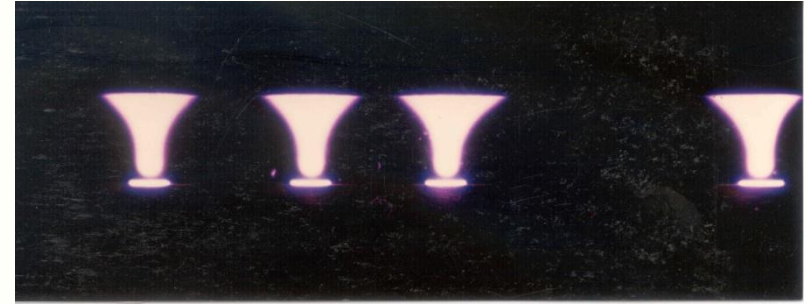
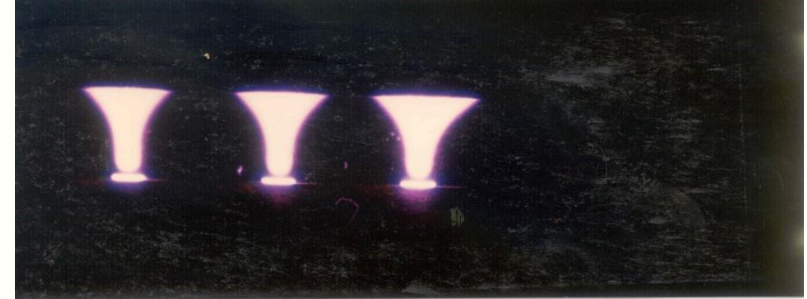


$U_0 = 0 - 1500 \text{ V}$, $R_0 = 20 - 1000 \text{ k}\Omega$, $a = 3 \text{ mm}$, $b = 10 \text{ mm}$, $d = 0.3 \text{ mm}$,
 $\rho_{\text{Si}} = 0.9 - 2.6 \text{ M}\Omega \text{ cm}$, gas: Ar, He + air, $p = 40 - 100 \text{ hPa}$, $l = 20 - 50 \text{ mm}$



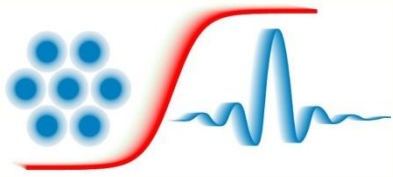
Quasi-1-Dimensional DC Gas-Discharge System: Experimental Cascade with Increasing Number of Stationary Current Filaments (A)

luminescence radiation distribution in the discharge slid for increasing driving voltage U_0



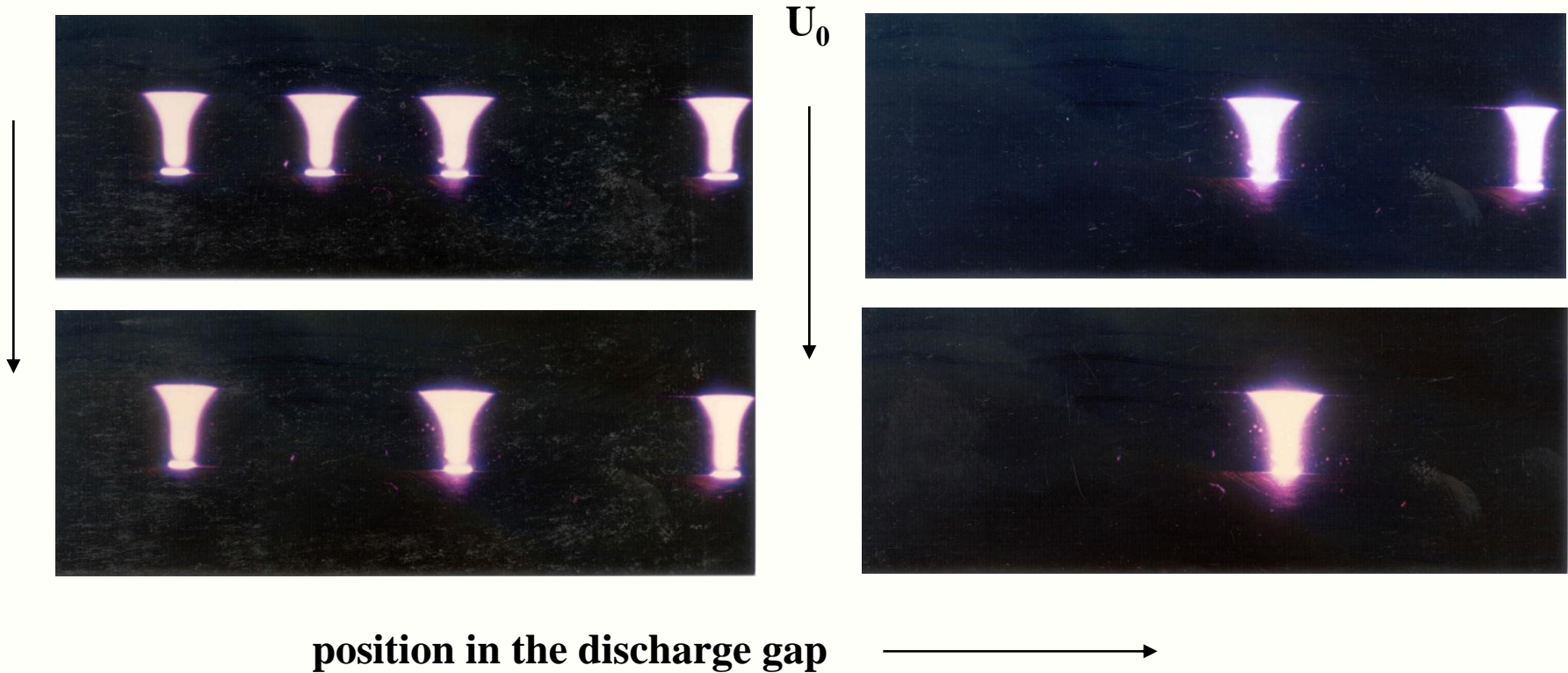
position in the discharge gap



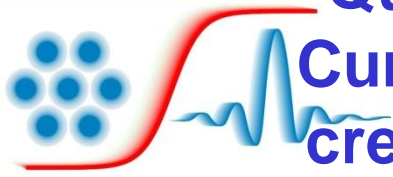


Quasi-1-Dimensional DC Gas-Discharge System: Experimental Cascade with Decreasing Number of Stationary Current Filaments (B)

luminescence radiation distribution in the discharge space for decreasing driving voltage

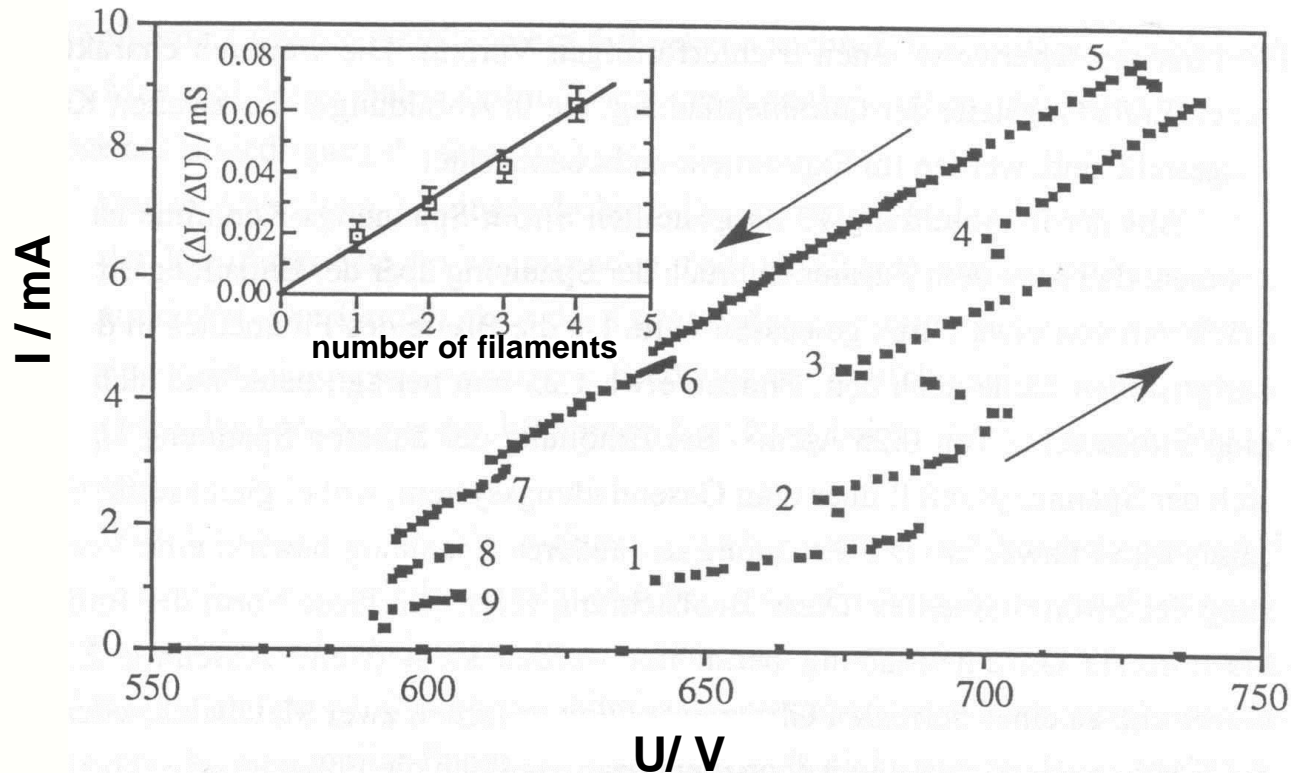


cathode: metal, anode: Si, $\rho=1.5\text{k}\Omega\text{ cm}$, gas: Ar, $p=170\text{hPa}$, $d=3.75\text{ mm}$, $R_0=164\text{k}\Omega$,
 $U_0=640\rightarrow 700\rightarrow 590\text{V}$



Quasi-1-Dimensional DC Gas-Discharge System IV: Current-Voltage Characteristic for Increasing and De- creasing Number of Stationary Current Filaments (C)

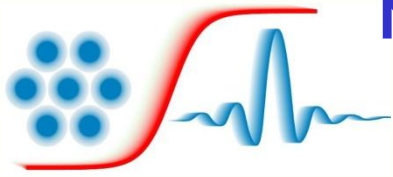
(total current)-(voltage drop at the device) characteristic for increasing and decreasing driving voltage U_0



Semiconductor: Si
 $d_{\text{SC}}=0.3\text{mm}$

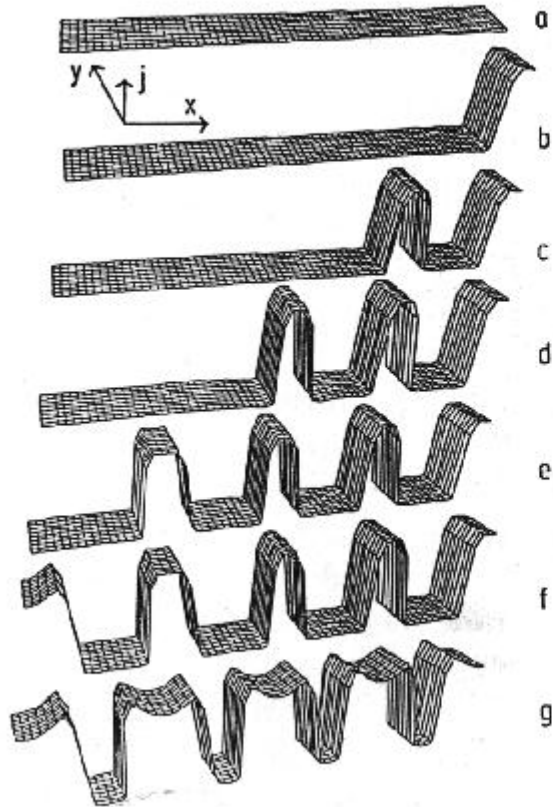
$\rho=1\text{k}\Omega\text{ cm}$
 $h_{\text{SC}}=10\text{mm}$

$l_{\text{gap}}=38\text{mm}$
gas: 170hPa Ar + 20hPa air
 $d_{\text{gap}}=3.75\text{mm}$



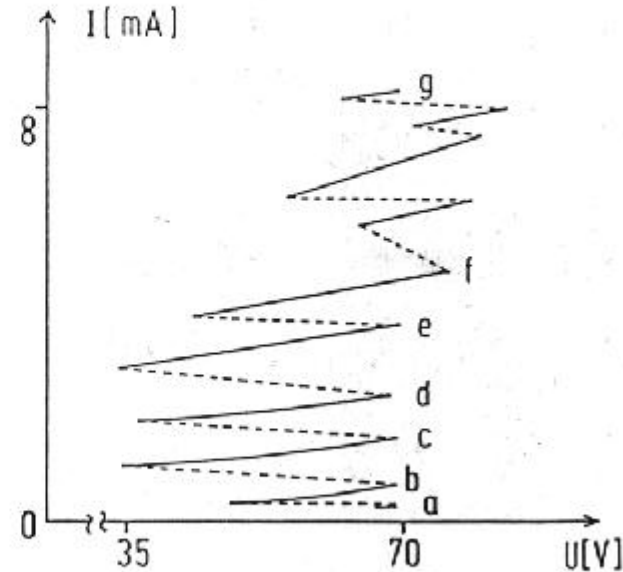
Numerical Solution of the Network Equation in \mathcal{R}^1 : Increasing Number of Stationary DSs when Increasing U_0

A)

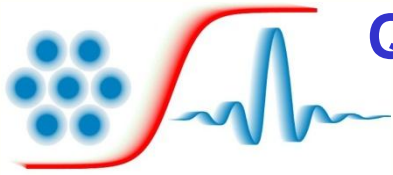


stationary patterns of the current density for U_0 increasing from a to g

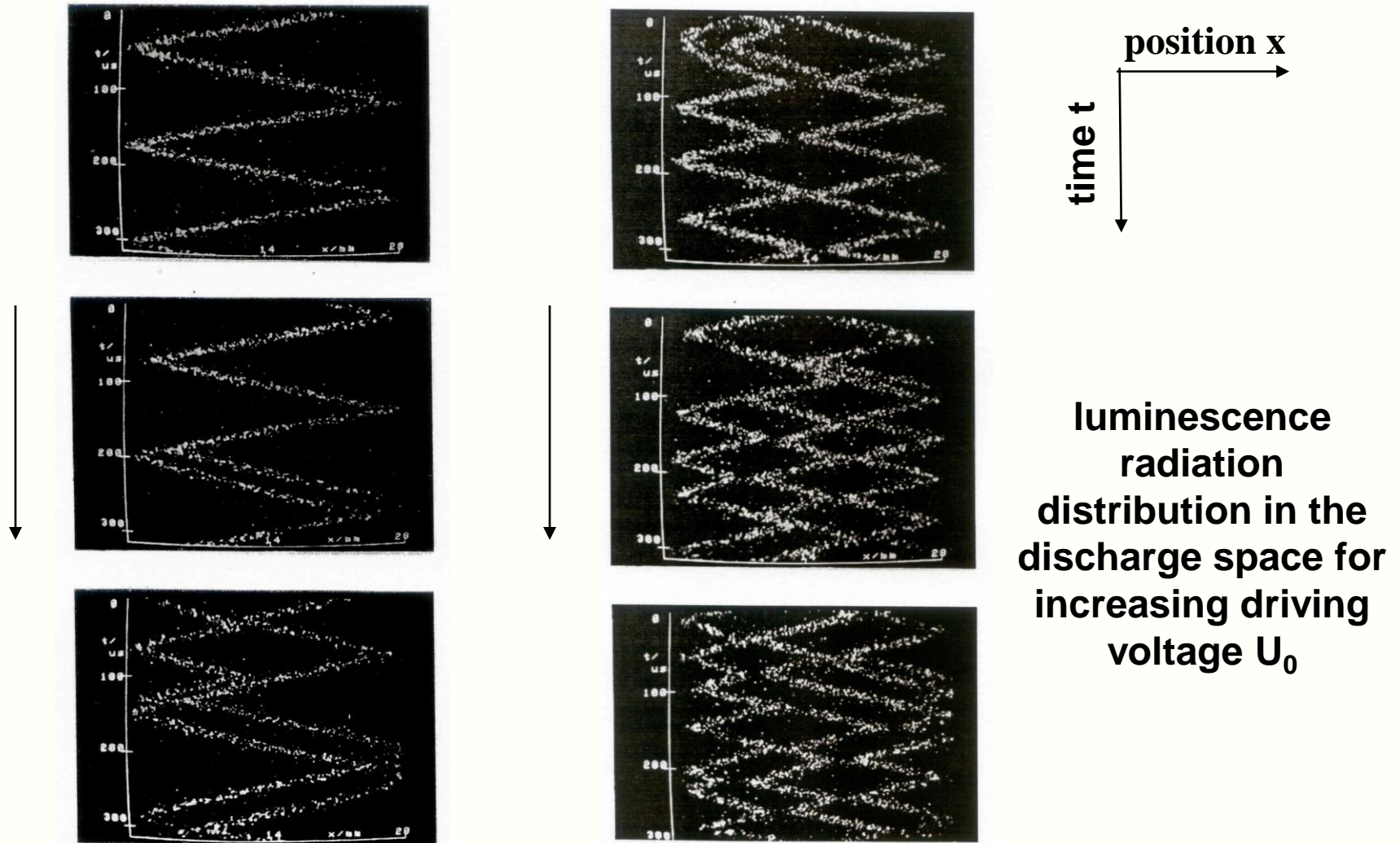
B)

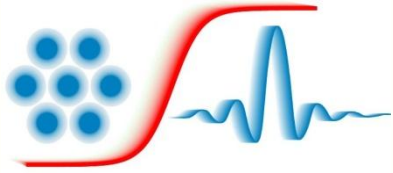


(total current)-(voltage drop at the device) characteristic for increasing U_0

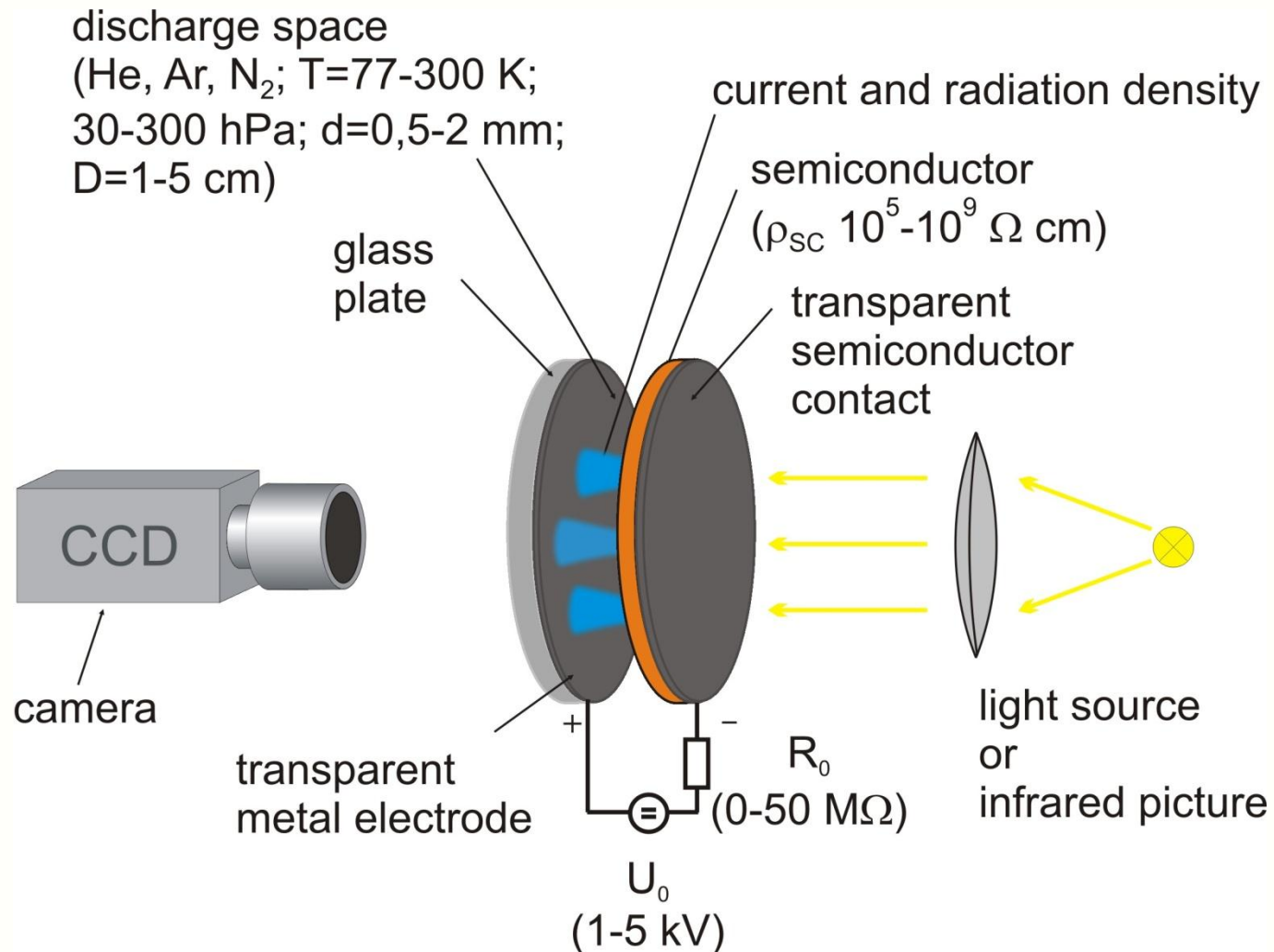


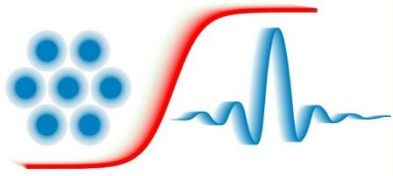
Quasi-1-Dimensional DC Gas-Discharge System: Cascade of Traveling Current Filaments



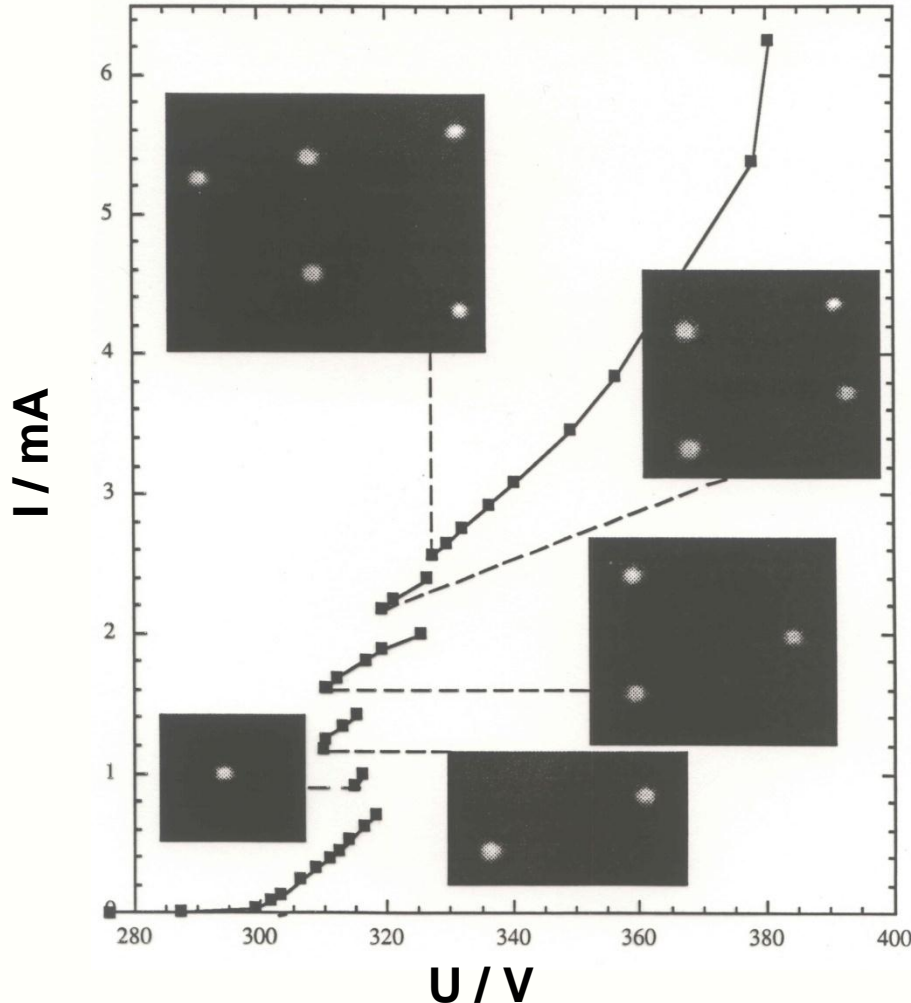


Quasi 2-Dimensional DC Gas-Discharge System: Experimental Set-Up





Quasi 2-Dimensional DC Gas-Discharge System: Cascade with Increasing Number of Stationary Current Filaments when Increasing U_0



the driving voltage U_0 is increased

I = total current

U = voltage drop at the device

anode: Si compensated with Au

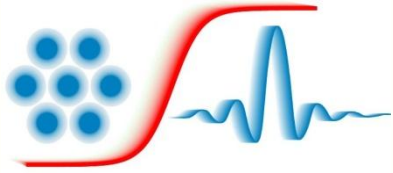
$\rho_{SC}=150 \text{ k}\Omega \text{ cm}$, $a_{SC}=0.5\text{mm}$, $T_{SC}=4^\circ\text{C}$,

cathode: optically transparent ITO,

gas: He, $p=300\text{hPa}$, $d=1.5\text{mm}$,

$l_x=l_y=11\text{mm}$, $R_0=86\text{k}\Omega$, $t_{exp}=20\text{ms}$

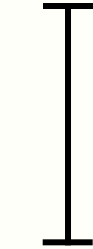
inserted images:
luminescence radiation distribution in
the discharge space



Quasi 2-Dimensional DC Gas-Discharge System: Travelling Isolated Current Filament (Movie)

luminescence
radiation
distribution in
the discharge
plane

1 cm

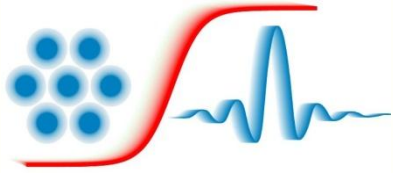


real
time



if not linked:
start movie
“Filament.avi”
in the folder

parameters: $U_0=2,7$ kV, $\rho_{SC}=4,95$ M Ω cm,
 $R_0=20$ M Ω , Gas: N₂, T=100 K, p=280 hPa,
D=30 mm, d=250 μ m, $a_{SC}=1$ mm, I=46 μ A



The Generalized FN Equation I: The Equation

$$u_t = d_u^2 \Delta u + f(u) - \kappa_3 v - \kappa_4 w + \kappa_1,$$

$$\tau v_t = d_v^2 \Delta v + u - v,$$

$$\theta w_t = d_w^2 \Delta w + u - w,$$

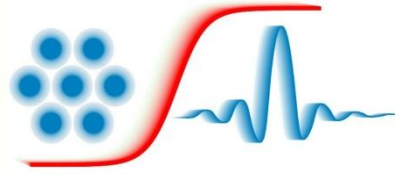
$$u = u(x, y, z, t),$$

$$v = v(x, y, z, t),$$

$$w = w(x, y, z, t),$$

$$f(u) = \lambda u - u^3$$

$$d_u, d_v, d_w, \tau, \theta, \kappa_3, \kappa_4 \geq 0, \lambda > 1$$



The Generalized FN Equation II: Physical Meaning of Dependent Variables and Parameters in the Case of Gas-Discharge

u *current density*

v *voltage drop at the high ohmic layer*

w *surface charge? temperature?*

τ *dielectric relaxation time normalized to reaction time of charge carriers*

θ *relaxation time normalized to reaction time of charge carriers*

d_u *ambipolar diffusion length*

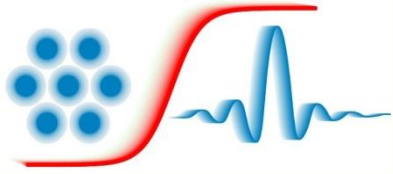
d_v *electrical pseudo diffusion length*

λ *parameter describing the current density vs. voltage characteristic*

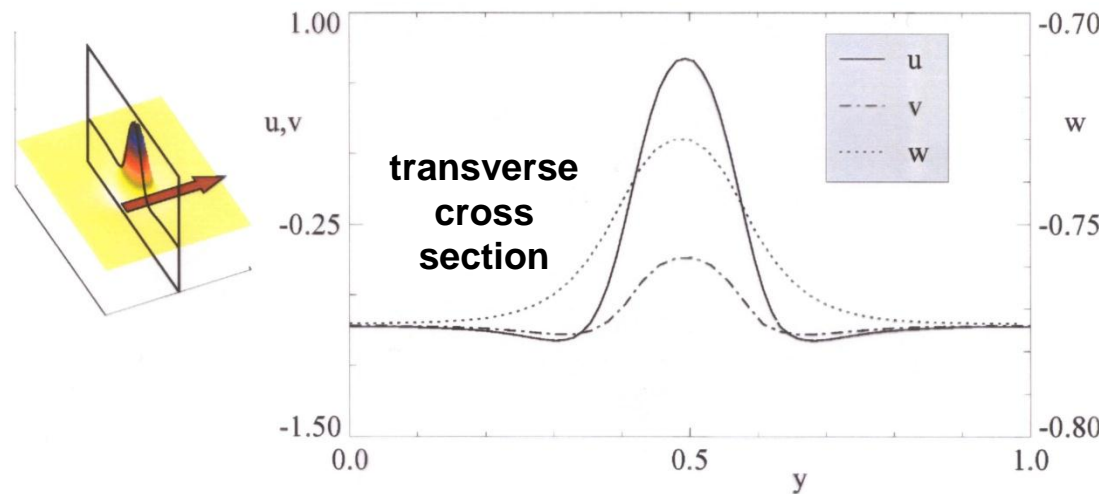
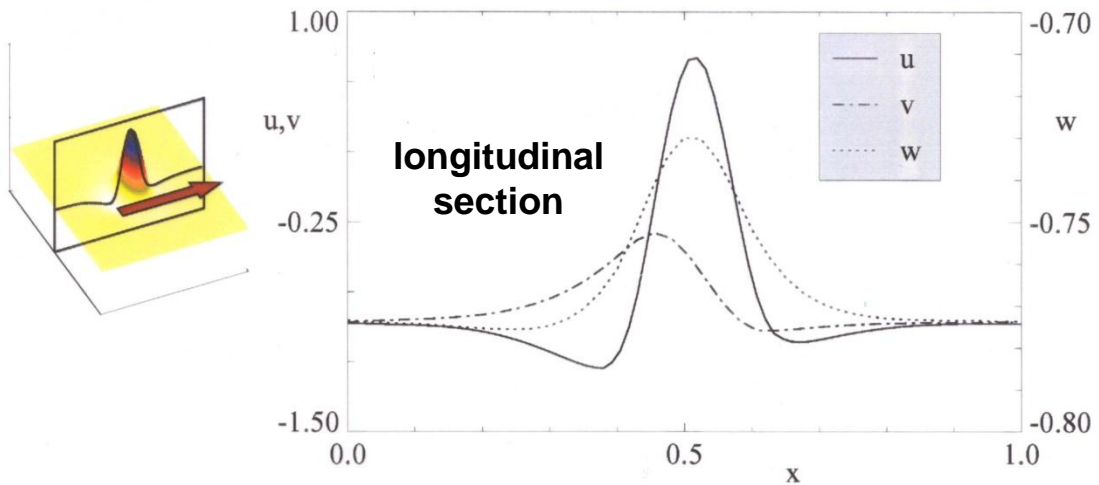
d_w *diffusion length u*

κ_1 *driving voltage*

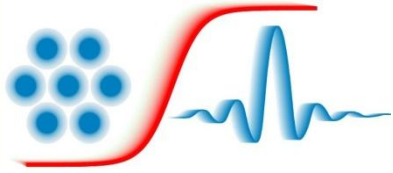
κ_3, κ_4 *physical origin not well known*



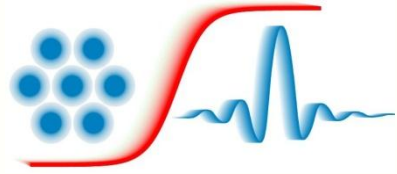
Numerical Solution for the Generalized FH Equation in \mathbb{R}^2 : Isolated Travelling DS



parameters: $D_u=10^{-3}$,
 $D_v=1.25 \cdot 10^{-3}$, $D_w=6.4 \cdot 10^{-2}$,
 $\kappa_1=-6.92$, $\kappa_3=1$, $\kappa_4=8.5$, $\lambda=2$,
 $\tau=25$, $\theta=1$, $\Delta x=1/38$

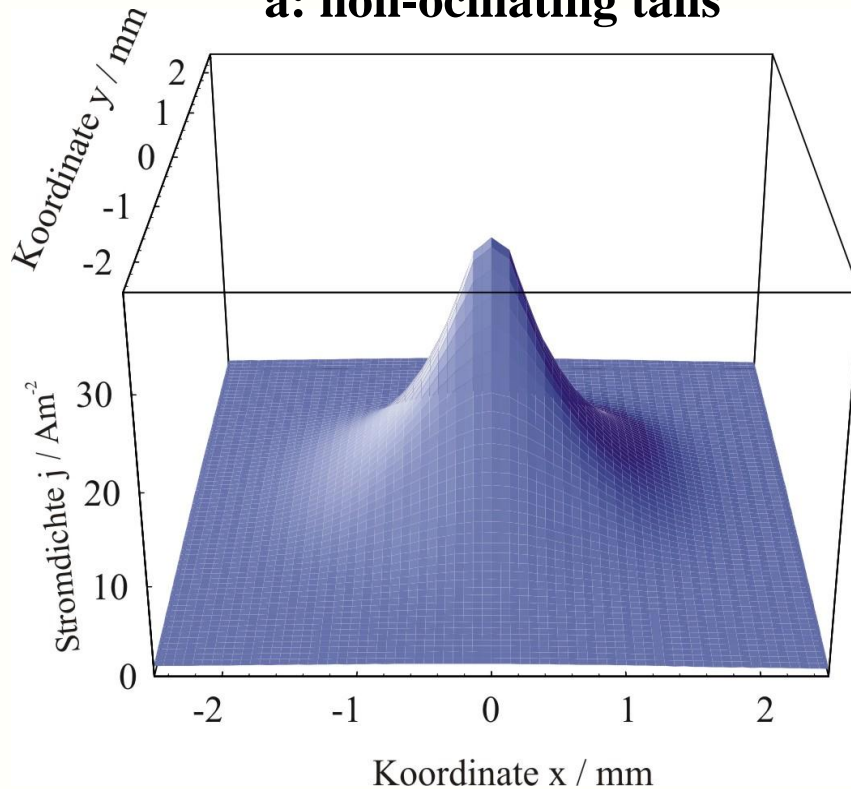


4. Additional Results on Dissipative Solitons (DSs) in DC Gas-Discharge Systems and Comparison with Solutions of the Generalized FitzHugh-Nagumo (FN) Equation



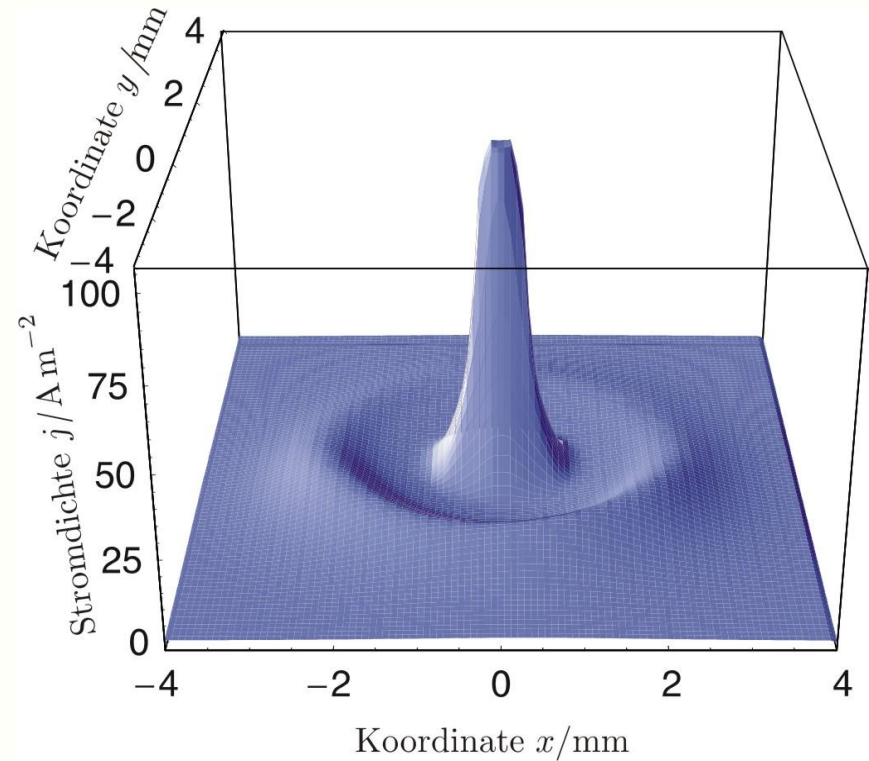
Experimental Quasi 2-Dimensional DC Gas-Discharge System: Current Filaments with Non-Oscillating and Oscillating Tails

a: non-oscillating tails

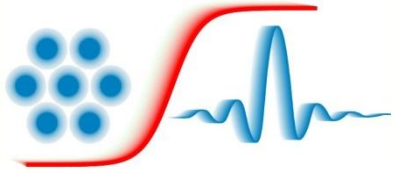


parameters: $U_0 = 2,74$ kV, $\rho_{SC} = 4,95$ M Ω cm, $R_0 = 20$ M Ω , Gas: N₂, T=100 K, p = 280 hPa, D=30 mm, d = 250 μ m, $a_{SC}=1$ mm, I = 46 μ A, $t_{exp}=20$ ms

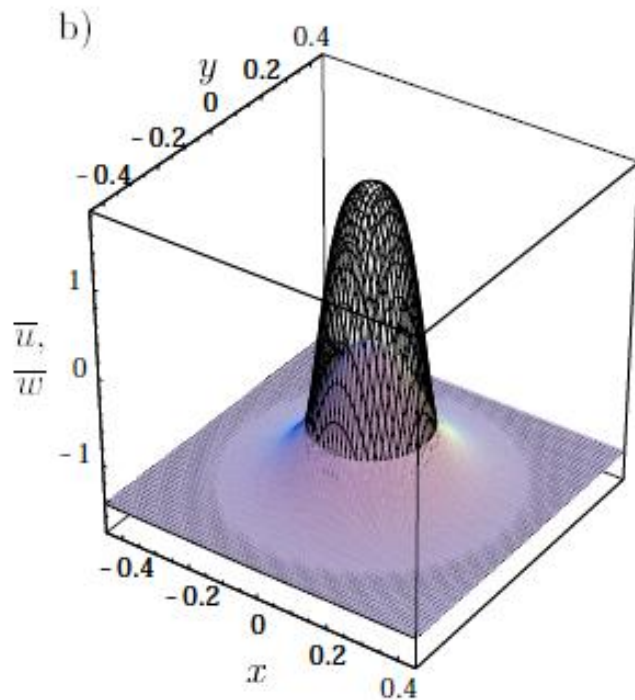
b: oscillating tails



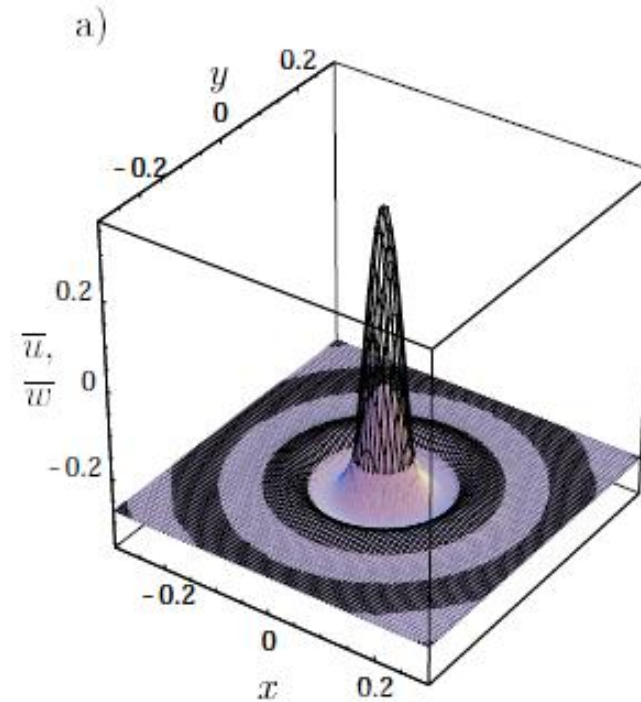
parameters: $U_0 = 3,6$ kV, $\rho_{SC} = 3,05$ M Ω cm, $R_0 = 4,4$ M Ω , Gas: N₂, T=100 K, p = 279 hPa, D=30 mm, d = 500 μ m, $a_{SC}=1$ mm, I = 200 μ A, $t_{exp}=20$ ms



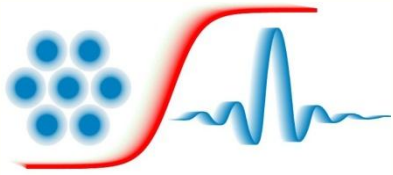
Numerical Solution for the Generalized FH Equation in \mathbb{R}^2 : Stationary DSs with Non-Oscillating and Oscillating Tails I



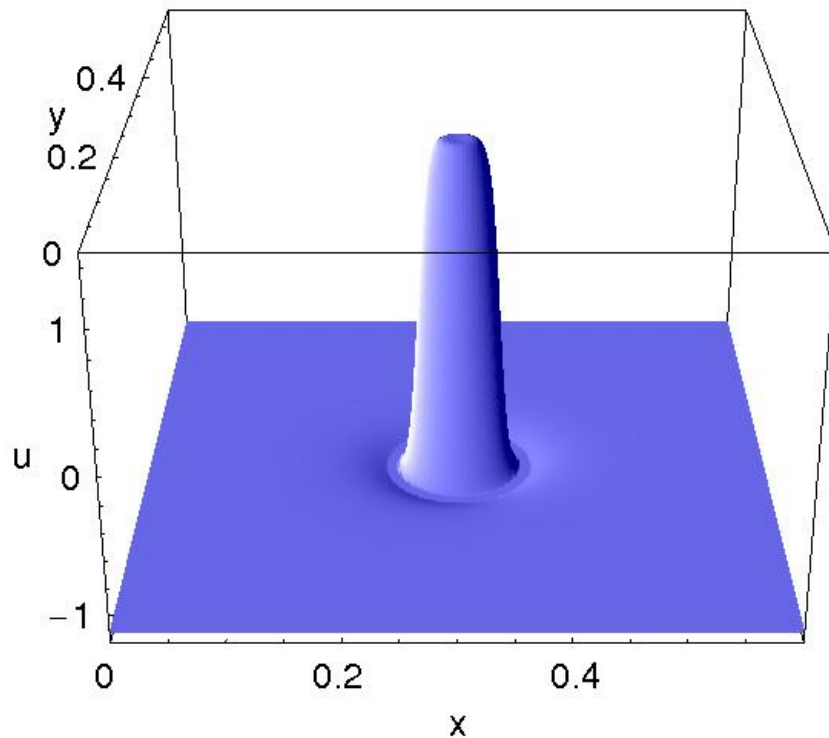
a: non-ocillating tails



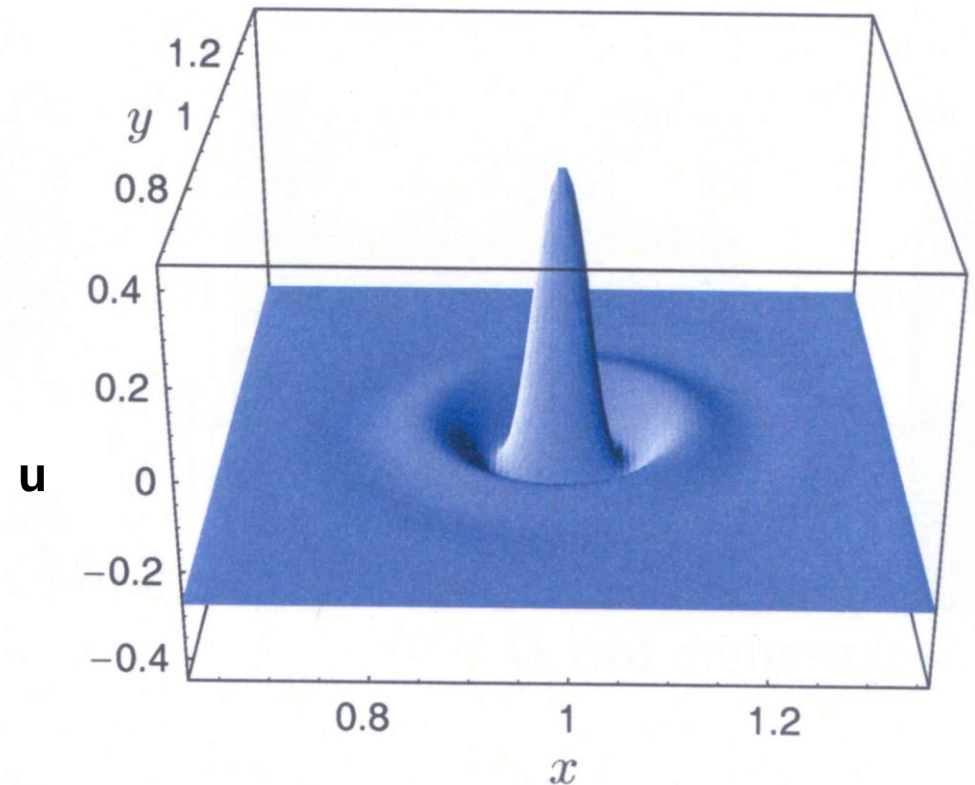
b: oscillating tails



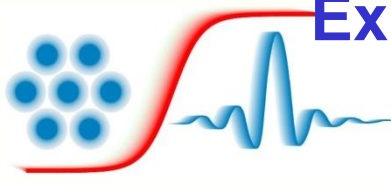
Numerical Solution for the Generalized FH Equation in \mathcal{R}^2 : Stationary DSs with Non-Oscillating and Oscillating Tails I



a: non-ocillating tails

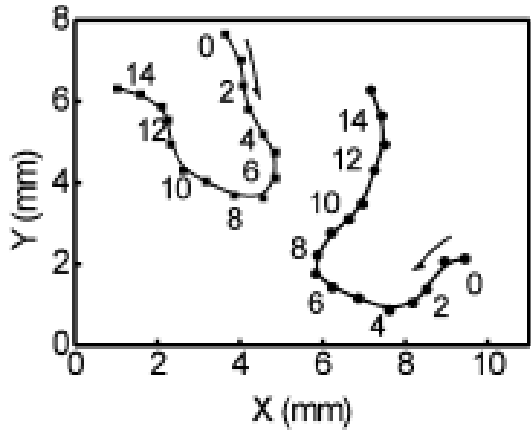


b: oscillating tails

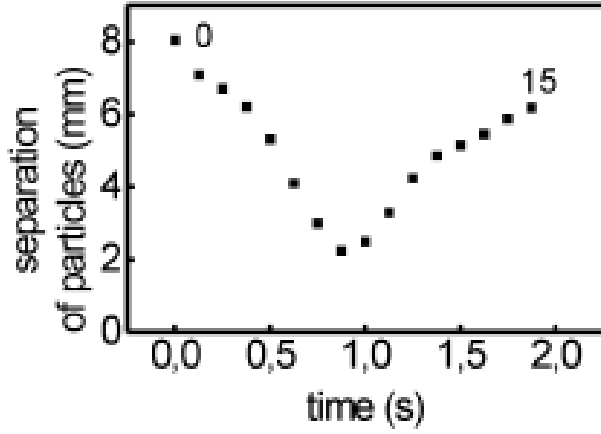


Experimental Quasi 2-Dimensional DC Gas-Discharge System: Examples for Travelling, Scattering and Cluster Formation of Filaments (“Molecules”)

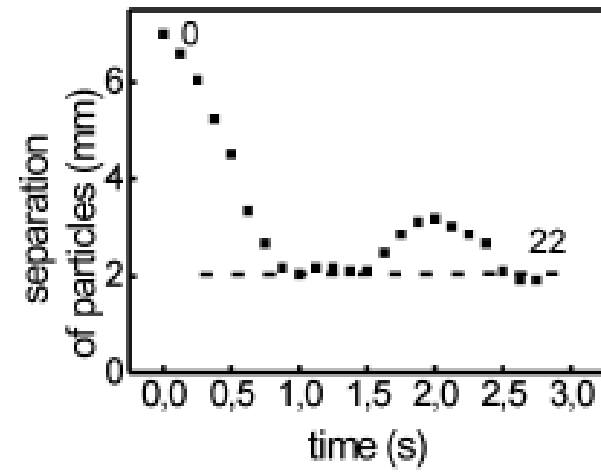
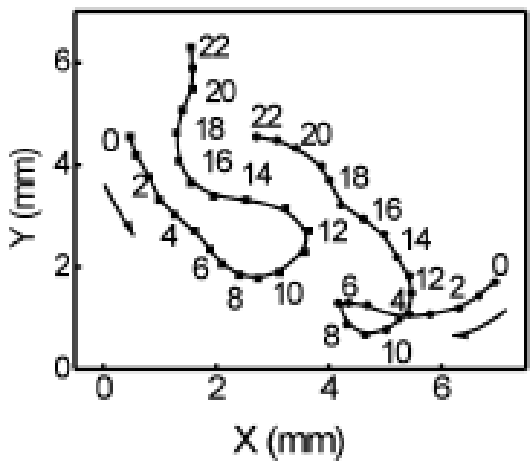
trajectories



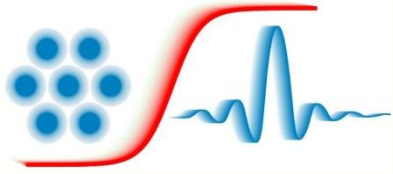
distance



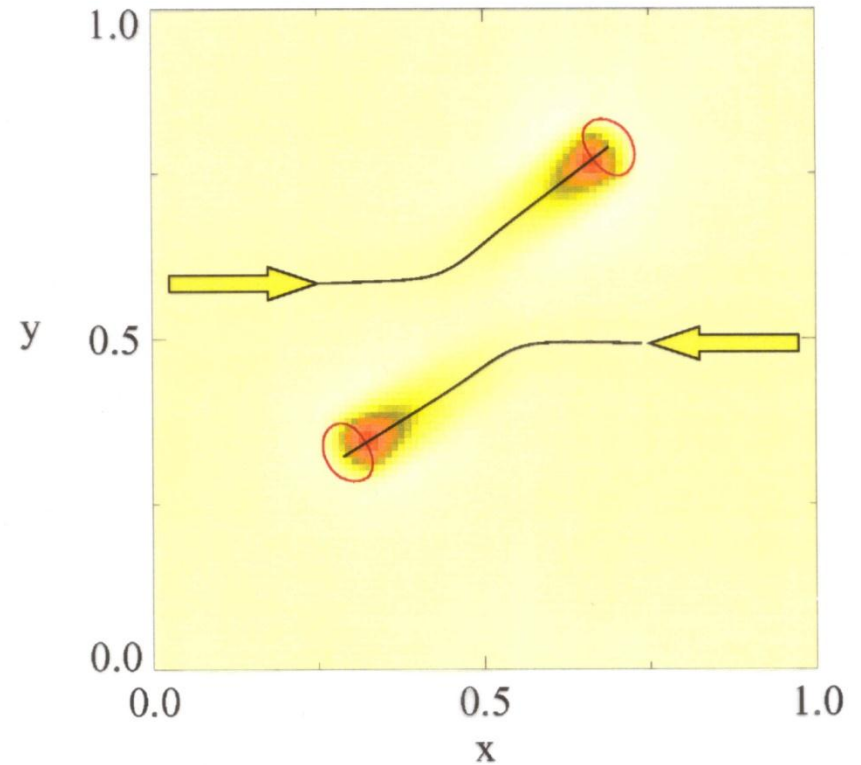
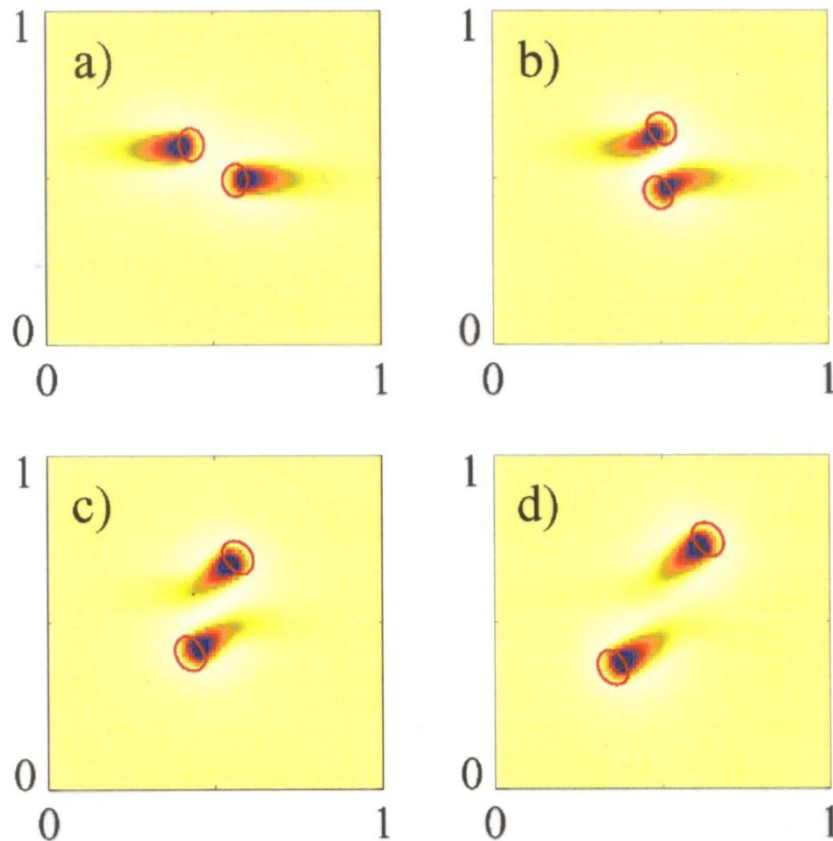
scattering



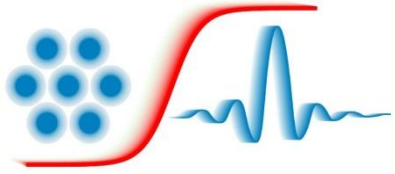
cluster formation



Numerical Solution for the Generalized FH Equation in \mathcal{R}^2 : Scattering of 2 DSs (Weak Interaction I)

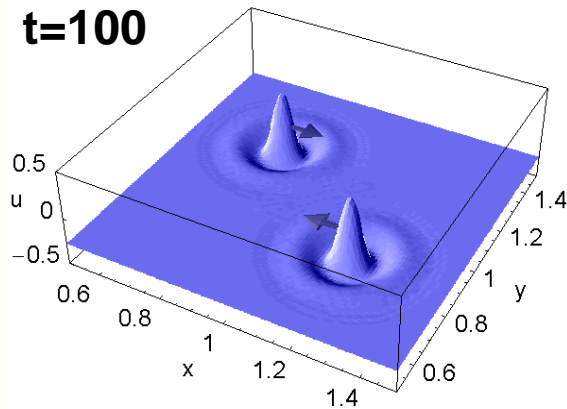


parameters: $D_u=1.55 \cdot 10^{-4}$, $D_v=1.95 \cdot 10^{-4}$, $D_w=0.05$, $\kappa_1=-8.715$, $\kappa_3=1$, $\kappa_4=8.44$, $\lambda=2$,
 $\tau=48$, $\theta=0.5$, $\Delta x=1/100$

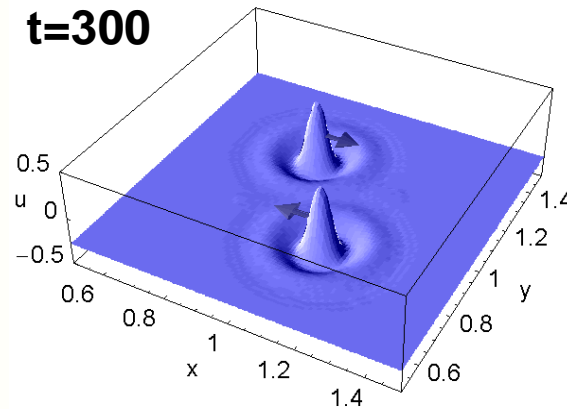


Numerical Solution for the Generalized FH Equation in \mathcal{R}^2 : Scattering of 2 DSs (Weak Interaction II)

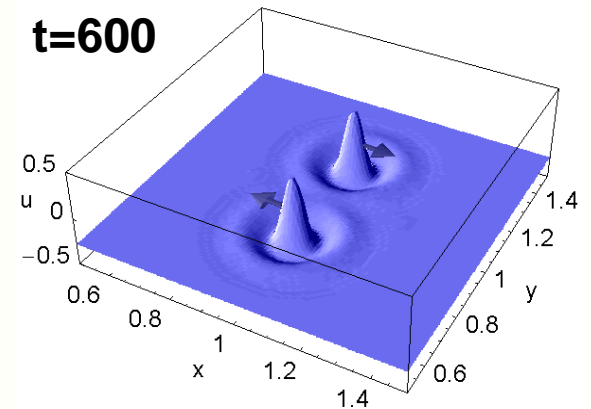
t=100



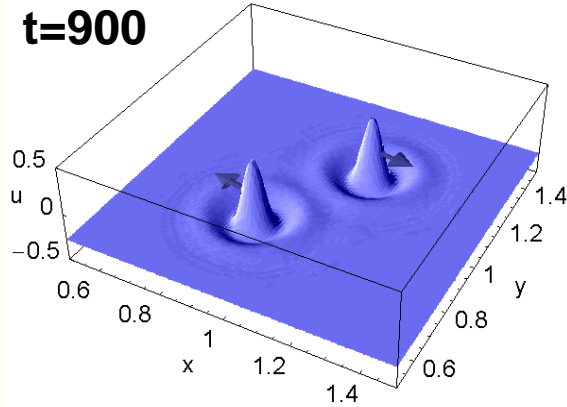
t=300



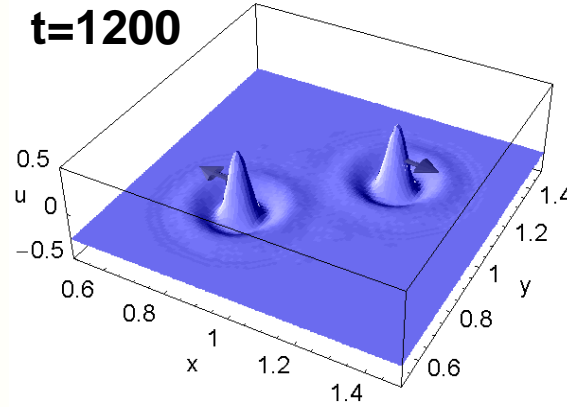
t=600



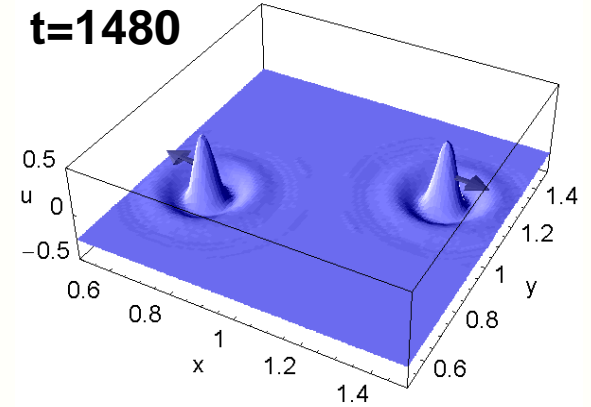
t=900



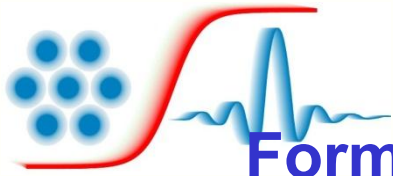
t=1200



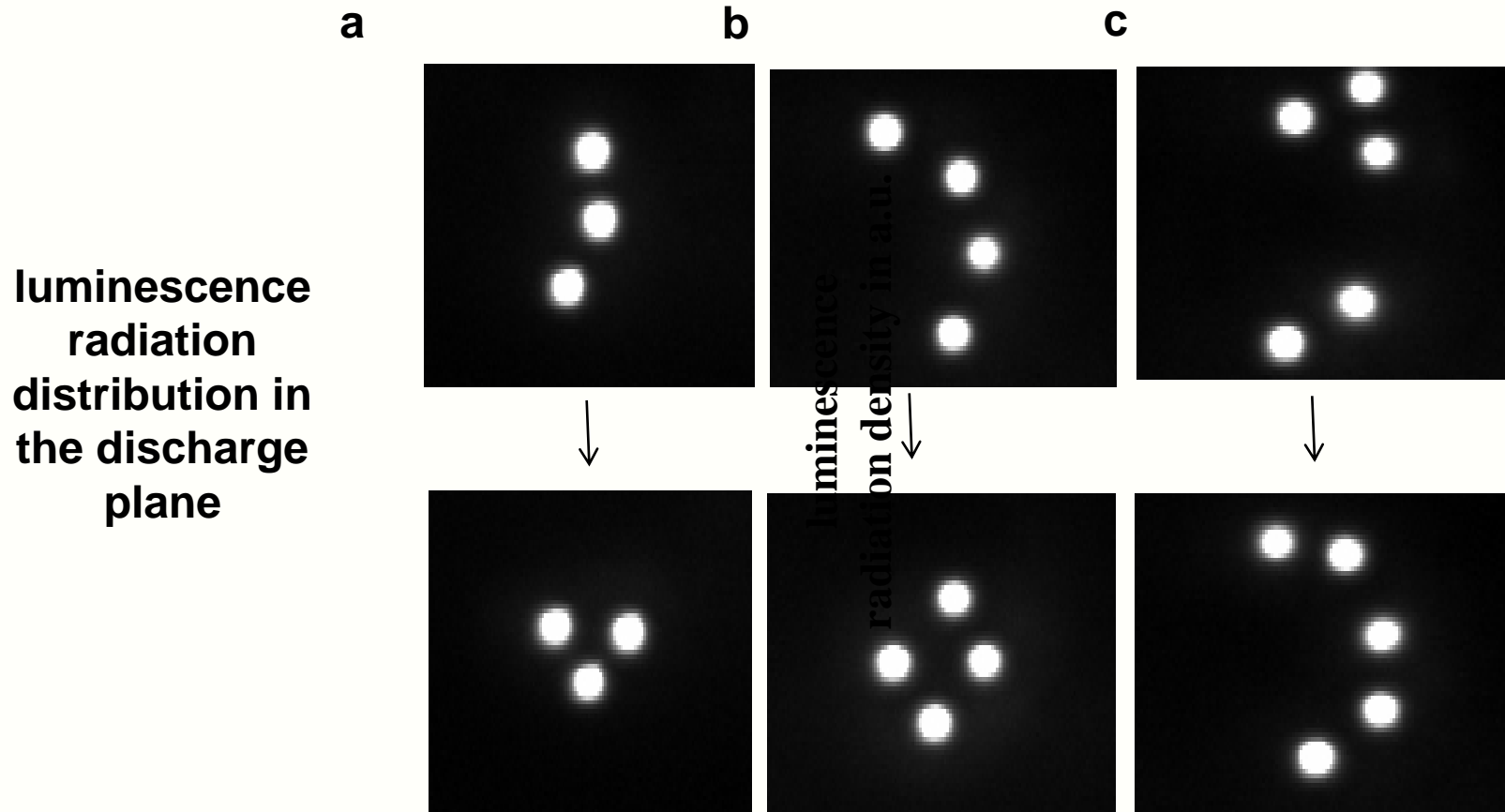
t=1480



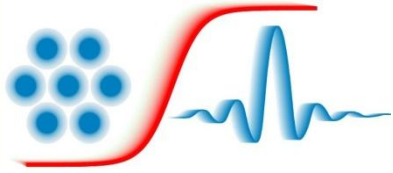
$\tau=3.35, \theta=0, D_u=1.1 \cdot 10^{-4}, D_v=0, D_w=9.64 \cdot 10^{-4}, \lambda=1.01, \kappa_1=-0.1, \kappa_3=0.3, \kappa_4=1.0$
 $\Omega=[0,1] \times [0,1], \Delta x=5 \cdot 10^{-3}, \Delta t=0.1.$



Experimental Quasi 2-Dimensional DC Gas-Discharge System: Formation of Clusters of Current Filaments (“Molecules”)

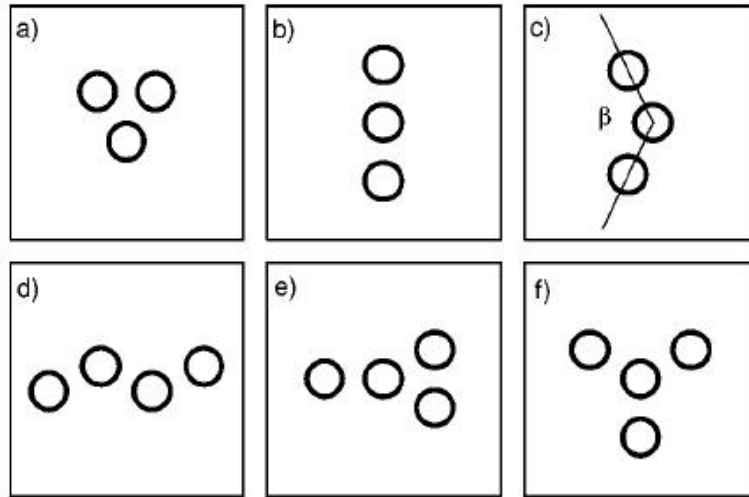


for supply voltage increasing from a to c clusters are observed with increasing number of filaments; for a given set of parameters the number of filaments is retained the but their configuration may change in time

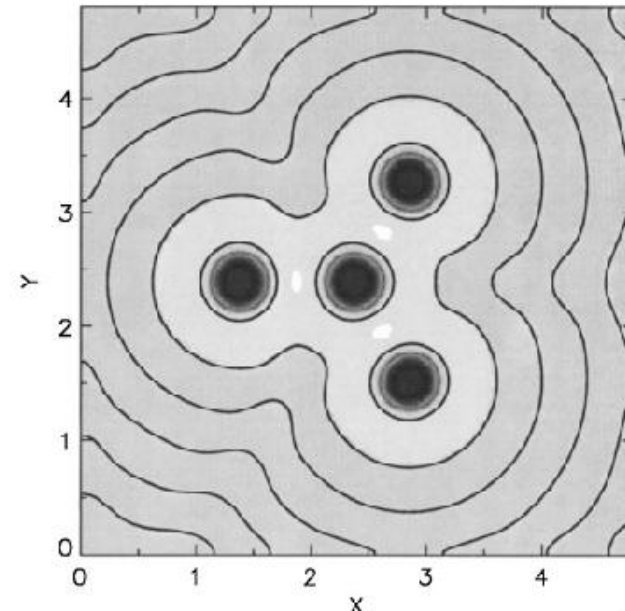


Numerical Solution for the Generalized FH Equation in \mathcal{R}^2 : Formation of Clusters of DSs (“Molecules”) I

A)

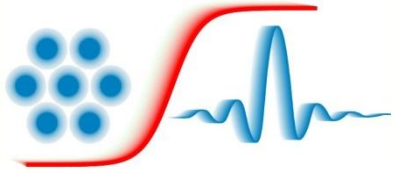


B)



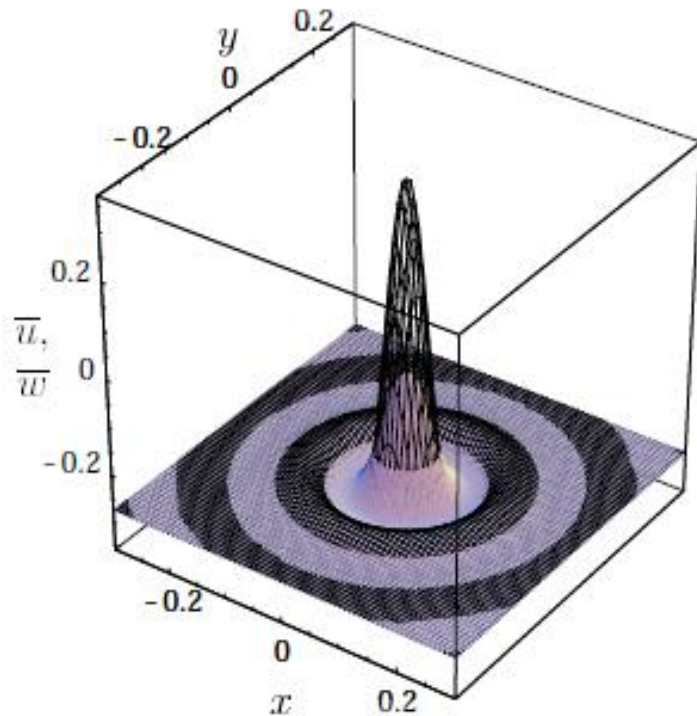
A: possible configurations for a fixed set of parameters; except for (b) the solutions are stable

B: u distribution for the case (f) in (A) reflecting the oscillatory behavior in the surrounding of individual DSs and related lock-in distances

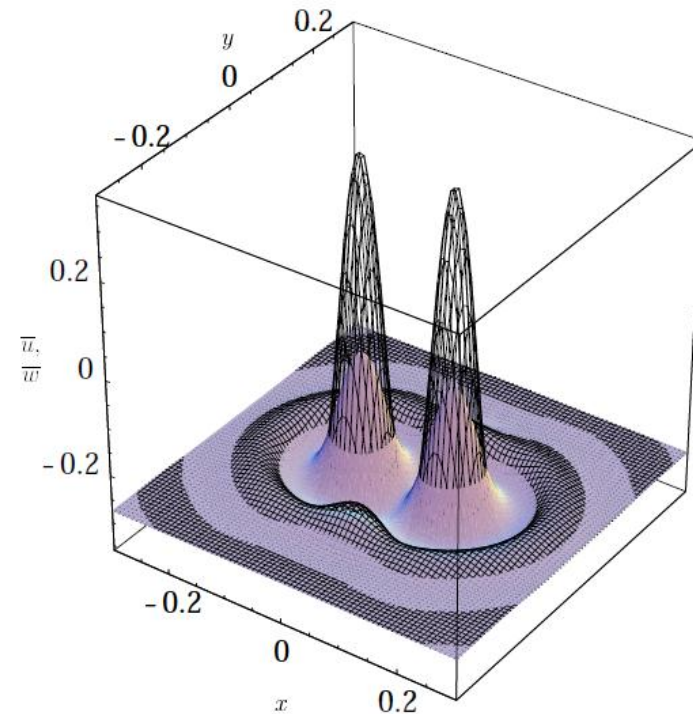


Numerical Solution for the Generalized FH Equation in \mathcal{R}^2 : Formation of Clusters of DSs (“Molecules”) II

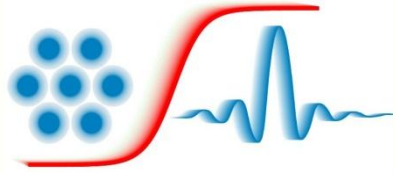
a: single DS



b: 2-DS-molecule

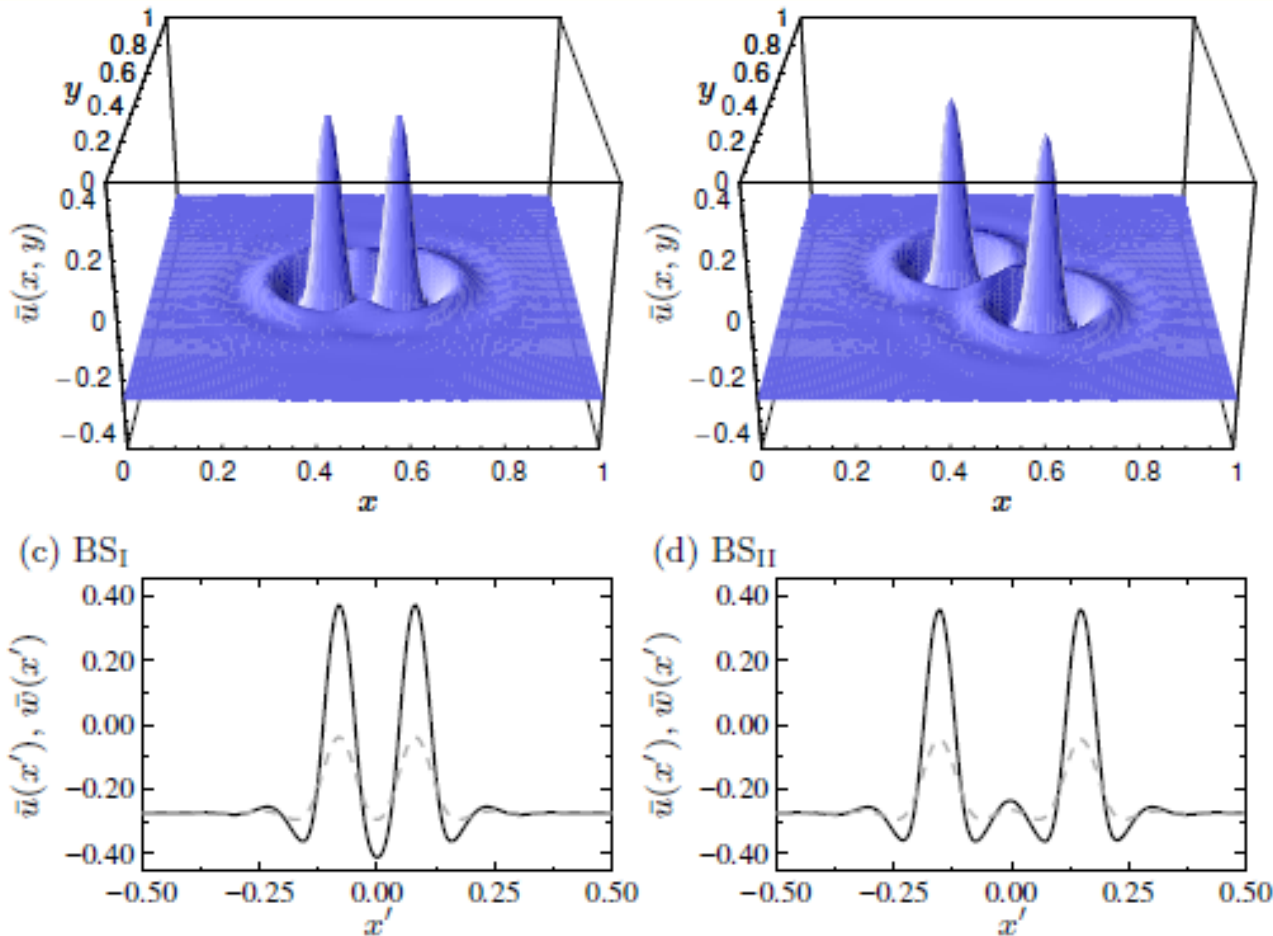


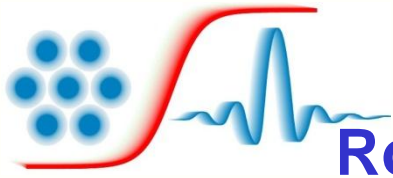
same parameters



Numerical Solution for the Generalized FH Equation in \mathcal{R}^2 : Formation of Clusters of DSs (“Molecules”) III

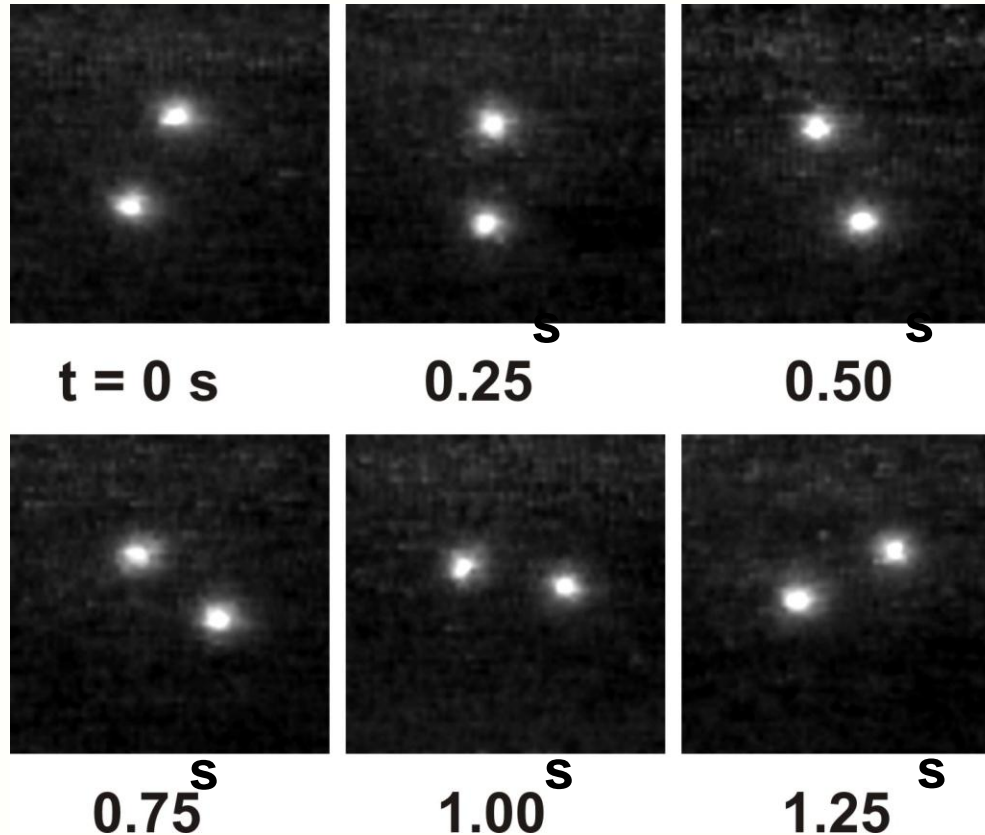
stable stationary clusters with different distance at a given set of parameters



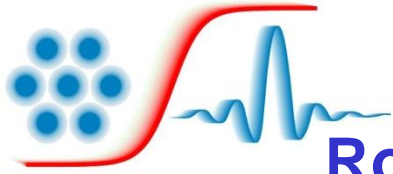


Experimental Quasi 2-Dimensional DC Gas-Discharge System: Rotating Cluster of 2 Current Filaments („Molecule“)

luminescence radiation distribution in the discharge plane



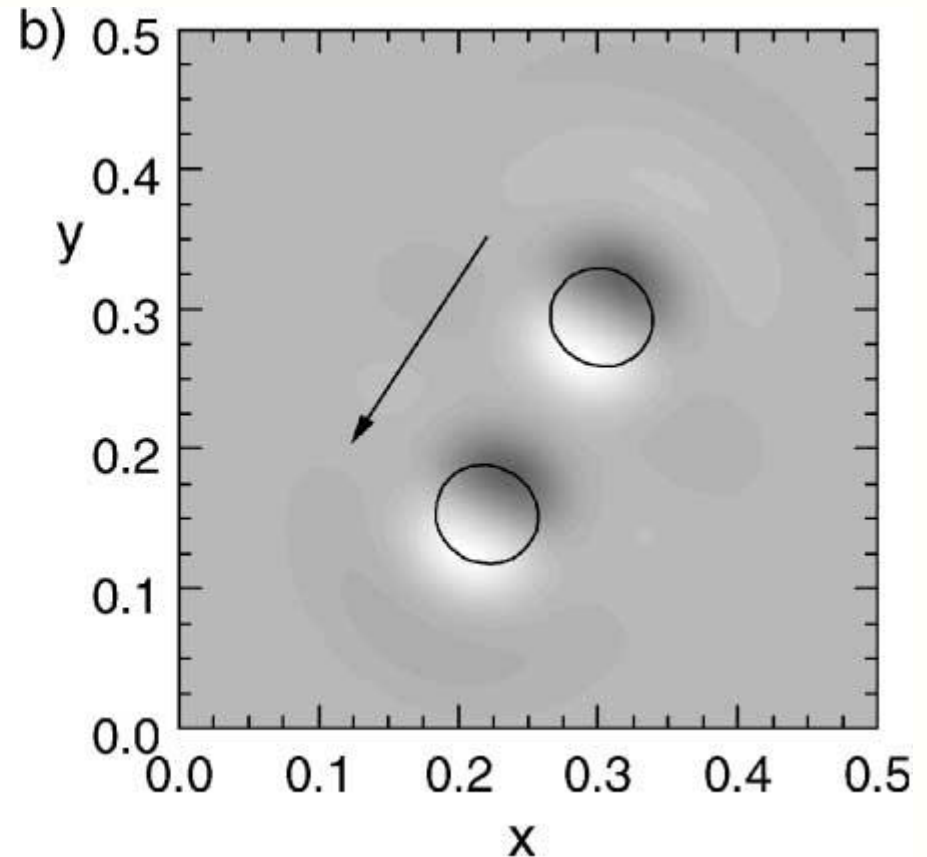
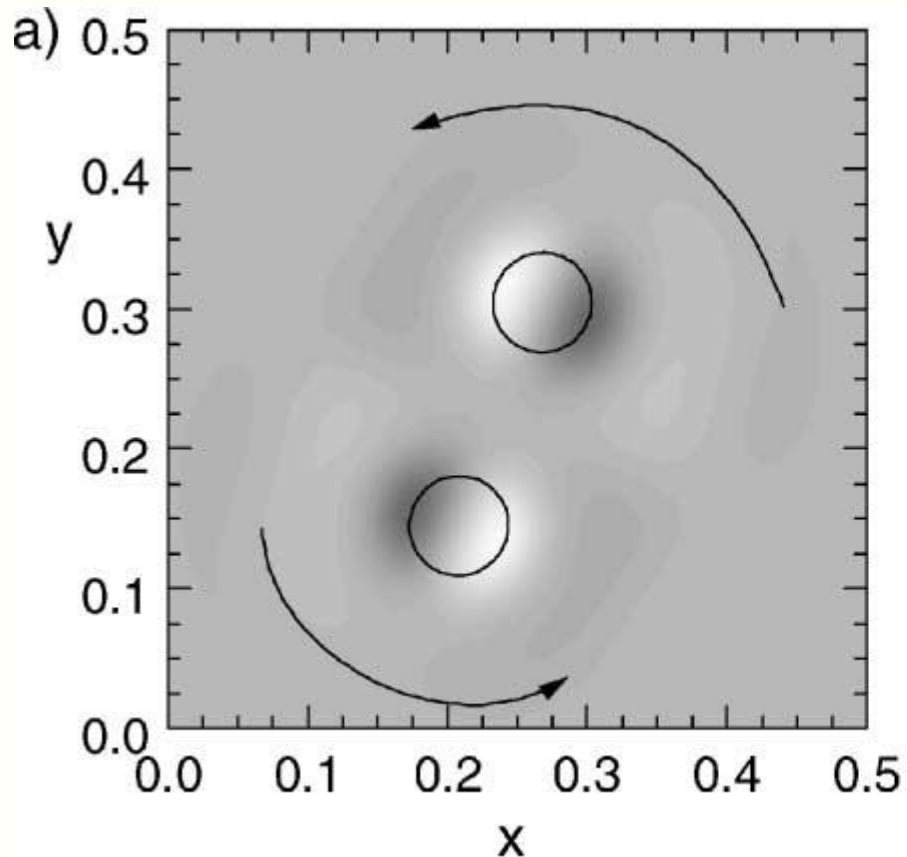
parameters: $U_0 = 1.9$ kV, $\rho_{SC} = 22.2$ M Ω cm, gas: N₂, T=90 K, $p = 2$ hPa, $d_{gap} = 0.8$ mm,
 $d_{SC} = 1$ mm, $a_{SC} = 20$ mm, $t_{exp} = 20$ ms

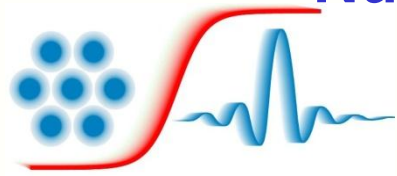


Numerical Solution for the Generalized FH Equation in \mathcal{R}^2 :

Rotating and Travelling Cluster of DSs (“Molecules”)

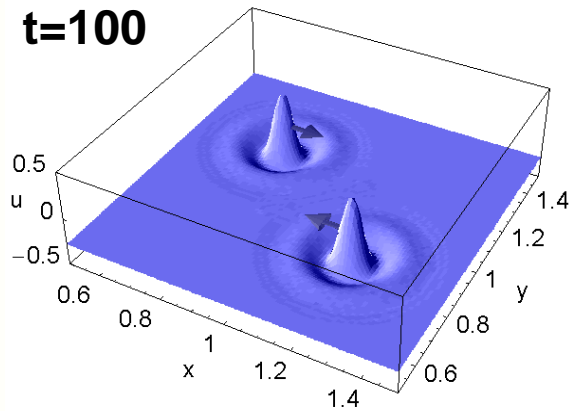
depending on the details of the collision process clusters with different dynamics may occur for a given set of parameters



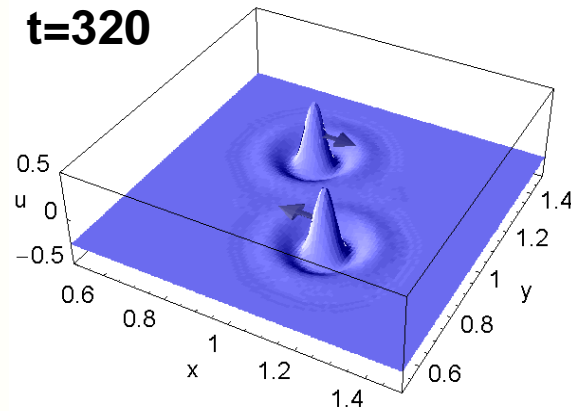


Numerical Solution for the Generalized FH Equation in \mathcal{R}^2 : Formation of a Rotating Cluster of 2 DSs (Molecule) in the Course of Collision

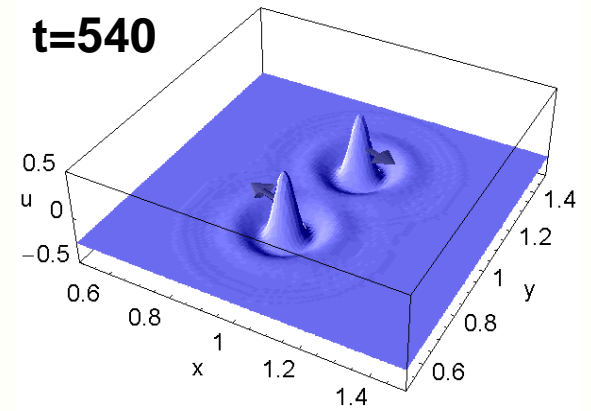
t=100



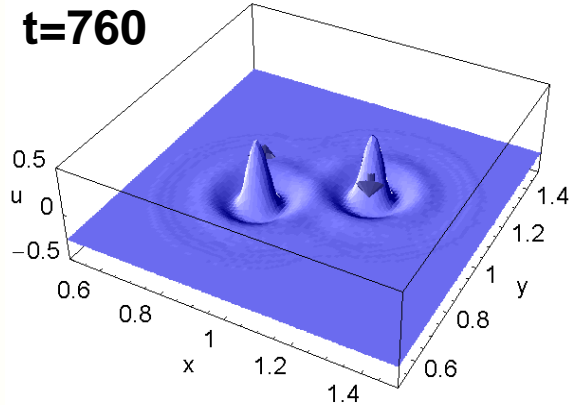
t=320



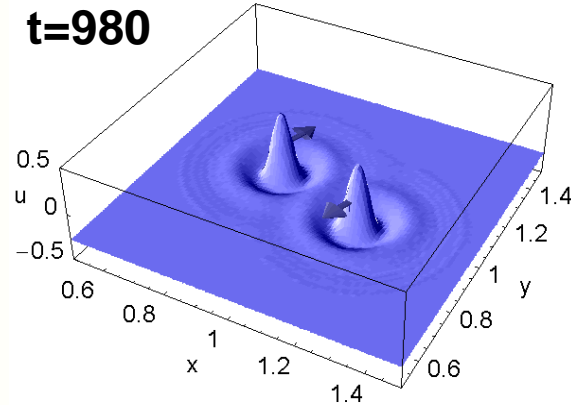
t=540



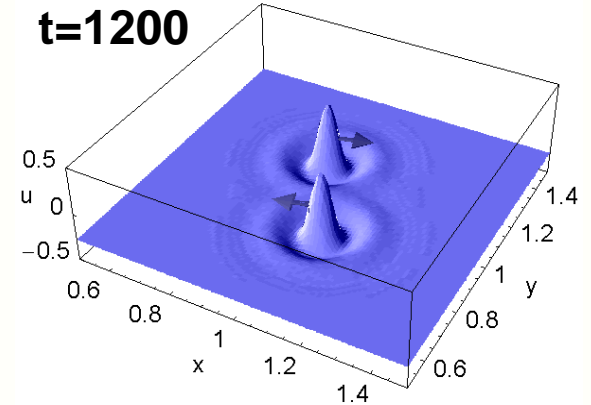
t=760



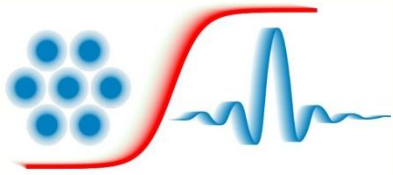
t=980



t=1200

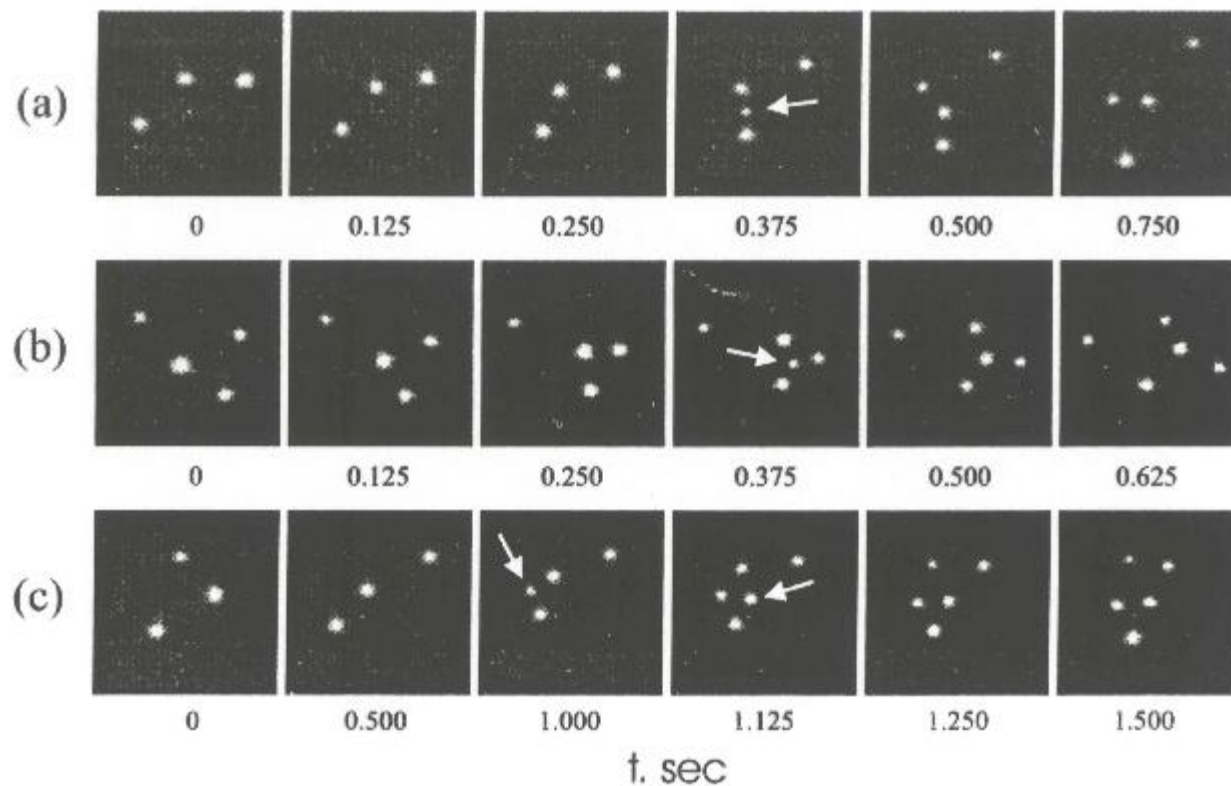


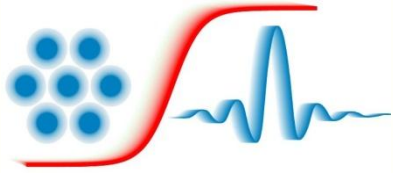
$\tau=3.35$, $\theta=0$, $D_u=1.1 \cdot 10^{-4}$, $D_y=0$, $D_w=9.64 \cdot 10^{-4}$, $\lambda=1.01$, $\kappa_1=-0.1$, $\kappa_3=0.3$, $\kappa_4=1.0$,
 $\Omega=[0,1] \times [0,1]$, $\Delta x=5 \cdot 10^{-3}$, $\Delta t=0.1$.



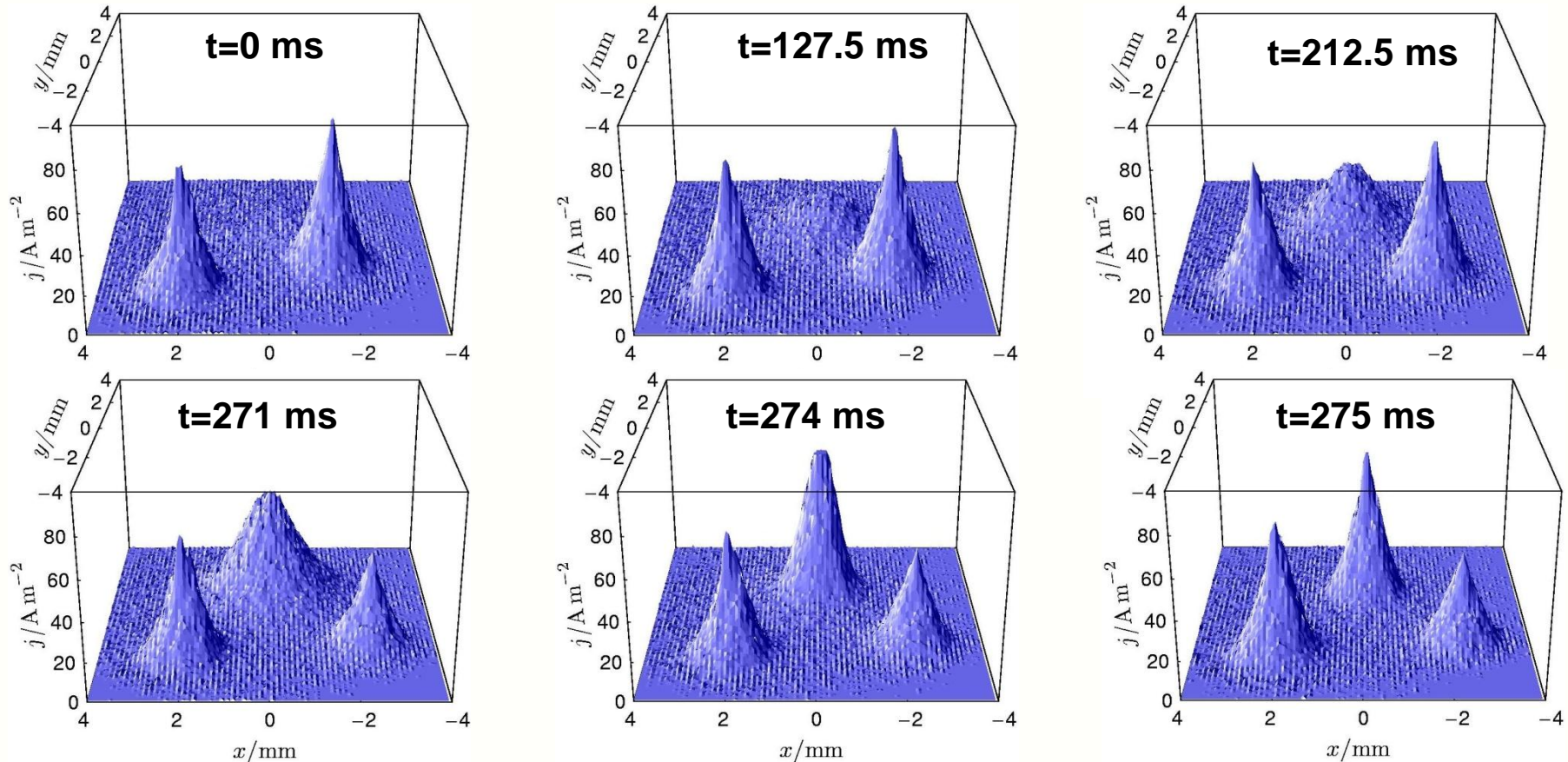
Experimental Quasi 2-Dimensional DC Gas-Discharge System: Current Filament Generation in the Course of Collision I

luminescence radiation distribution in the discharge plane

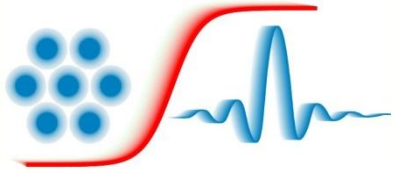




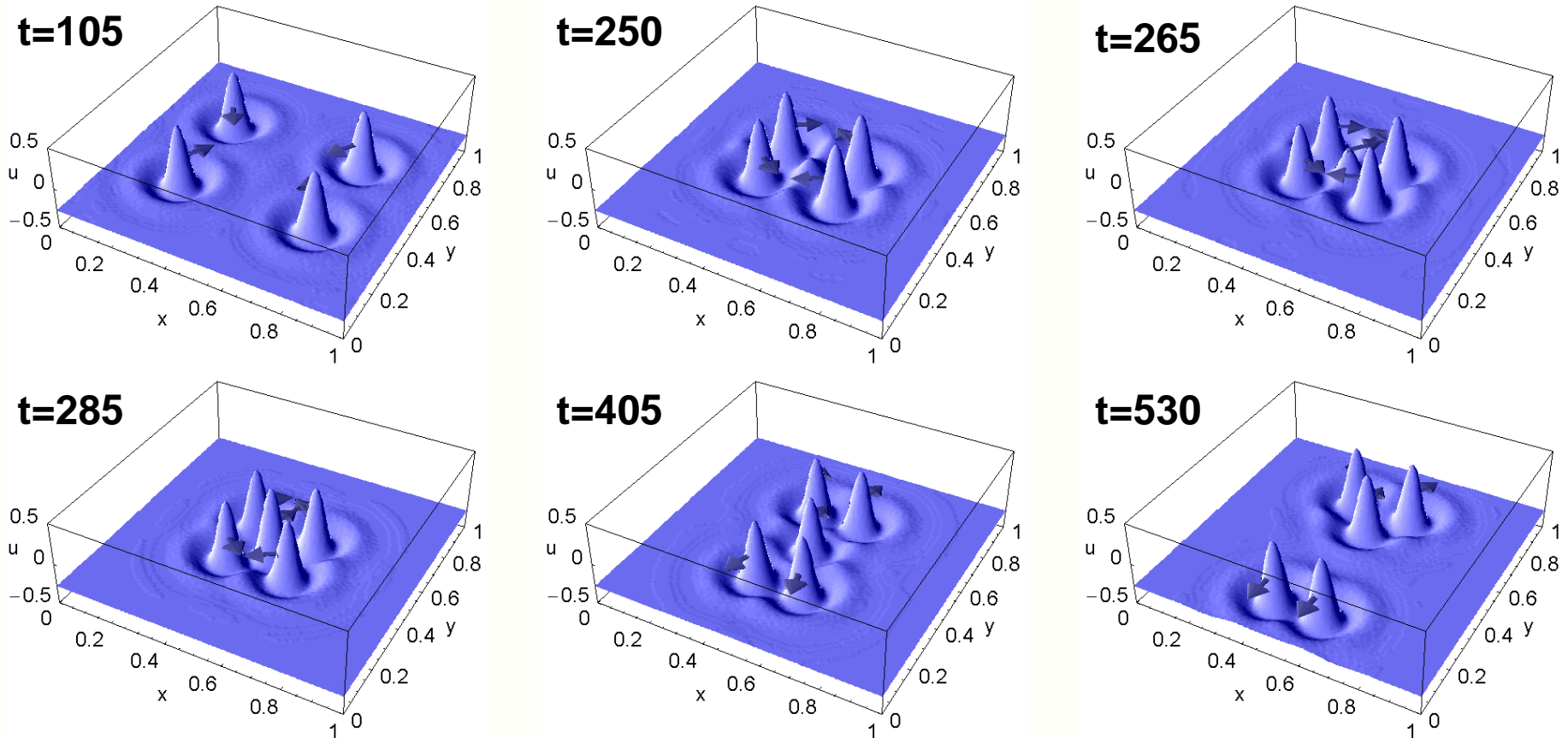
Experimental Quasi 2-Dimensional DC Gas-Discharge System: Current Filament Generation in the Course of Collision I



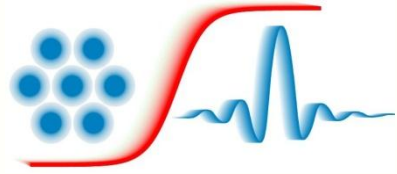
parameters: $U_0=3,8$ kV, $\rho_{SC}=4,14$ M Ω cm, $R_0=20$ M Ω , Gas: N_2 , $T=100$ K, $p=290$ hPa, $D=30$ mm, $d=500$ μm , $a_{SC}=1$ mm, $I=100\text{-}250$ μA , $t_{\text{exp}}=0,2$ ms, $f_{\text{rep}}=2$ kHz



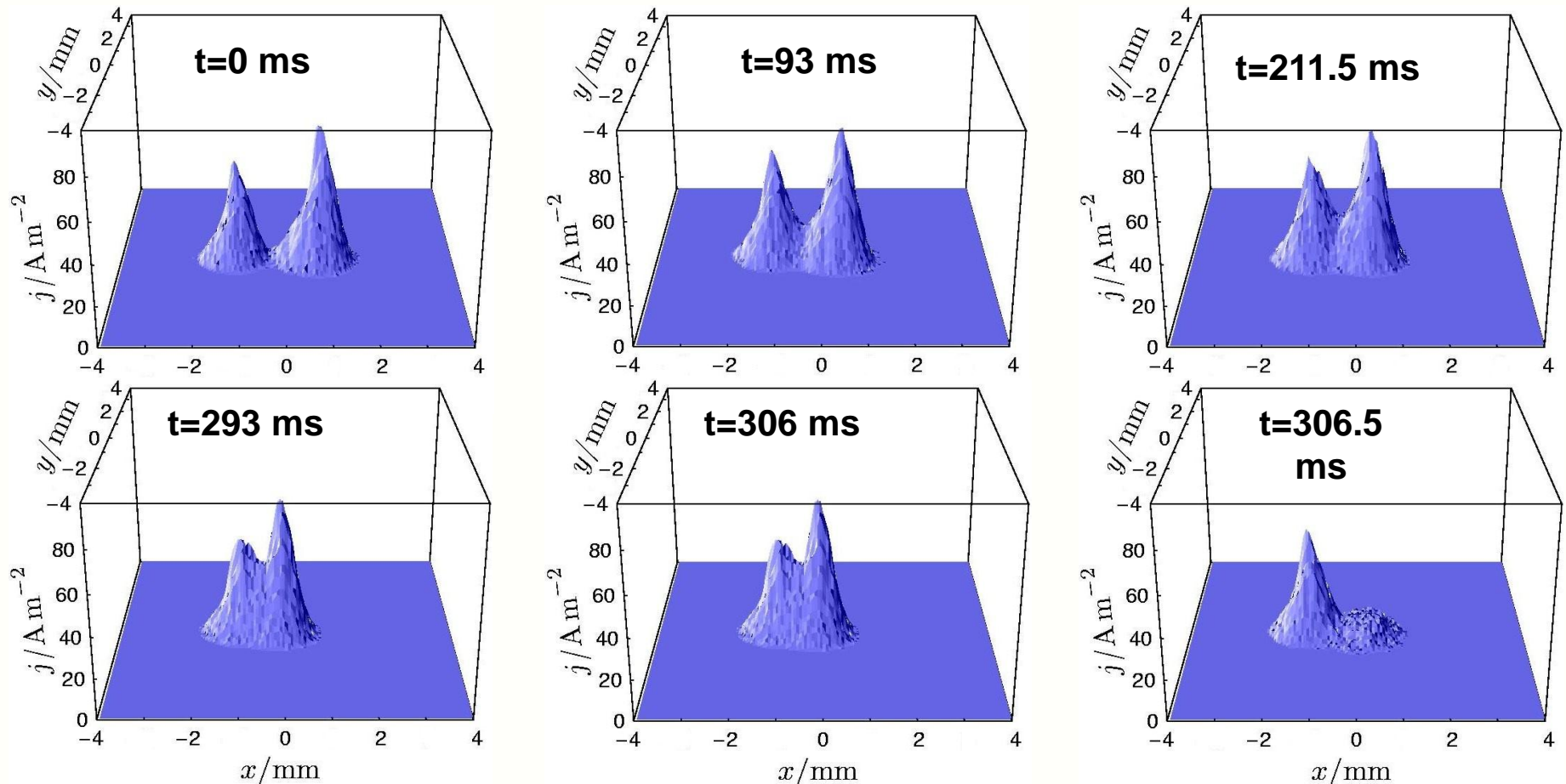
Numerical Solution for the Generalized FH Equation in \mathcal{R}^2 : DS Generation in the Course of Collision



$\tau=3.47$, $\theta=0$, $D_u=1.1 \cdot 10^{-4}$, $D_v=0$, $D_w=9.64 \cdot 10^{-4}$, $\lambda=1.01$, $\kappa_1=-0.1$, $\kappa_3=0.3$,
 $\kappa_4=1.0$, $\Omega=[0,1] \times [0,1]$, $\Delta x=5 \cdot 10^{-3}$, $\Delta t=0.1$.



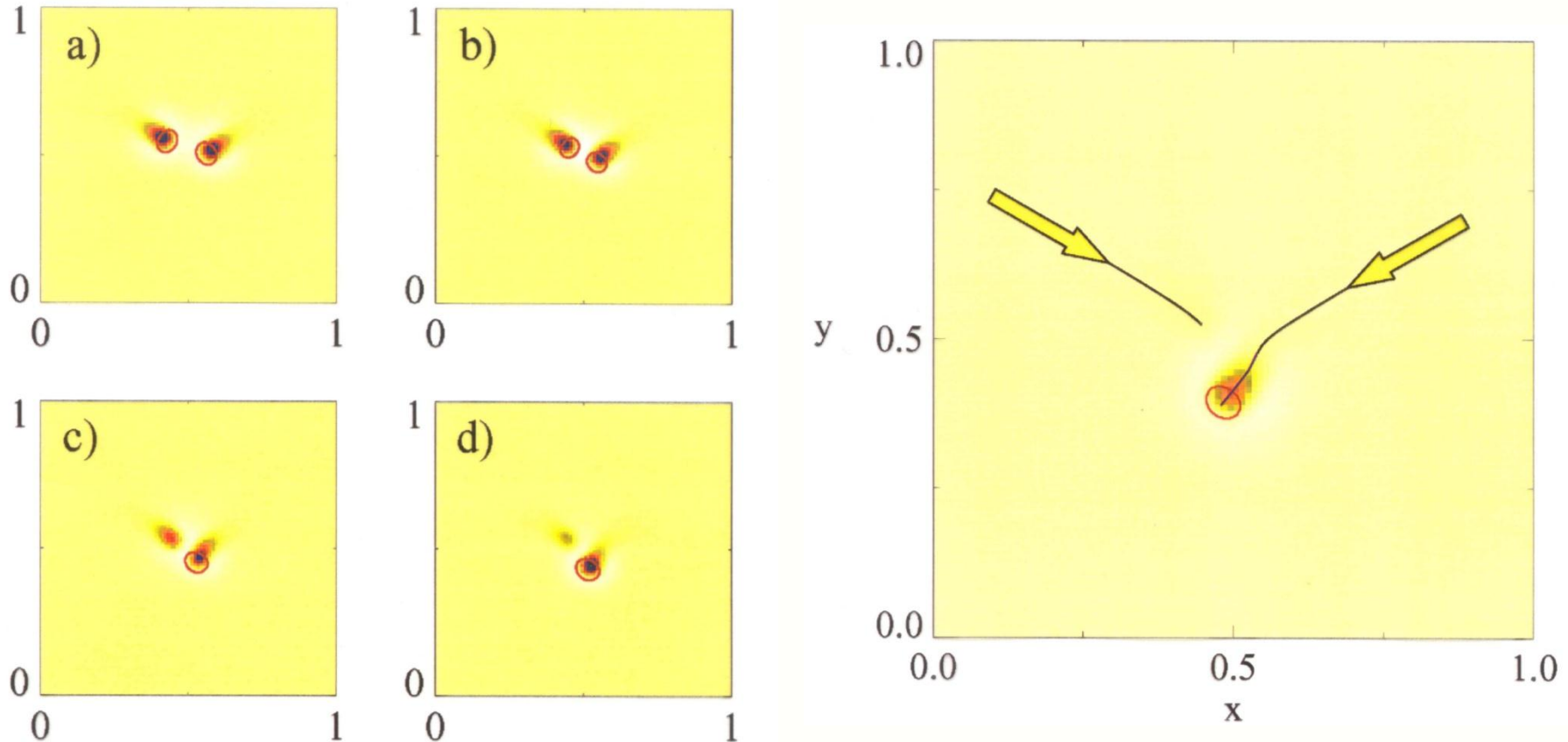
Experimental Quasi 2-Dimensional DC Gas-Discharge System: Current Filament Annihilation in the Course of Collision



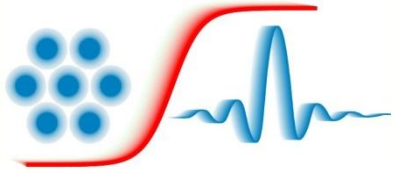
parameters: $U_0=3,8$ kV, $\rho_{\text{SC}}=4,14$ $\text{M}\Omega$ cm, $R_0=20$ $\text{M}\Omega$, Gas: N_2 , $T=100$ K, $p=290$ hPa, $D=30$ mm, $d=500$ μm , $a_{\text{SC}}=1$ mm, $I=100\text{-}250$ μA , $t_{\text{exp}}=0,2$ ms, $f_{\text{rep}}=2$ kHz



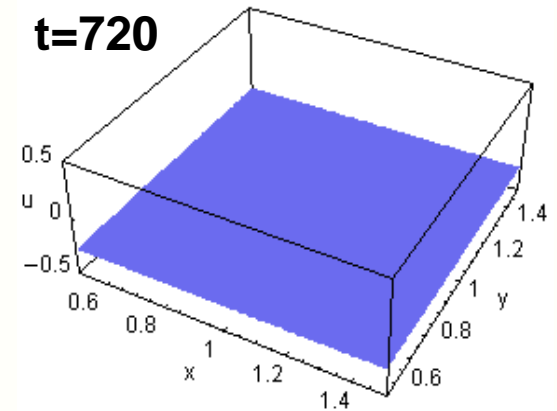
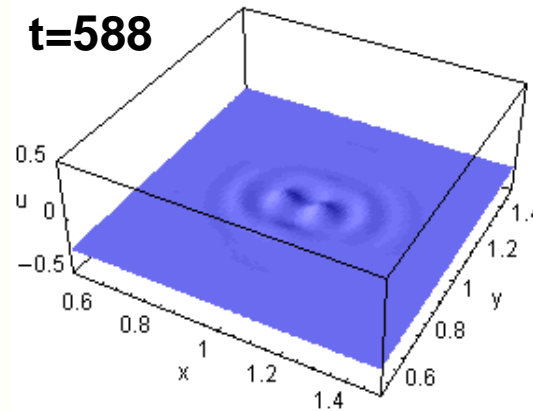
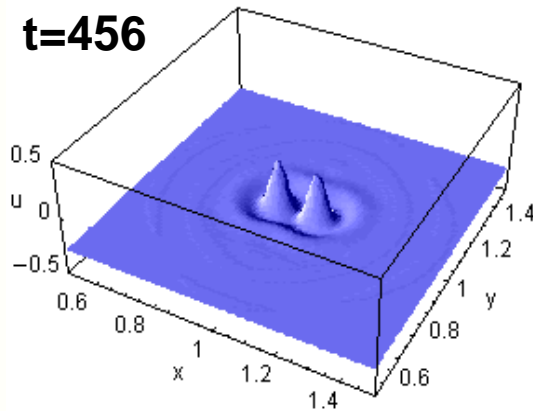
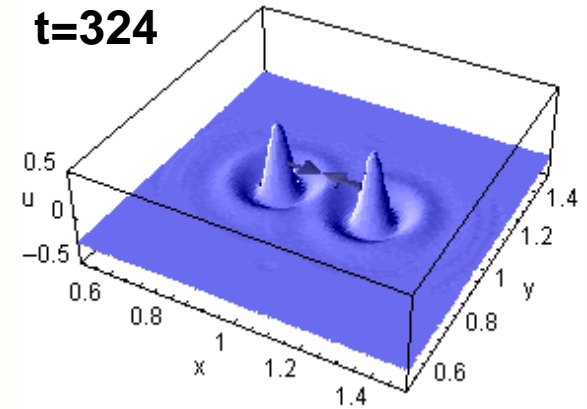
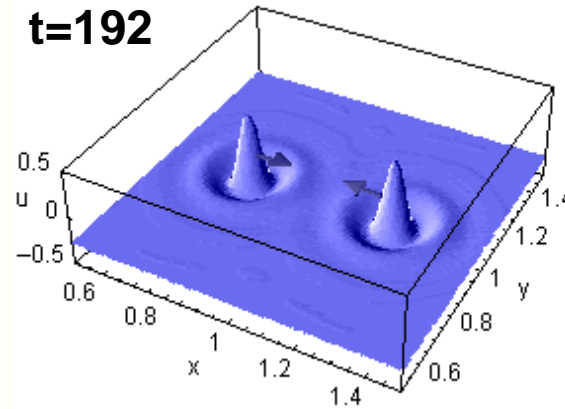
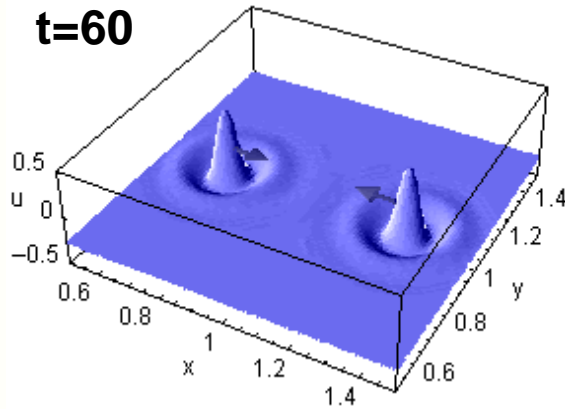
Numerical Solution for the Generalized FH Equation in \mathcal{R}^2 : DS Annihilation in the course of Collision I



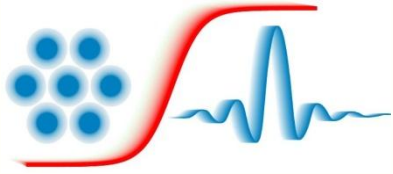
parameters: $D_u=1.55 \cdot 10^{-4}$, $D_v=1.95 \cdot 10^{-4}$, $D_w=9.6 \cdot 10^{-4}$, $\kappa_1=-8.92$, $\kappa_3=1$, $\kappa_4=8.5$, $\lambda=2$,
 $\tau=25$, $\theta=1$, $\Delta x=1/100$



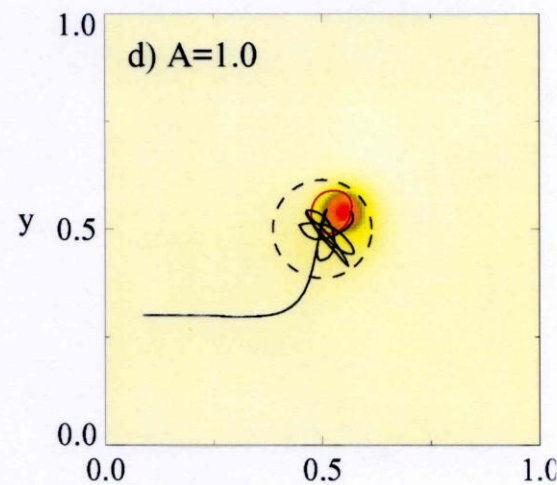
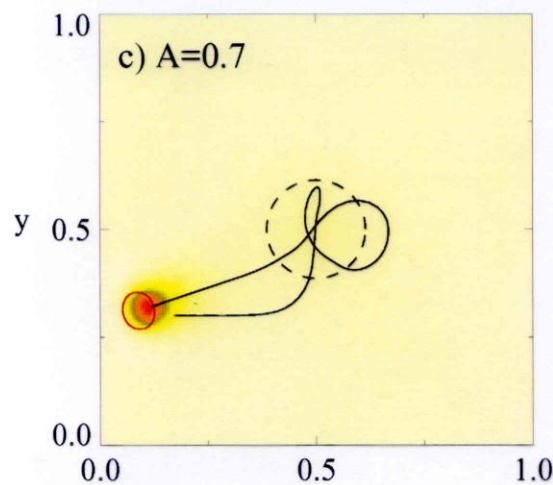
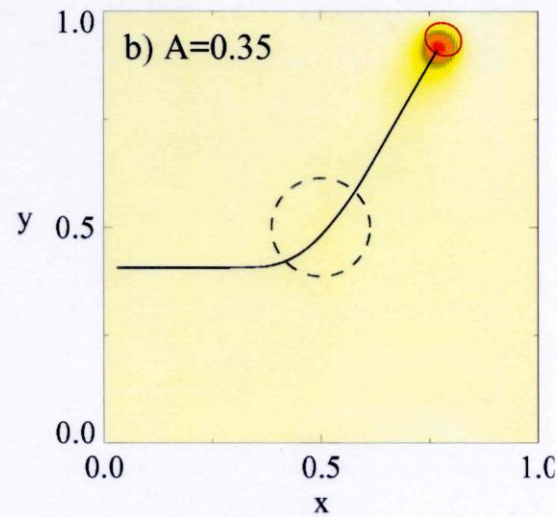
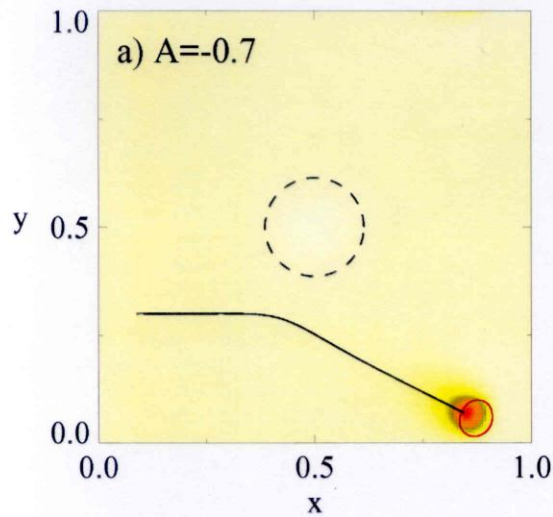
Numerical Solution for the Generalized FH Equation in \mathcal{R}^2 : DS Annihilation in the course of Collision II



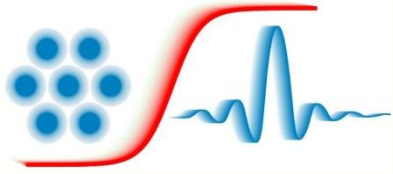
$\tau=3.59$, $\theta=0$, $D_u=1.1 \cdot 10^{-4}$, $D_v=0$, $D_w=9.64 \cdot 10^{-4}$, $\lambda=1.01$, $\kappa_1=-0.1$, $\kappa_3=0.3$, $\kappa_4=1.0$,
 $\Omega=[0,1] \times [0,1]$, $\Delta x=5 \cdot 10^{-3}$, $\Delta t=0.1$.



Numerical Solution for the Generalized FH Equation in \mathcal{R}^2 : Scattering of a DS at an Inhomogeneity

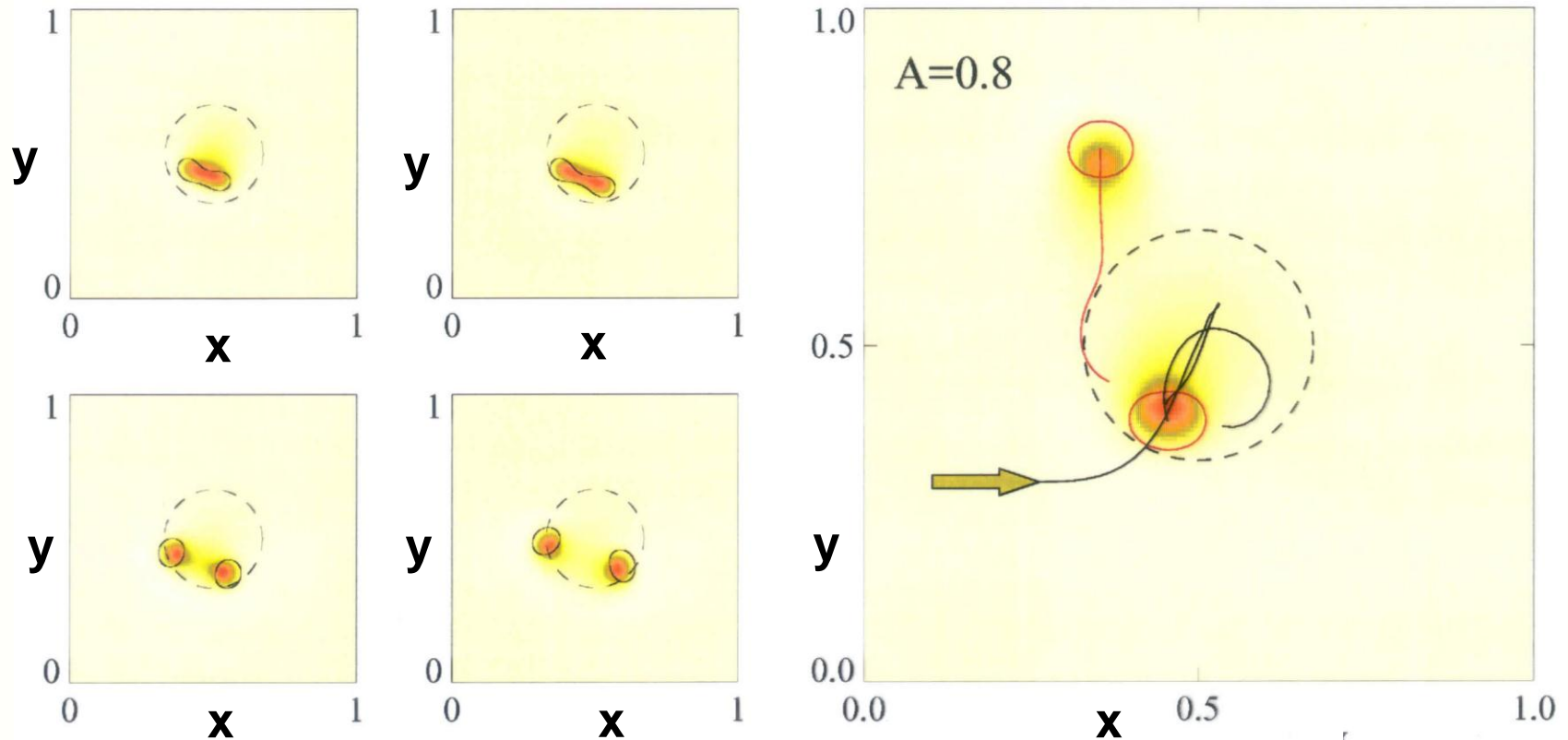


parameters: $D_u=1.21 \cdot 10^{-4}$,
 $D_v=1.70 \cdot 10^{-3}$, $D_w=1.46 \cdot 10^{-2}$,
 $\kappa_3=8$, $\lambda=4.3$, $\tau=60$, $\theta=0$, $\Delta x=1/160$
 $\kappa_1=-12.34+A \exp(-53.3((x-0.5)^2+(y-0.5)^2))$



Numerical Solution for the Generalized FH Equation in \mathcal{R}^2 : Generation of a DS at an Inhomogeneity

inhomogeneity as particle generator

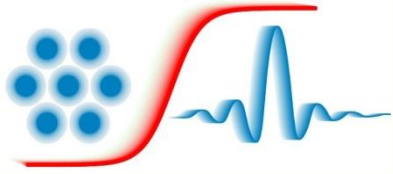


$$D_u=1.21 \cdot 10^{-4}, D_v=1.70 \cdot 10^{-3}, D_w=1.46 \cdot 10^{-2}, \lambda=4.3, \kappa_3=8, \tau=60, \theta=0, \Delta x=1/160$$

$$\kappa_1=-12.34+ A \exp(-35.5 ((x-0.5)^2+(y-0.5)^2))$$



5. Theoretical Foundation of a Particle Concept



Particle Equation for Interacting DSs Derived from the Generalized FH Equation near to the Travelling Bifurcation Point

$$\partial_t \underline{p}_i = \underline{\alpha}_i + \frac{1}{\kappa_3} F(|\underline{p}_i - \underline{p}_j|) \frac{\underline{p}_i - \underline{p}_j}{|\underline{p}_i - \underline{p}_j|},$$

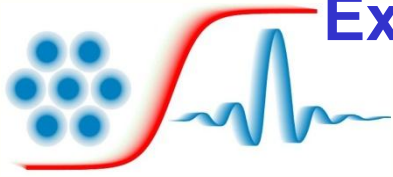
$$\partial_t \underline{\alpha}_i = (\tau - \tau_c) \kappa_3^2 \underline{\alpha}_i + \frac{1}{\kappa_3} \frac{\langle (\partial_x^2 u_0)^2 \rangle}{\langle (\partial_x u_0)^2 \rangle} |\underline{\alpha}_i|^2 \underline{\alpha}_i + F(|\underline{p}_i - \underline{p}_j|) \frac{\underline{p}_i - \underline{p}_j}{|\underline{p}_i - \underline{p}_j|},$$

\underline{p}_i = position coordinate of the i th DS

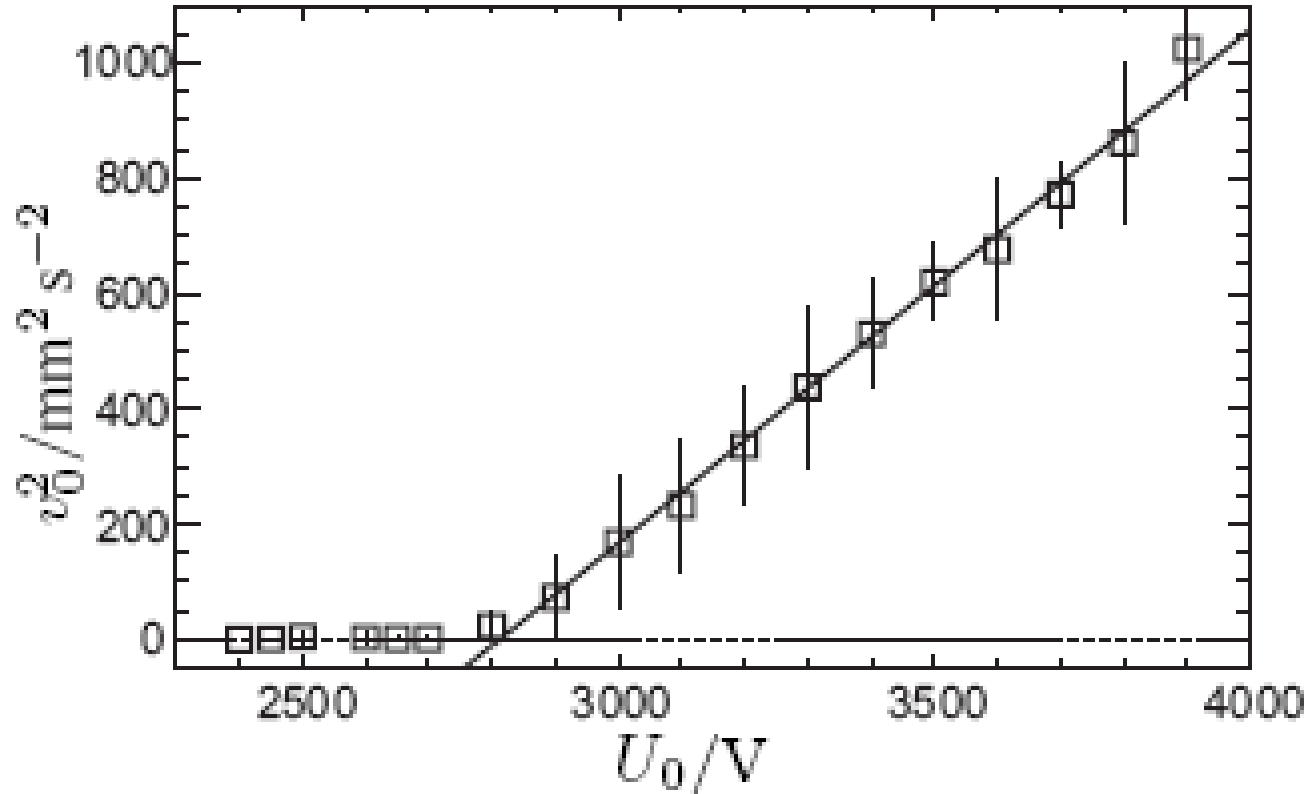
$\underline{\alpha}_i$ = displacement of the slow inhibitor with respect to the activator

u_0 = solution for the activator component at the bifurcation point

$F(|\underline{p}_i - \underline{p}_j|)$ = interaction function



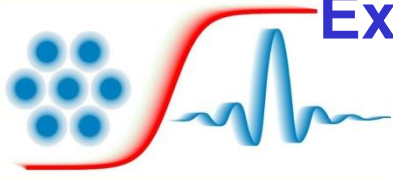
Experimental Quasi 2-Dimensional DC Gas-Discharge System: Comparison of the Experimental to the Theoretical Scaling Law for the Velocity of DSs I



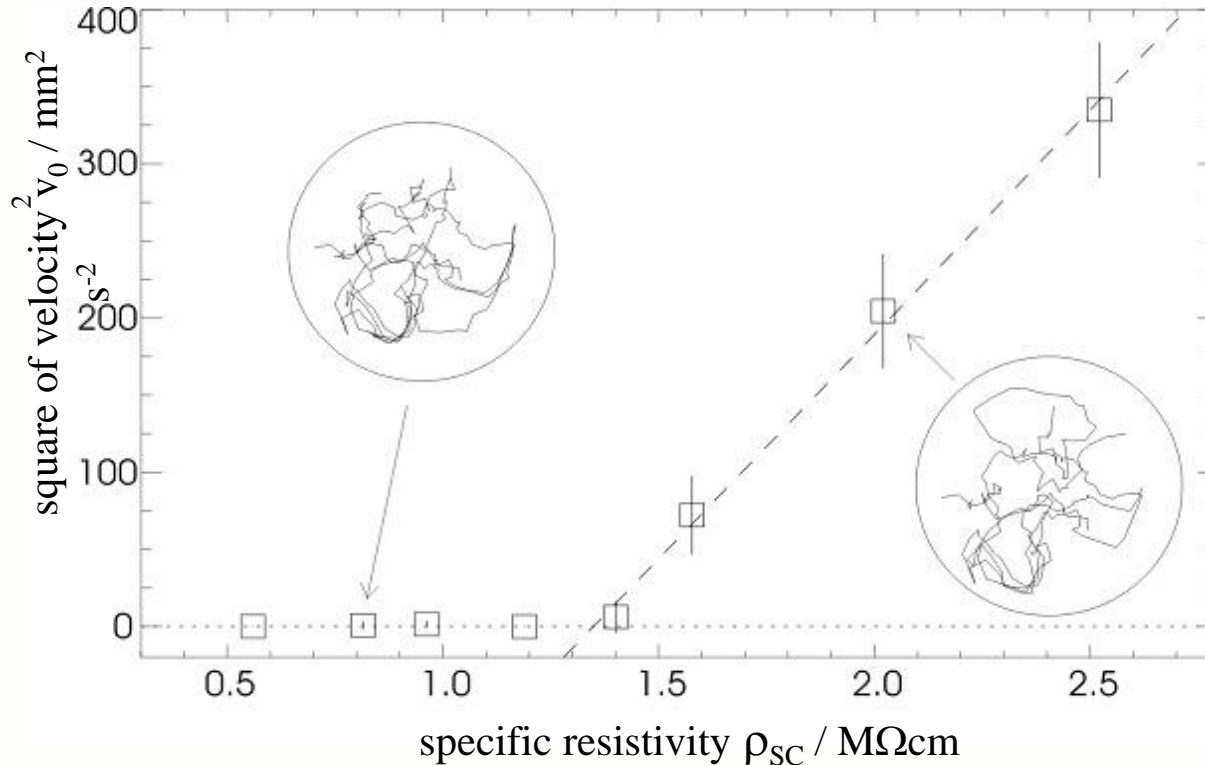
parameters:

$\rho_0=1,83 \text{ M}\Omega$, $R_0=10 \text{ M}\Omega$, Gas: N_2 , $T=100 \text{ K}$,
 $p=280 \text{ hPa}$, $d=500 \text{ }\mu\text{m}$,
 $t_{\text{exp}}=20 \text{ ms}$

square of the intrinsic velocity of a single filament plotted as a function of the driving voltage U_0 (points); the velocity is evaluated by advanced statistical data processing; the straight line is deduced from the generalized FN equation



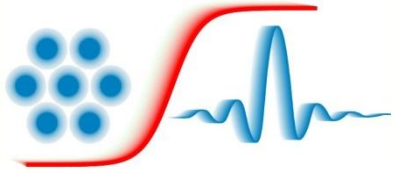
Experimental Quasi 2-Dimensional DC Gas-Discharge System: Comparison of the Experimental to the Theoretical Scaling Law for the Velocity of DSs II



parameters:

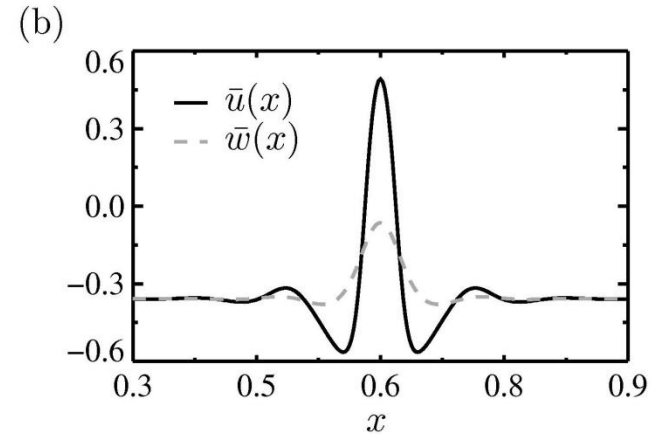
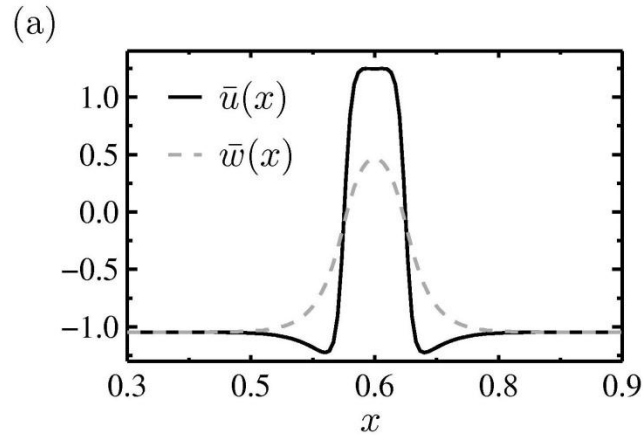
$U_0=3,7 \text{ kV}$, $R_0=10 \text{ M}\Omega$,
Gas: N_2 , $T=100 \text{ K}$,
 $p=286 \text{ hPa}$, $D=30 \text{ mm}$,
 $d=750 \mu\text{m}$, $a_{SC}=1 \text{ mm}$,
 $I=107 \mu\text{A}$, $t_{\text{exp}}=20 \text{ ms}$,
 $f_{\text{rep}}=50 \text{ Hz}$

square of the intrinsic velocity of a single filament plotted as a function of the specific resistivity of the semiconductor wafer (points); the velocity is evaluated by advanced statistical data processing; the straight line is deduced from the generalized FN equation; typical experimental trajectories are displayed in the insert

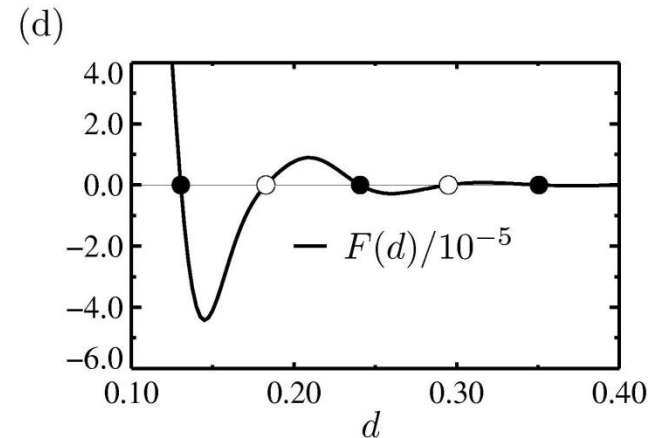
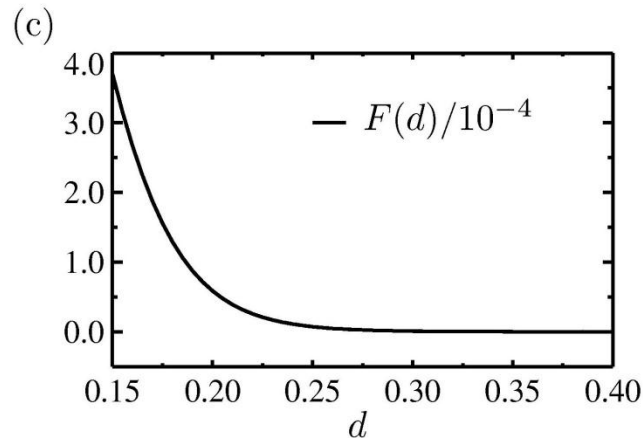


Numerical Solution for the Generalized FH Equation in \mathcal{R}^1 : Interaction Function for Dissipative Solitons

$$\bar{\mathbf{u}}(\mathbf{x}) = \bar{\mathbf{v}}(\mathbf{x})$$

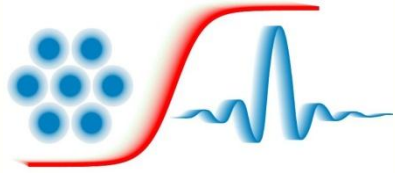


$$\mathbf{F}(\mathbf{d}) = \mathbf{F}(|\vec{\mathbf{p}}_2 - \vec{\mathbf{p}}_1|)$$



a+c: $\tau=0, \theta=0, D_u=0.5 \cdot 10^{-4}, D_v=0, D_w=10^{-3}, \lambda=3.0, \kappa_1=-0.1, \kappa_3=1.0, \kappa_4=1.0, \Omega=[0,1.2],$
 $\Delta x=5 \cdot 10^{-3},$

b+d: $D_u=1.1 \cdot 10^{-4}, D_w=9.64 \cdot 10^{-4}, \lambda=1.71, \kappa_1=-0.15, \Delta x=2.5 \cdot 10^{-3},$ others as in a+c.

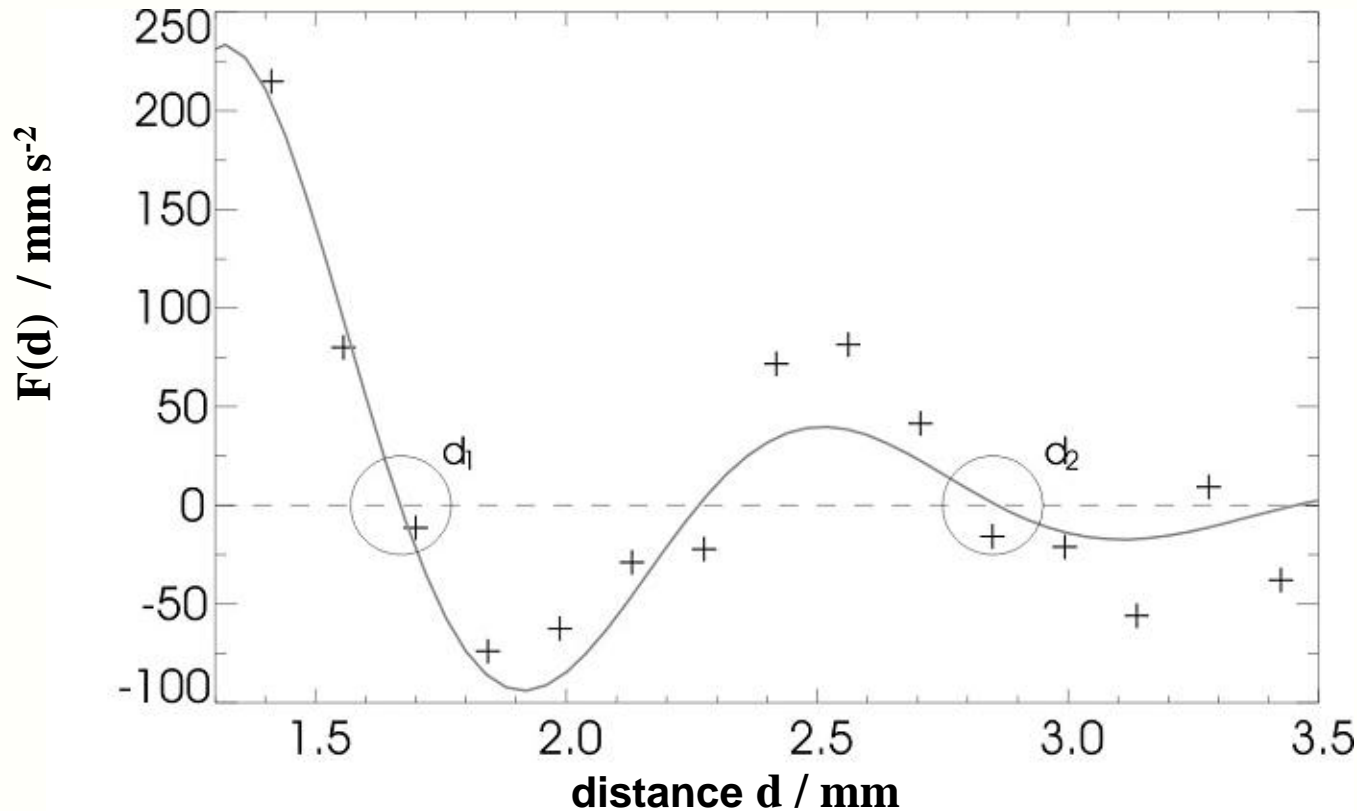


Selected Additional Theoretical Work on the Generalized FitzHugh-Nagumo Equation

- M. Krupa, B. Sandstede, P. Szmolyan, “Fast and slow waves in the FitzHugh-Nagumo equation”,
J. Differential Equations, vol. 133, p. 49 (1997)**
- Y. Nishiura et al., “Dynamics of Travelling Pulses in Heterogeneous Media”,
Chaos vol. 17, p. 037104 (2007), see also references therein**
- P. v. Heijster, A. Doelman, T. S. Kaper , “Pulse dynamics in 3-component systems: Stability and bifurcations”,
Physica D vol. 237 p. 3335 (2008), see also references therein**
- P. v. Heijster, B. Sandstede, “Planar radial spots in a three component FitzHugh Nagumo equation”,
math.bu.edu/people/heijster/PAPERS/planar_spots.pdf, see also references therein**



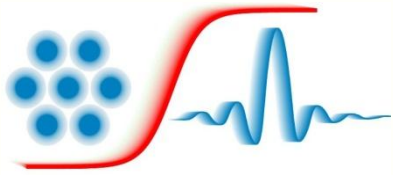
Experimental Quasi 2-Dimensional DC Gas-Discharge System: Determination of the Interaction Function of the Particle Equation from Experimental Trajectories



parameters:

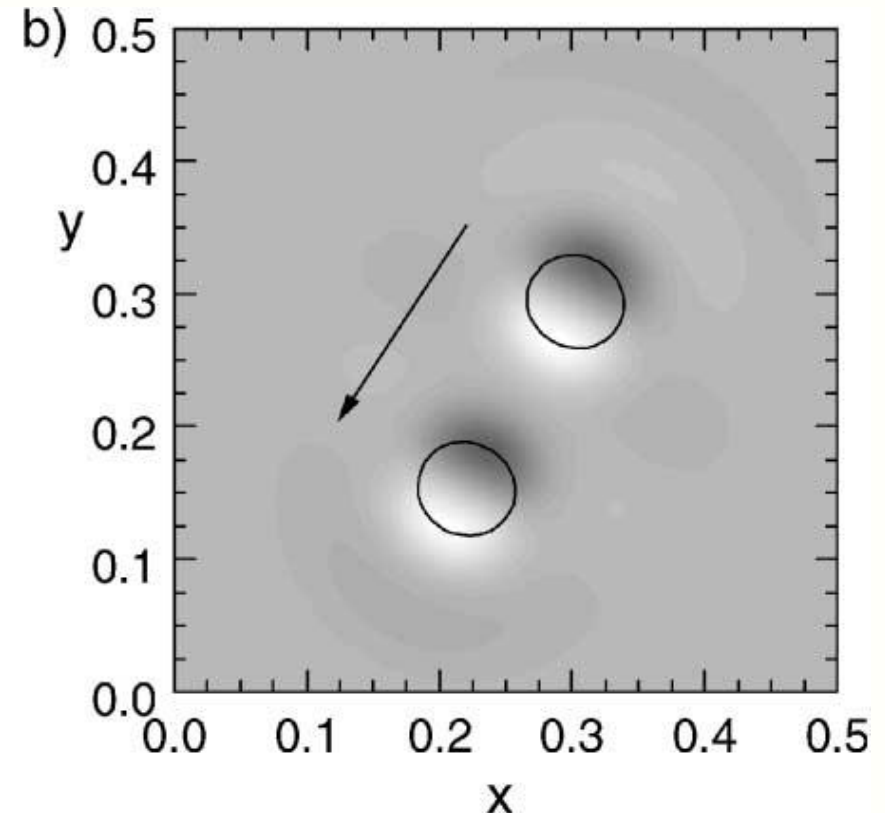
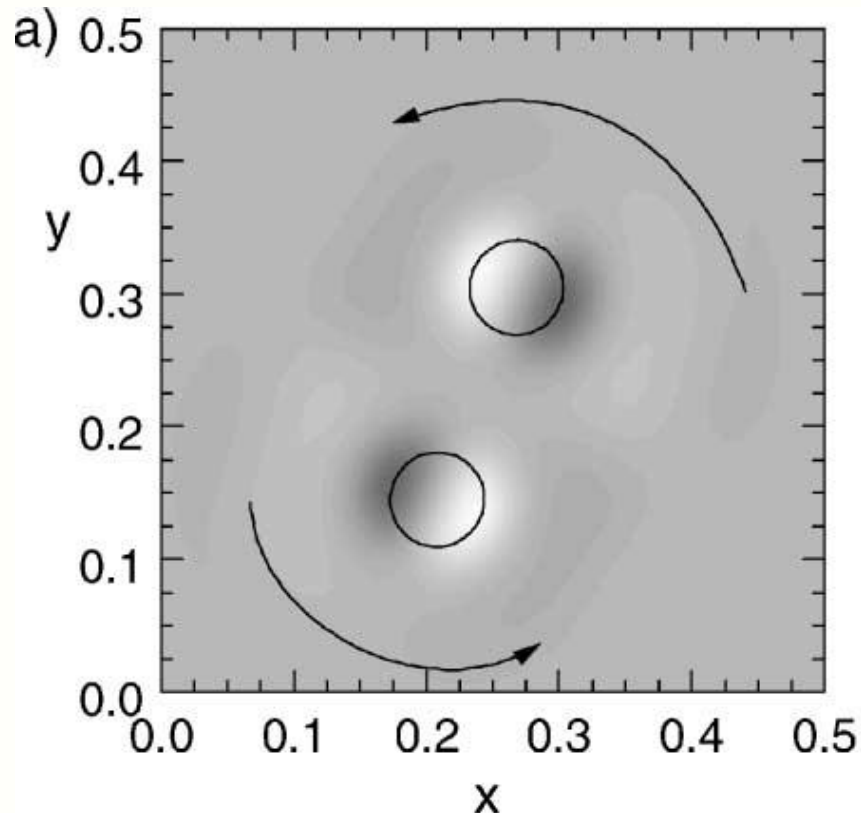
$U_0 = 4600 \text{ V}$
 $\rho_{\text{SC}} = 3,55 \text{ M}\Omega\text{cm}$
 $R_0 = 4,4 \text{ M}\Omega$,
Gas: N_2
 $T = 100 \text{ K}$
 $p = 283 \text{ hPa}$
 $D = 30 \text{ mm}$,
 $d = 550 \mu\text{m}$
 $a_{\text{SC}} = 1 \text{ mm}$
 $I = 233 \mu\text{A}$
 $t_{\text{exp}} = 20 \text{ ms}$,
 $f_{\text{rep}} = 50 \text{ Hz}$

the interaction function $F(d) = F(|\underline{p}_i - \underline{p}_j|)$ is evaluated by advanced statistical data processing; the continuous line is deduced from the generalized FN equation



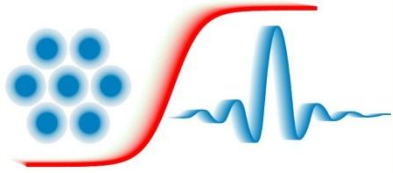
Numerical Solution for the Particle Equation in \mathcal{R}^2 : Scattering Diagram for a Fixed Set of Parameters III

solution of the generalized FN equation:



parameters of the original equation:

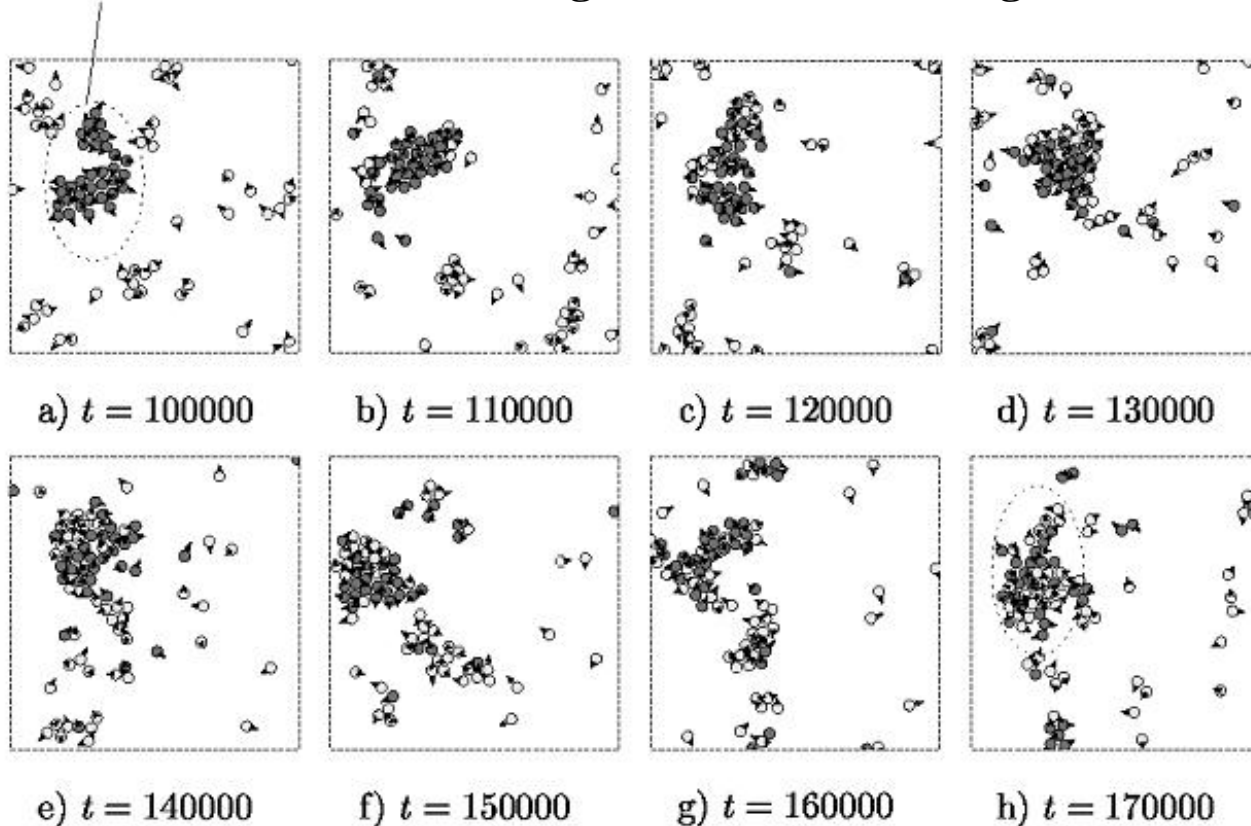
$$D_u=1.1 \cdot 10^{-4}, D_v=0, D_w=9.64 \cdot 10^{-4}, \kappa_1=-0.1, \kappa_3=0.3, \kappa_4=1, \lambda=1.01, \tau=3.43, \theta=0$$



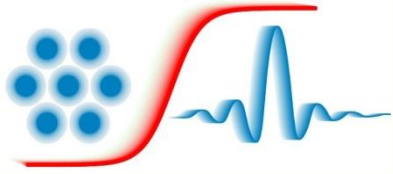
Numerical Solution for the Particle Equation in \mathcal{R}^2 : Possible Behaviour of a Many-DS System

marked solitons in the largest cluster at starting time

N=81

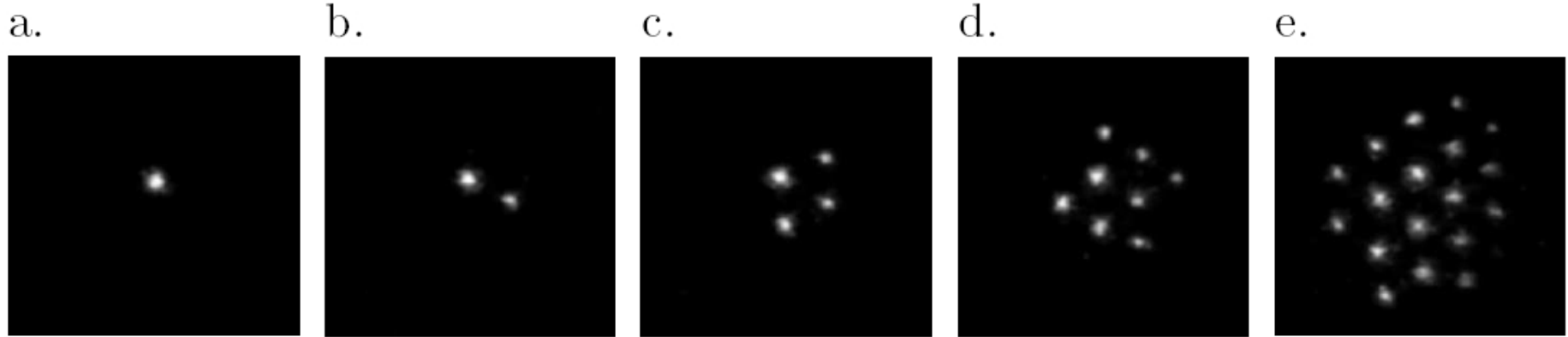


$\tau=3.34$, $\theta=0$, $D_u=1.1 \cdot 10^{-4}$, $D_v=0$, $D_w=9.64 \cdot 10^{-4}$, $\lambda=1.01$, $\kappa_1=-0.1$, $\kappa_3=0.3$, $\kappa_4=1.0$,
 $\Omega=[-2,2] \times [-2,2]$, $\Delta x=5 \cdot 10^{-3}$, $\Delta t=0.1$; F(d), Q = 1950.

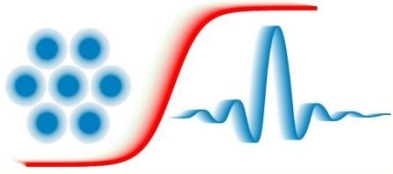


Experimental Quasi 2-Dimensional DC Gas-Discharge System: Self Completion I

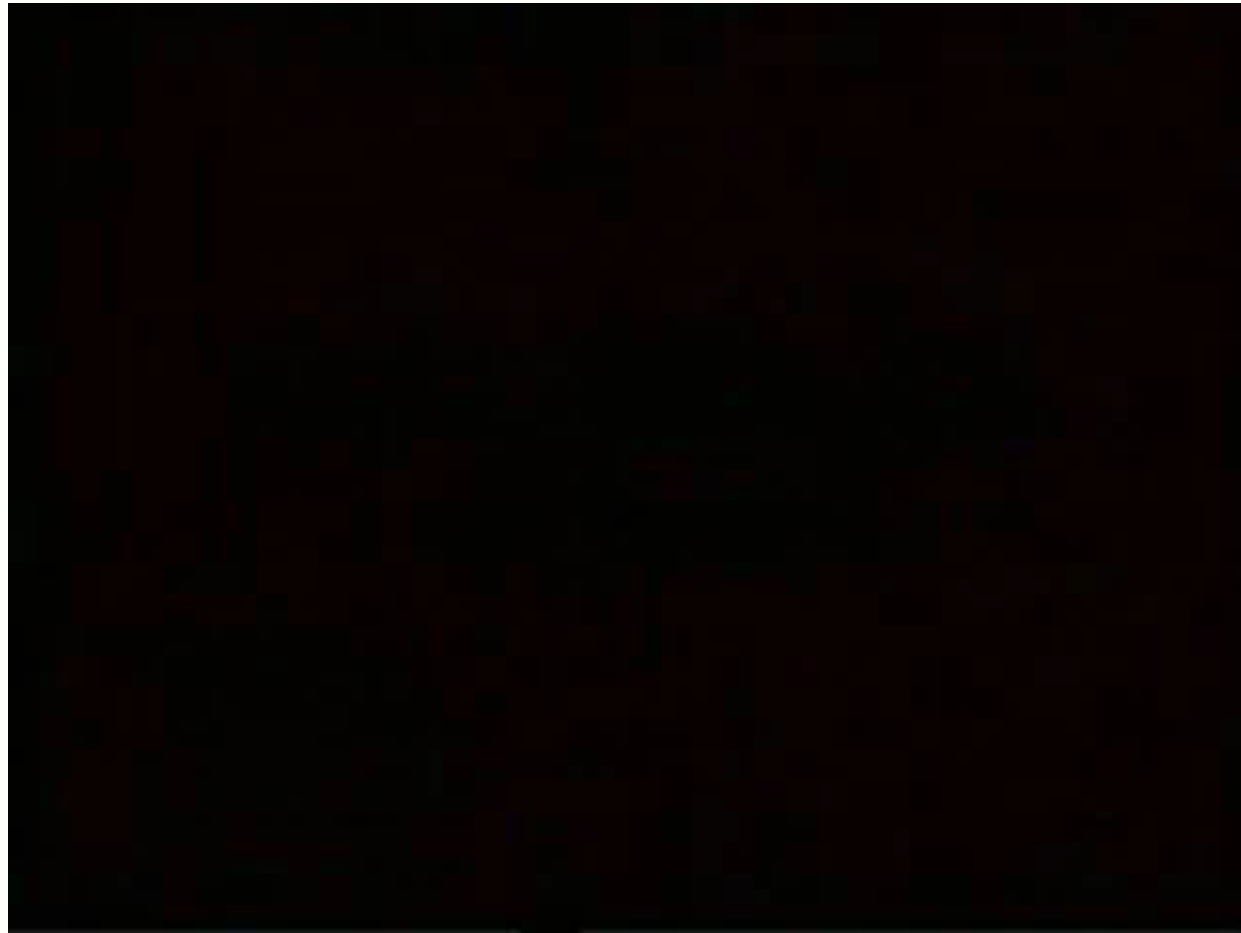
luminescence radiation distribution in the discharge plane



at fixed parameter in the course of time near to an existing localized structure an increasing number of such objects is generated at well defined distance to existing ones until the plain is fully covered with a hexagonal pattern



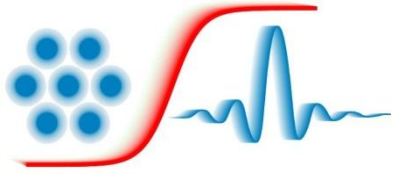
Experimental Quasi 2-Dimensional DC Gas-Discharge System: Self Completion II (Movie)



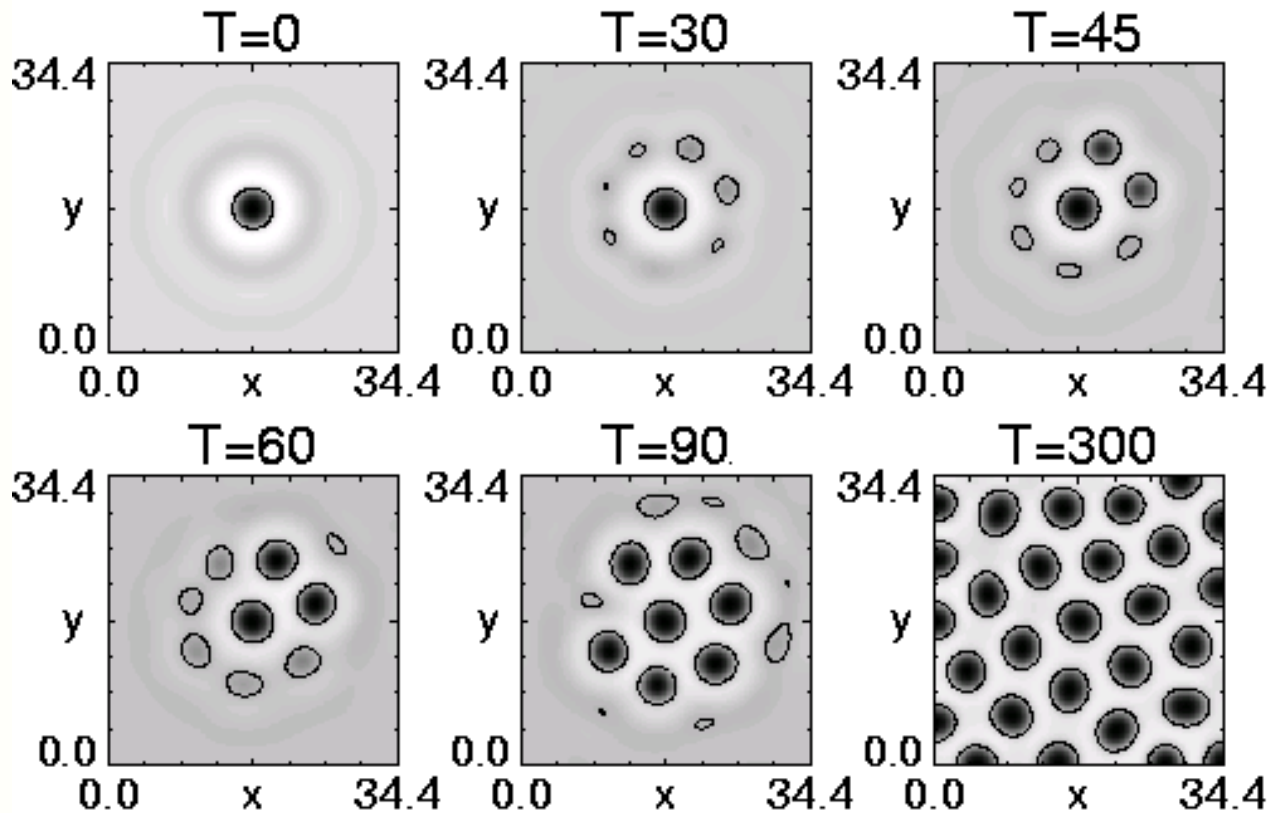
**luminescence
radiation
distribution in
the discharge
plane**

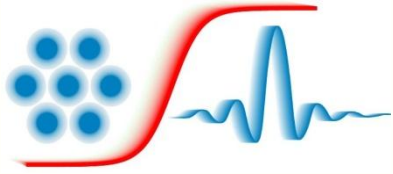
**if not linked:
start movie
“self
completion.
wmv”
in the folder**

radiation distribution in the discharge plane

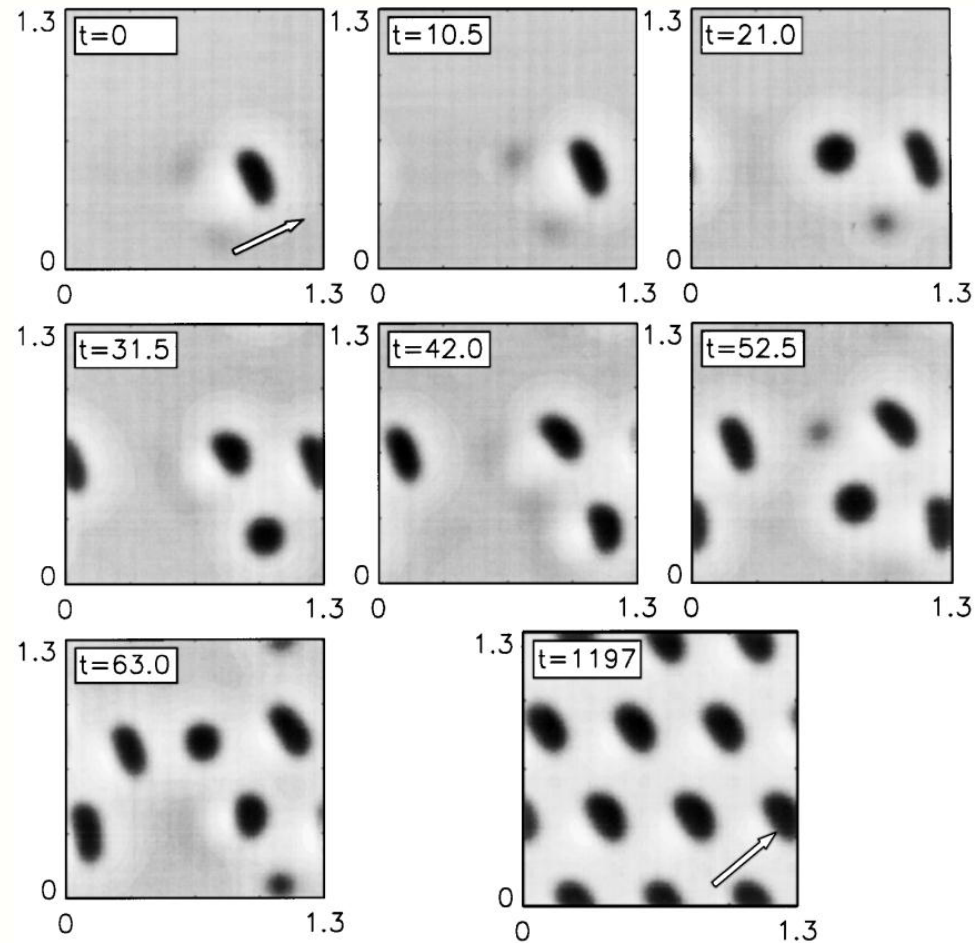


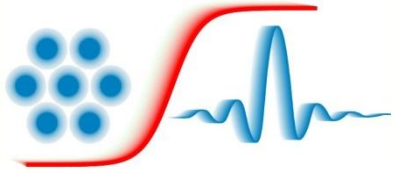
Numerical Solutions for the Generalized FH Equation in \mathbb{R}^2 : Self-Completion I



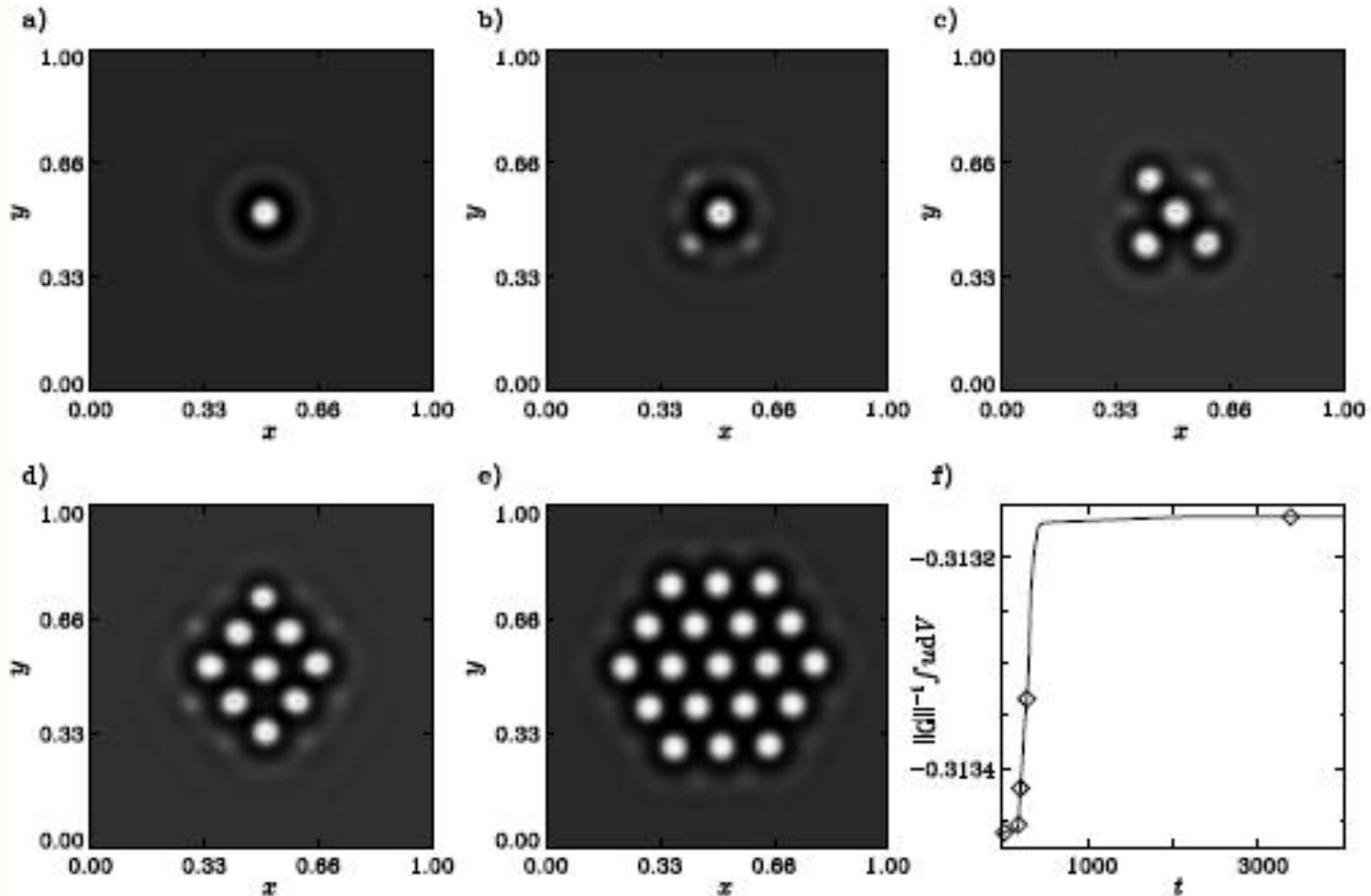


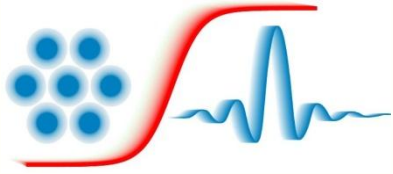
Numerical Solutions for the Generalized FH Equation in \mathcal{R}^2 : Self-Completion II



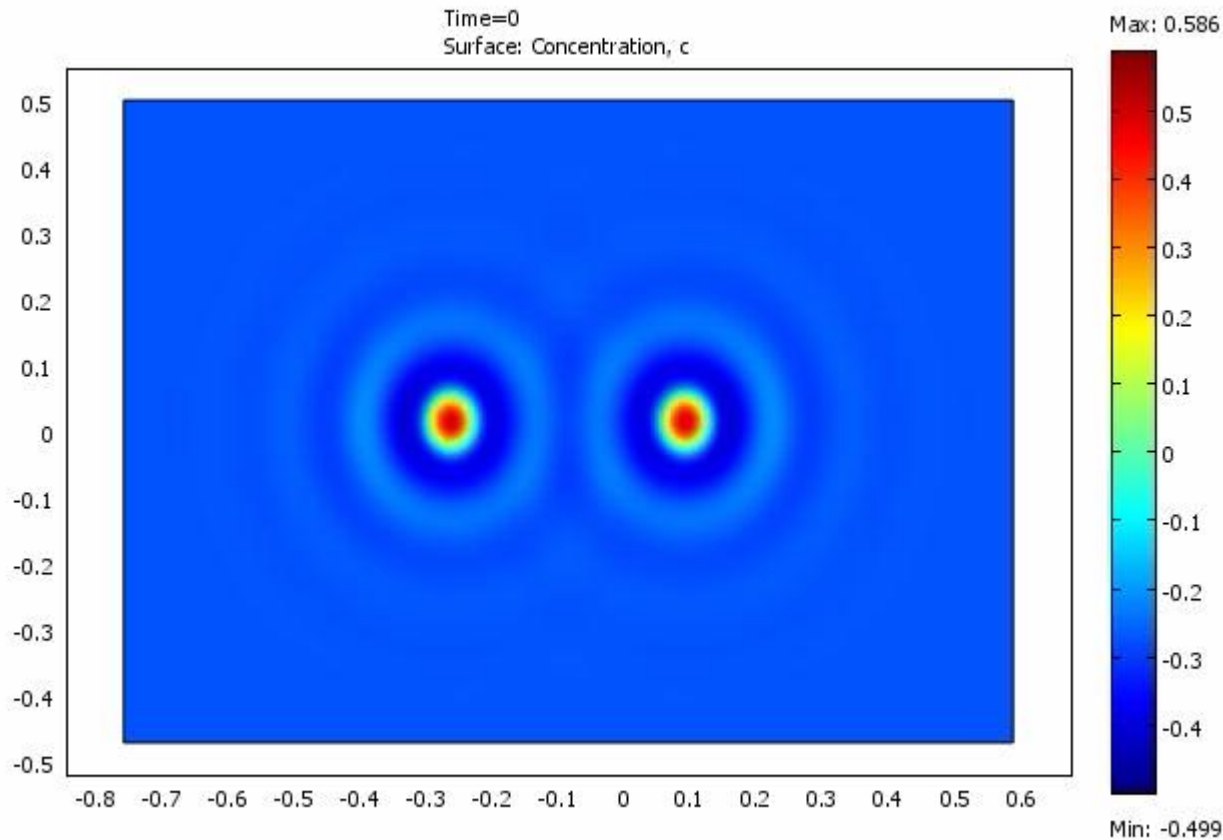


Numerical Solutions for the Generalized FH Equation in \mathbb{R}^2 : Self-Completion III

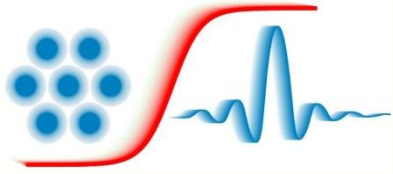




Numerical Solutions for the Generalized FH Equation in \mathbb{R}^2 : Self-Completion IV (Movie)

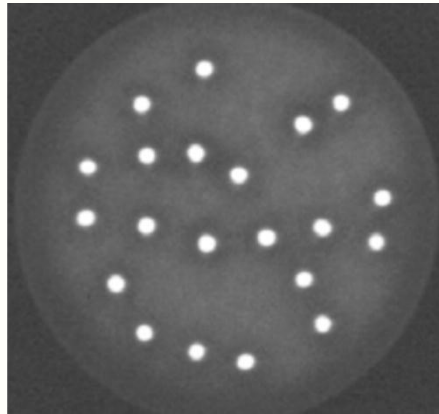


if not linked:
start movie
“self_completion.
avi” in the folder

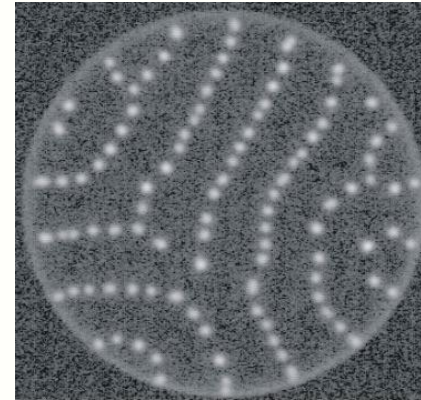


Experimental Quasi 2-Dimensional DC Gas-Discharge System: Various Many-Filament Patterns I

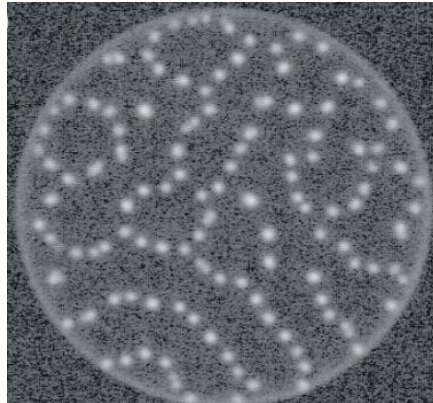
„gas“



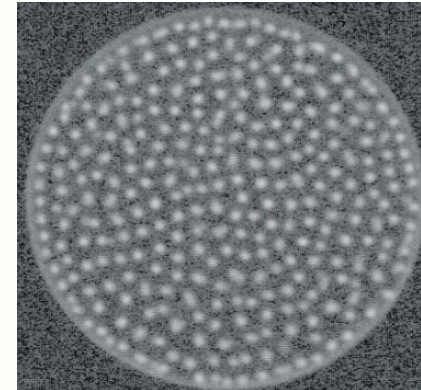
dynamic
chains



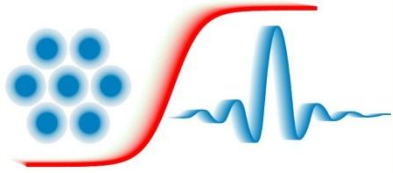
dynamic
net



„liquid“

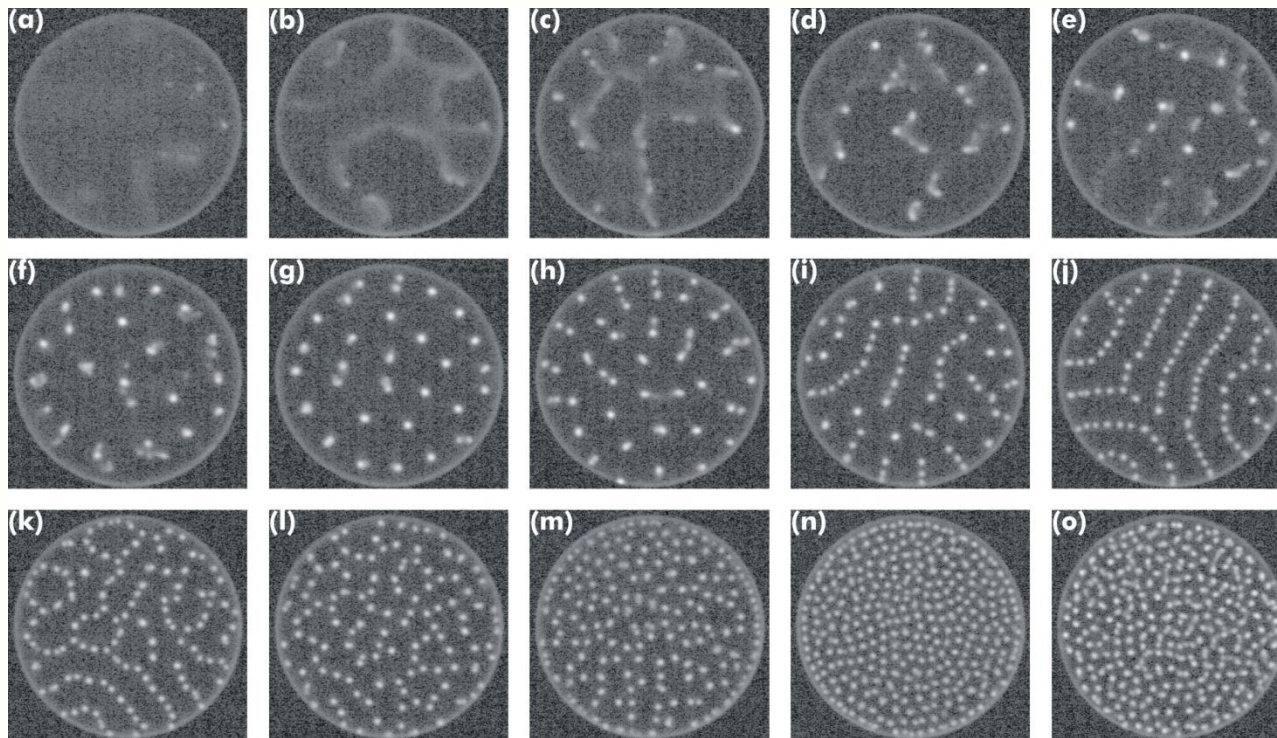


luminescence radiation distribution in the discharge plane

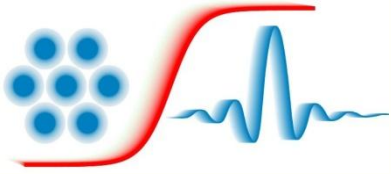


Experimental Quasi 2-Dimensional DC Gas-Discharge System: Various Many-Filament Patterns II

sequence of patterns: luminescence radiation distribution in the discharge plane

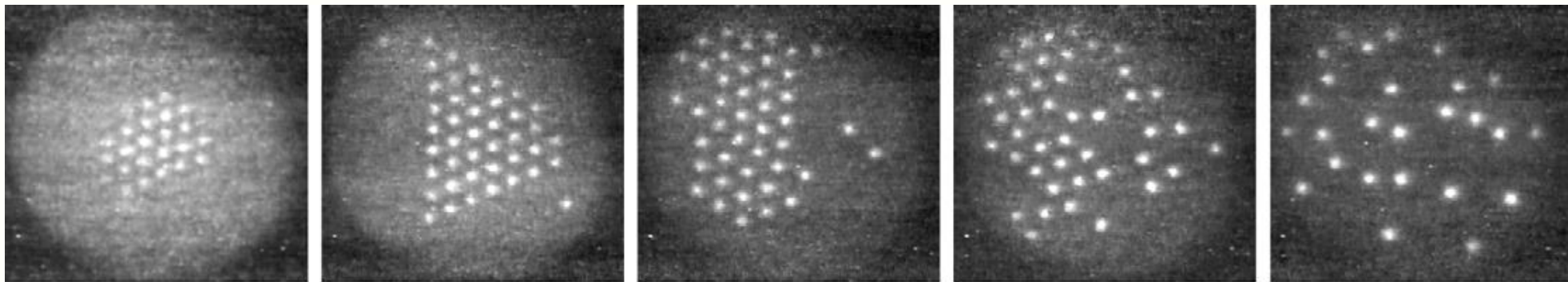


parameters: increasing voltage from 572 V to 1208 V for a) to o)
 $p=41$ hPa, $D=3$ cm, $d=0.05$ cm, $t_{\text{exp}}=40$ ms



Experimental Quasi 2 Dimensional DC Gas-Discharge System: Various Many-Filament Patterns III

luminescence radiation distribution in the discharge plane:
equilibrium between „crystal“- and „liquid“-many-filament patterns



A

B

C

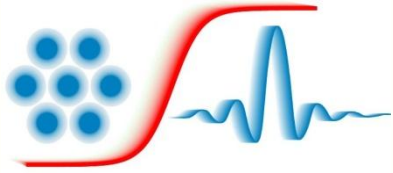
D

E

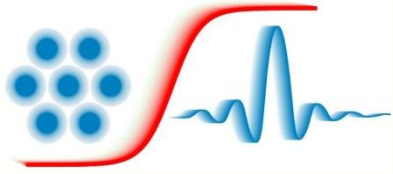
radiation distribution in the discharge plane



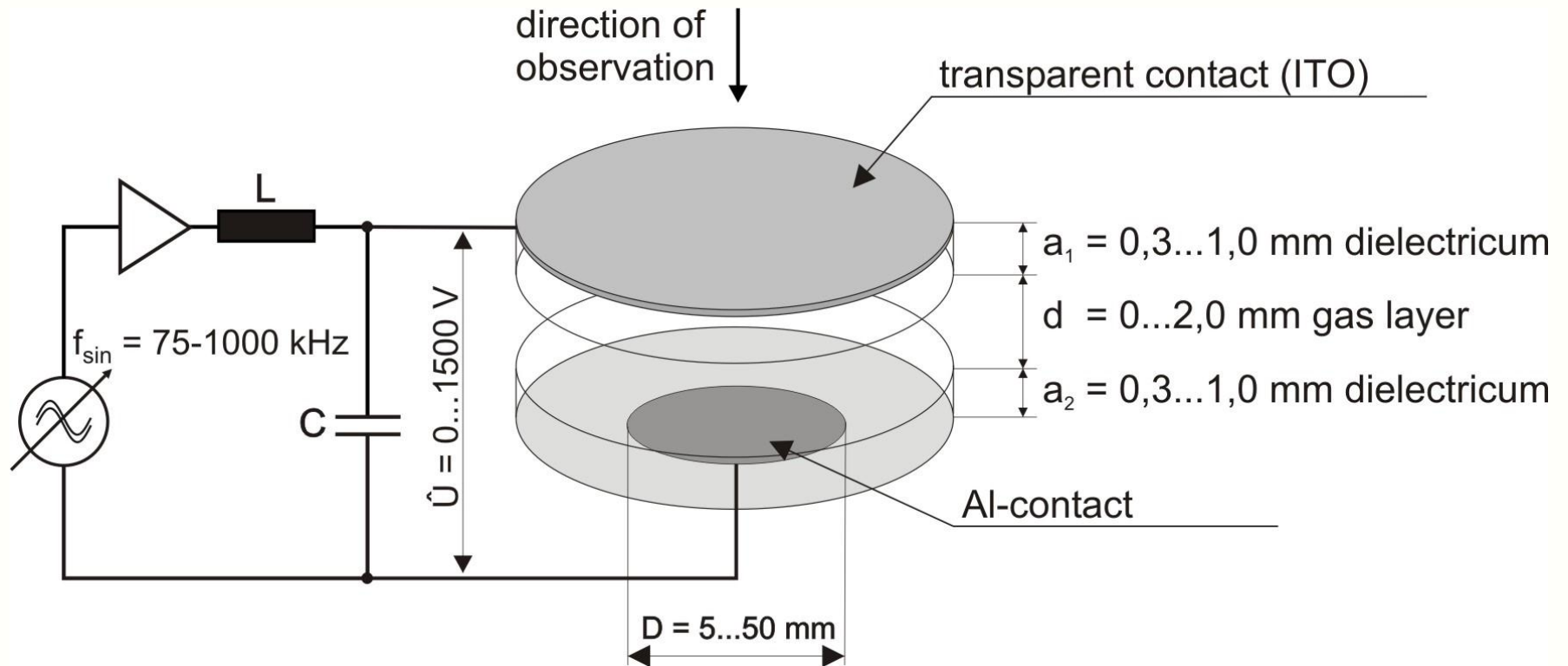
increasing conductivity of the semiconductor wafer

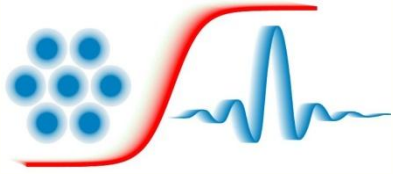


6. Experimental Results on Dissipative Soliton (DS) in AC Gas-Discharge Systems



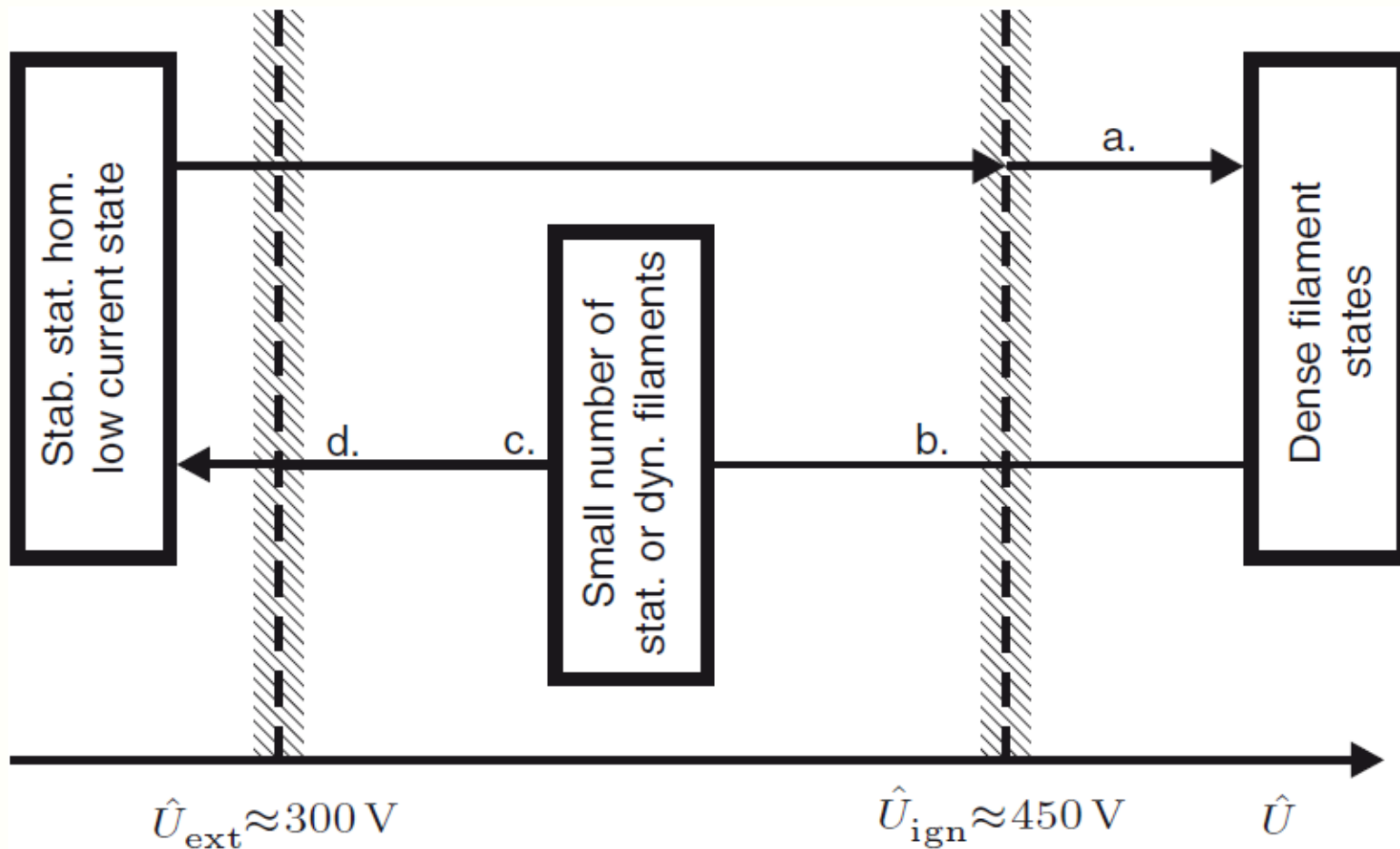
Experimental Quasi 2-Dimensional AC Gas-Discharge System: Experimental Set-Up

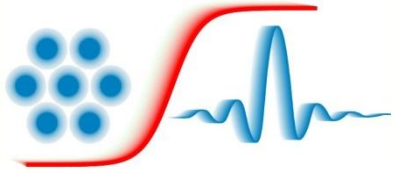




Experimental Quasi 2-Dimensional AC Gas-Discharge System: Bifurcation Sequence of Current Filament Patterns I

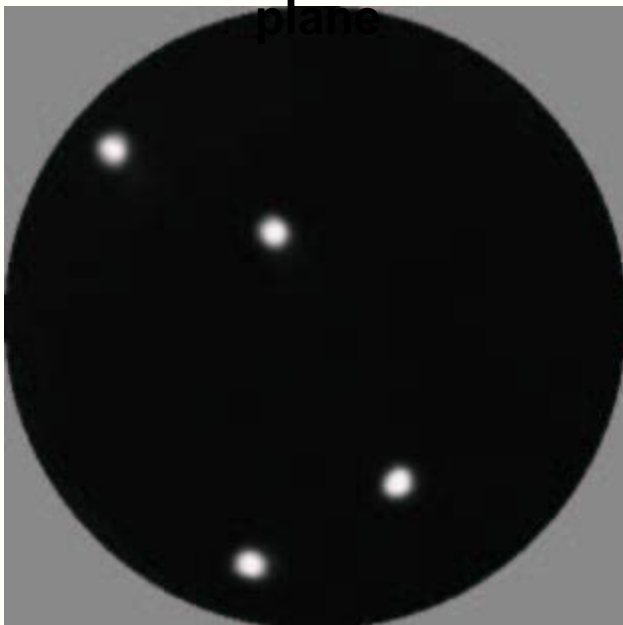
schematic formation of patterns for increasing and decreasing driver amplitude



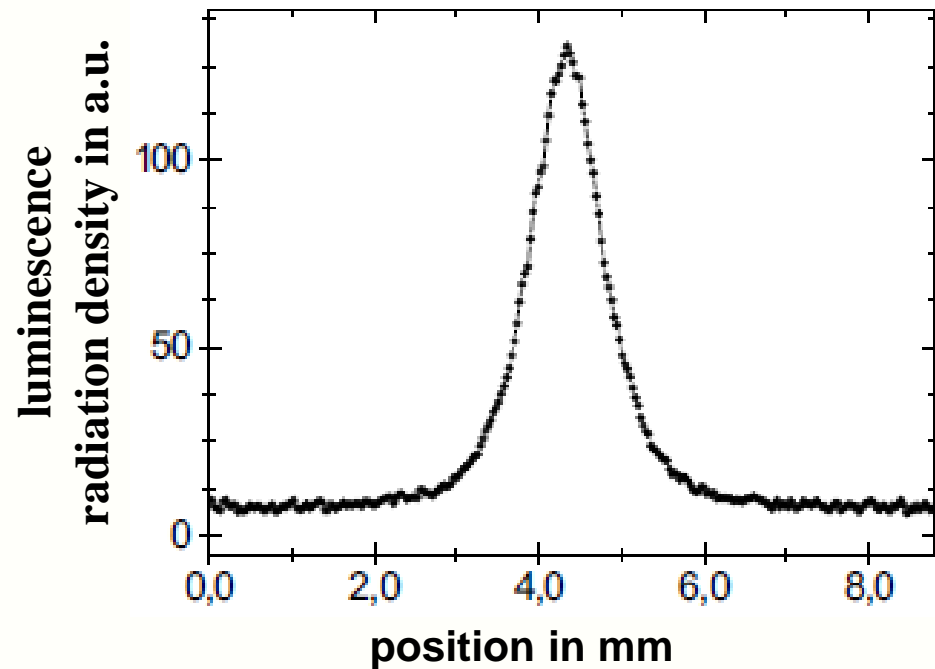


Experimental Quasi 2-Dimensional AC Gas-Discharge System: Isolated Stationary Current Filaments

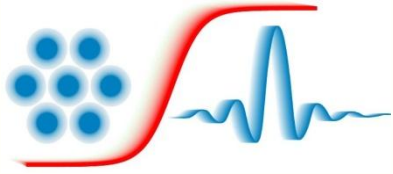
luminescence radiation
distribution in the discharge



cross section of a filament

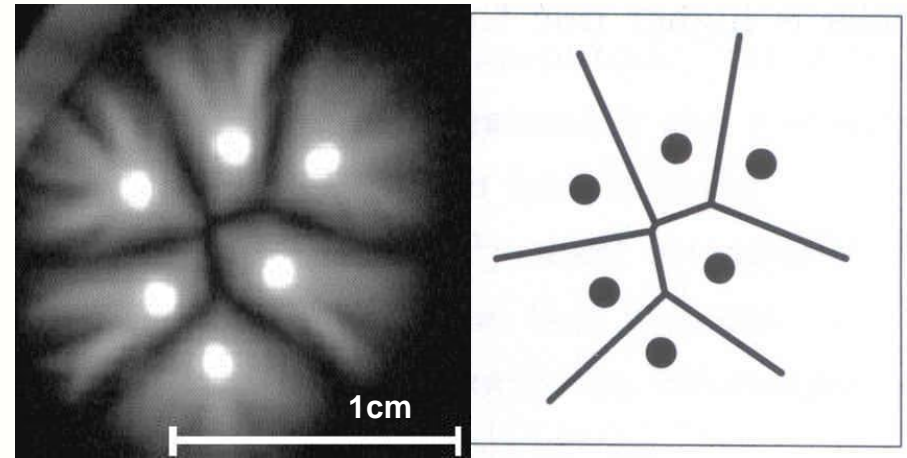
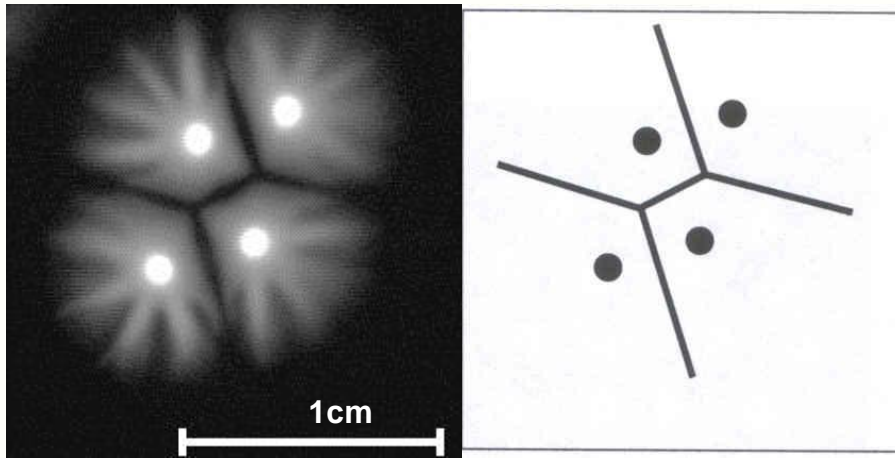


electrodes: a transparent ITO glass plate and a borosilicate glass plate (0.5 mm thick); room temperature, gas: He, $p=133\text{hPa}$, $d\sim 0.5\text{mm}$, $D\sim 20\text{mm}$, $f=200\text{kHz}$, $U=380\text{V}$, $t_{\text{exp}}=40\text{ms}$

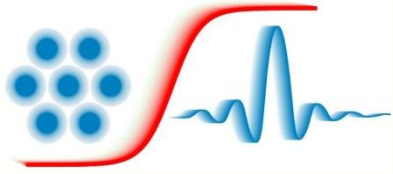


Experimental Quasi 2-Dimensional AC Gas-Discharge System: Self-Organized Voronoi Diagrams I

luminescence radiation distribution in the discharge plane

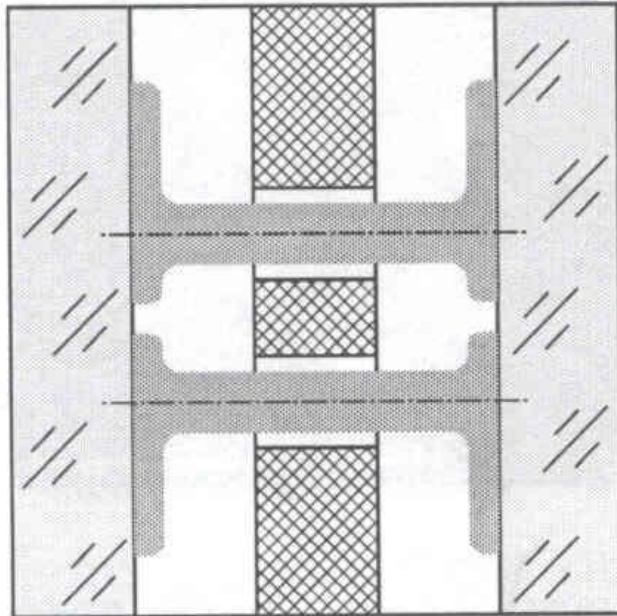


parameters: gas: N_2 , $p=122$ hPa, $d=2.6$ mm, $D=40$ mm, $a_1=0.5$ mm, $a_2=1$ mm,
 $\hat{U}=2050$ V, $t_{\text{exp}}=40$ ms, $f=50$ kHz



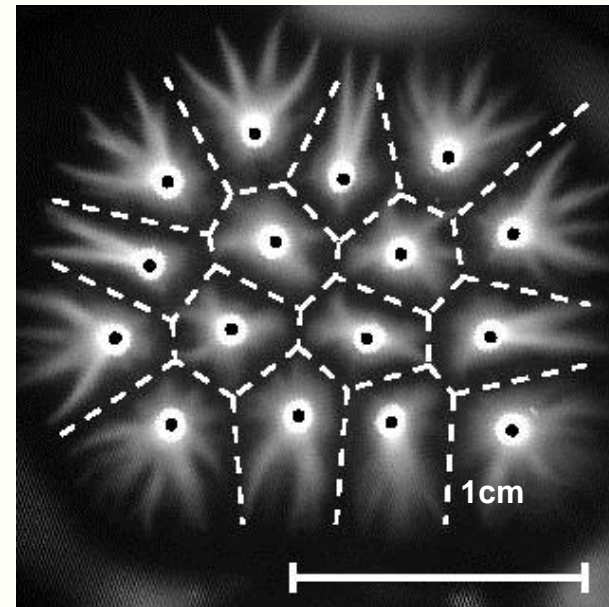
Experimental Quasi 2-Dimensional AC Gas-Discharge System: Self-Organized Voronoi Diagrams II

cross section of the
set-up



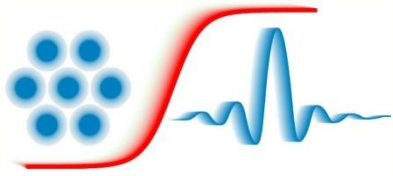
the positions of filaments are
predefining

luminescence radiation distribution in
the discharge plane



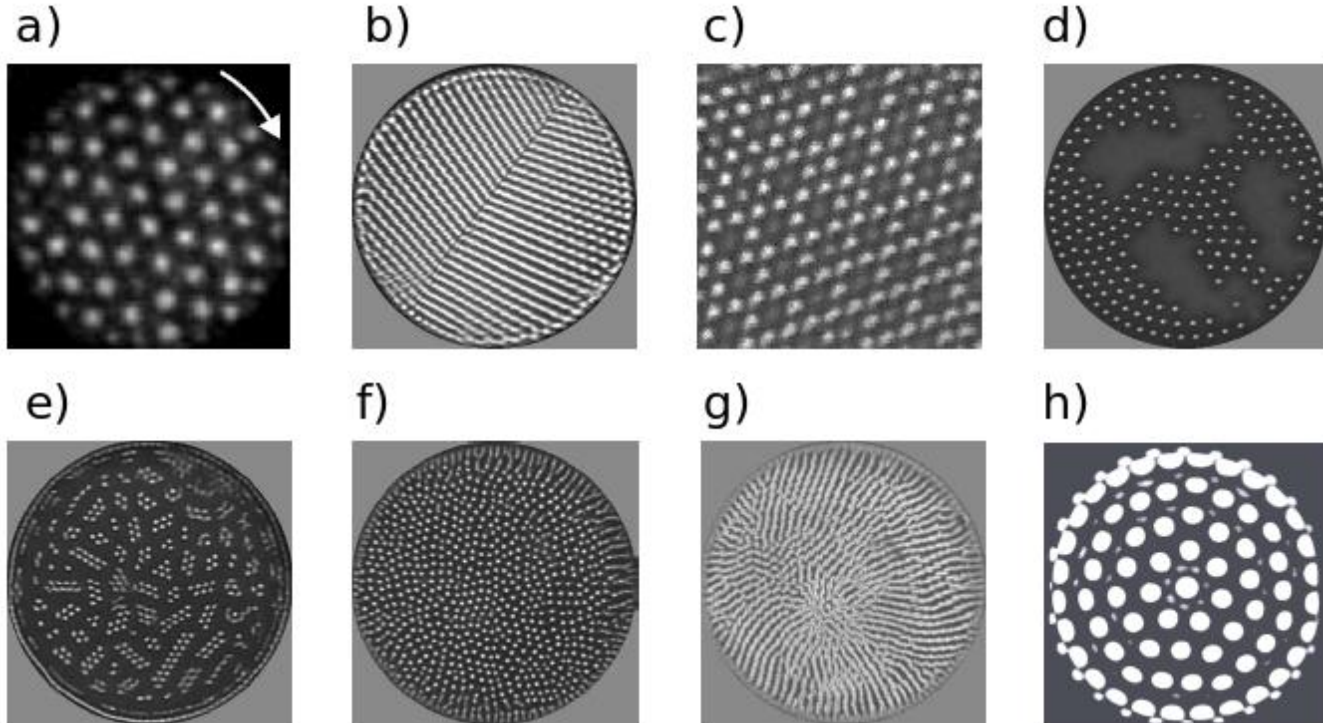
top view

parameters: gas: N_2 , $p=122$ hPa, $d=3.0$ mm (1mm inner isolating layer),
 $D=40$ mm, $a_1=0.7$ mm, $a_2=0.1$ mm, $\hat{U}=2100$ V, $t_{\text{exp}}=40$ ms, $f=50$ kHz

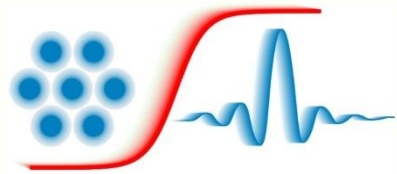


Experimental Quasi 2-Dimensional AC Gas-Discharge Systems: Selection of Snap- Shots of Dynamic Many-Filament Patterns

luminescence
radiation
distribution in
the discharge
plane



(a) rotating hexagonal pattern with “point-defects”; (b) travelling hexagonal pattern with a “grain boundary”; (c) superlattice with hexagonal symmetry; (d) pattern consisting of domains made of filaments; (e) pattern in the form of dynamical clusters (“molecular gas”); (f, g) generation of filaments at the boundary and annihilation while travelling to the centre, short (f) and long (g) exposure time; (h) filaments travelling on concentric rings with constant angular velocity, neighbouring rings may rotate in opposite direction

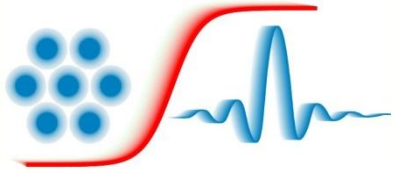


Experimental Quasi 2-Dimensional AC Gas-Discharge Systems: Travelling Filaments and Annihilation (Movie)

**luminescence
radiation
distribution in the
discharge plane**

**if not linked:
start movie
“Spots and their
Interaction.wmv”
In the folder**

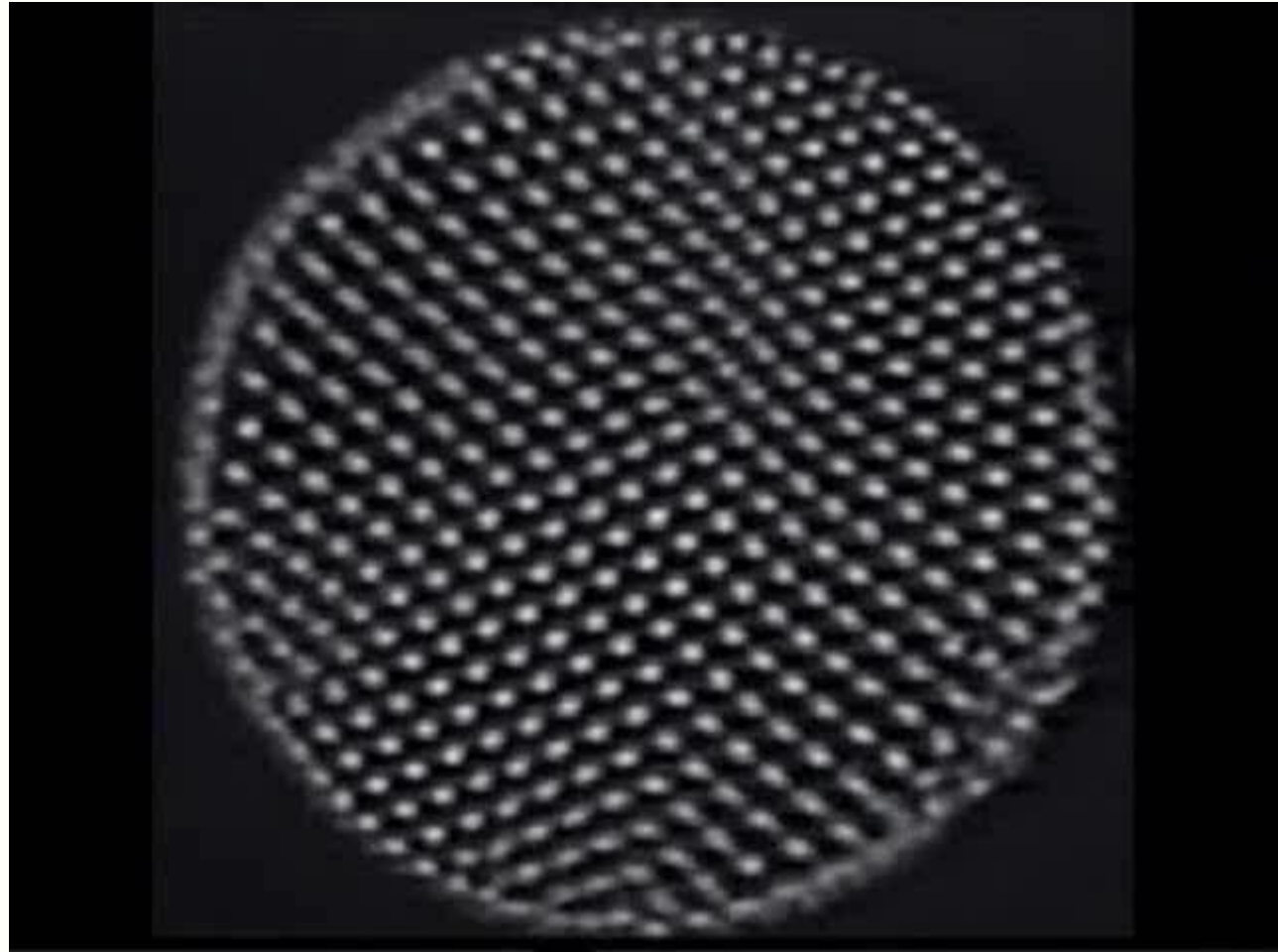


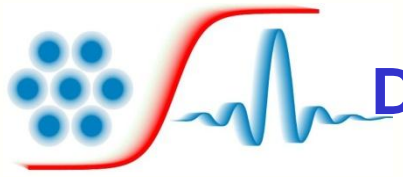


Experimental Quasi 2-Dimensional AC Gas-Discharge Systems: Dynamic Hexagonal Filament Pattern (Movie)

**luminescence
radiation
distribution in the
discharge plane**

**if not linked:
start movie
“Interaction of
Filaments at
High
Density.wmv”
in the folder**

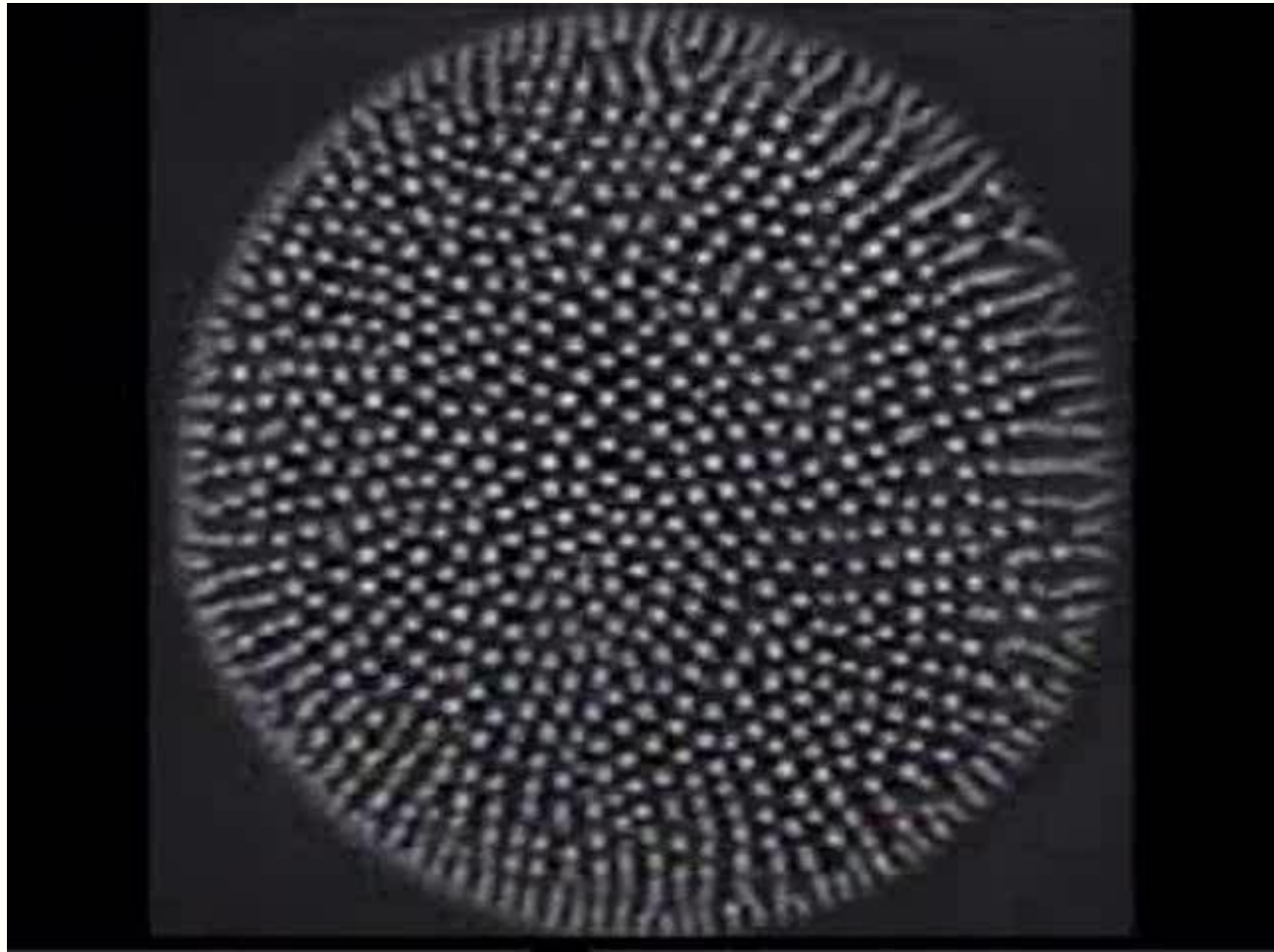


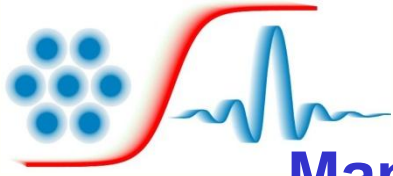


Experimental Quasi 2-Dimensional AC Gas-Discharge Systems: Many-Filament Pattern with Generation and Annihilation (Movie)

**luminescence
radiation
distribution in the
discharge plane**

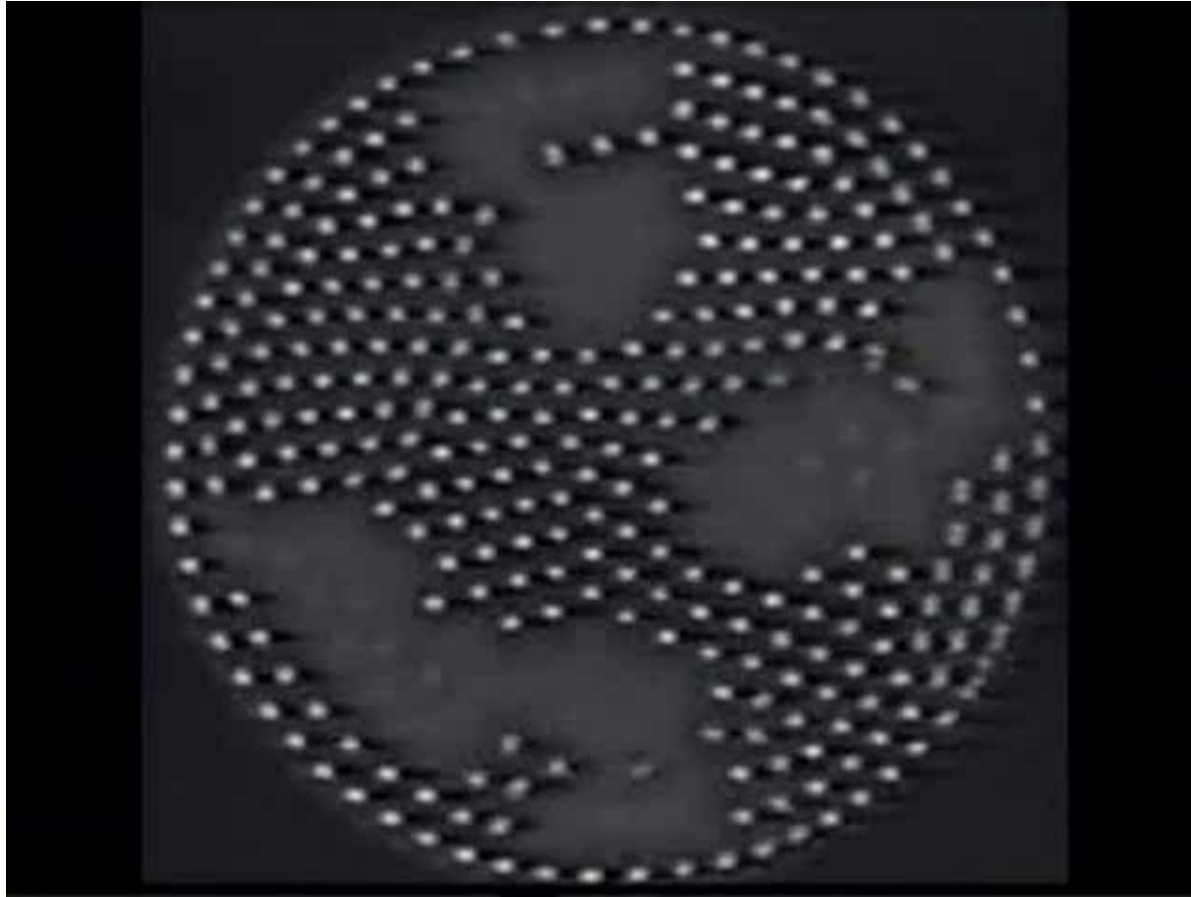
**if not linked: start
movie
“Generation and
Annihilation.wmv”
in the folder**



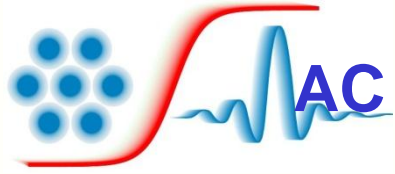


Experimental Quasi 2-Dimensional AC Gas-Discharge Systems: Dynamic Many-Filament Pattern with Domain Structure (Movie)

luminescence
radiation
distribution in the
discharge plane



start movie
from folder
“Domains of
High Density
Of Filament
Patterns.wmv”



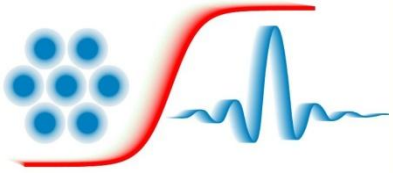
Experimental Quasi 2-Dimensional AC Gas-Discharge Systems: Dynamic Many-Filament Pattern Resembling a “Molecule Gas” (Movie)

**luminescence
radiation
distribution in the
discharge plane**

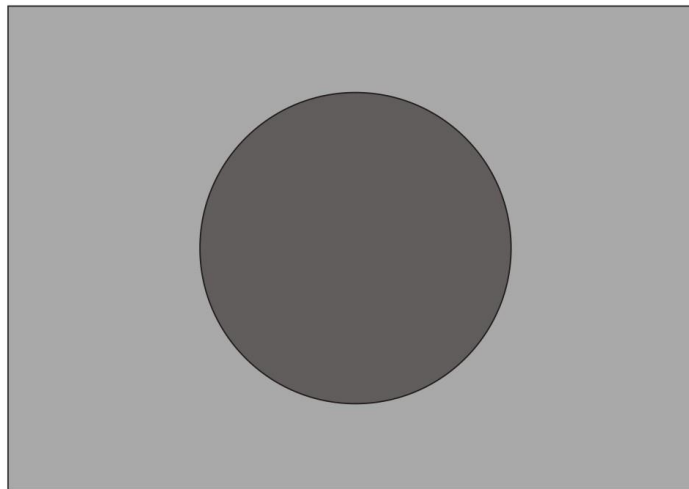
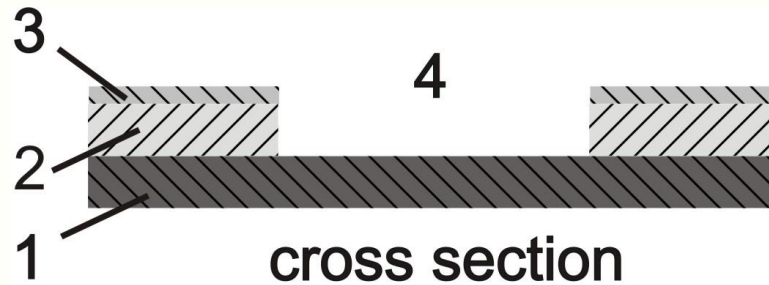
**if not linked: start
movie**

**“Interaction of
Clusters of High
Density.wmv”
in the folder**

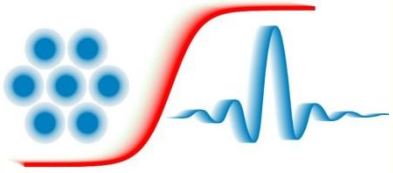




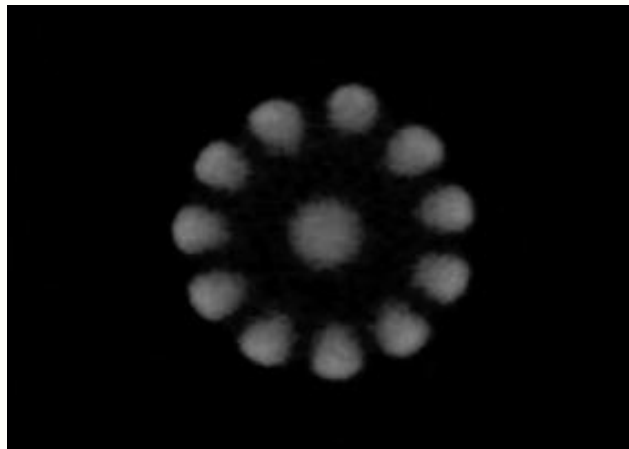
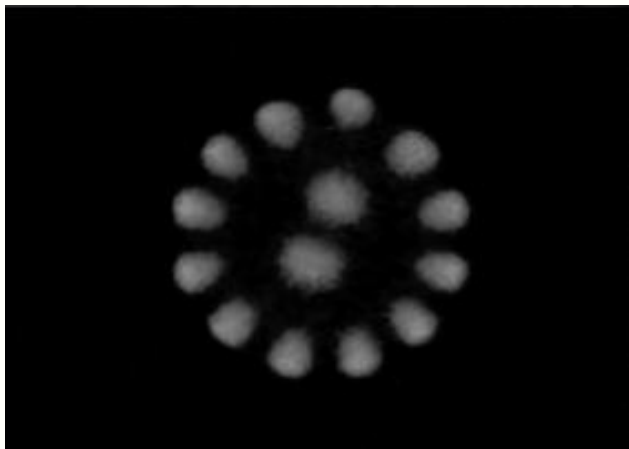
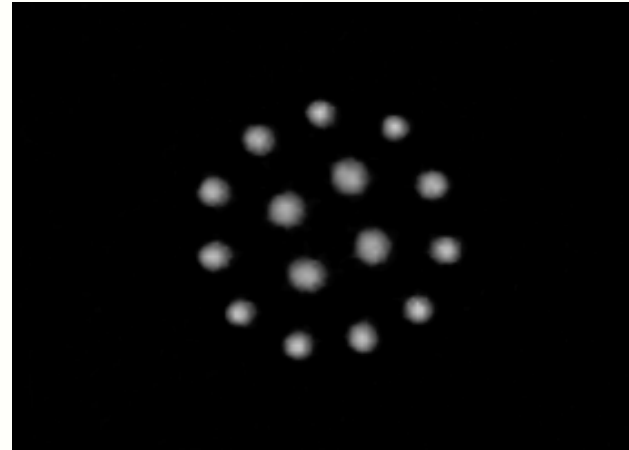
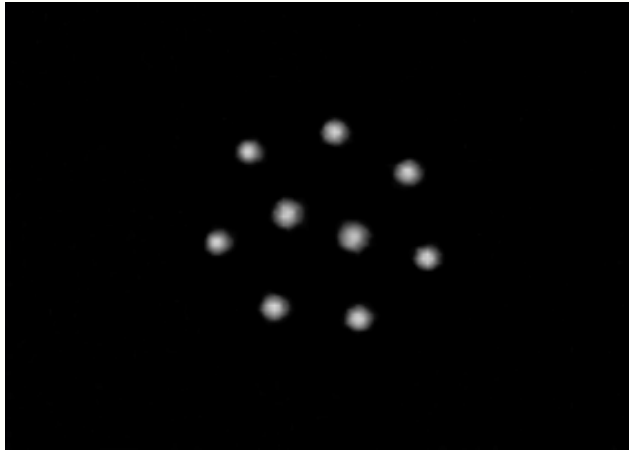
Cathode Spots in Gas-Discharge Systems I: Experimental Set-Up

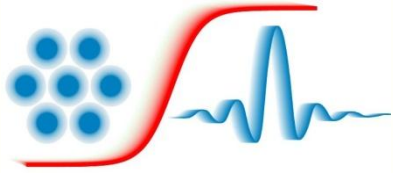


- 1. cathode**
- 2. dielectric spacer**
- 3. hollow anode**
- 4. gas space**

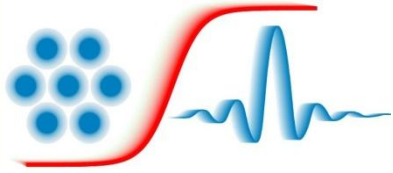


Cathode Spots in Gas-Discharge Systems II: Observed Spot Patterns

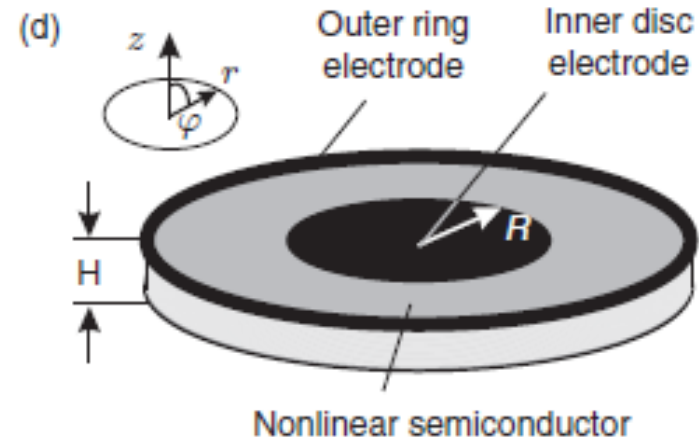
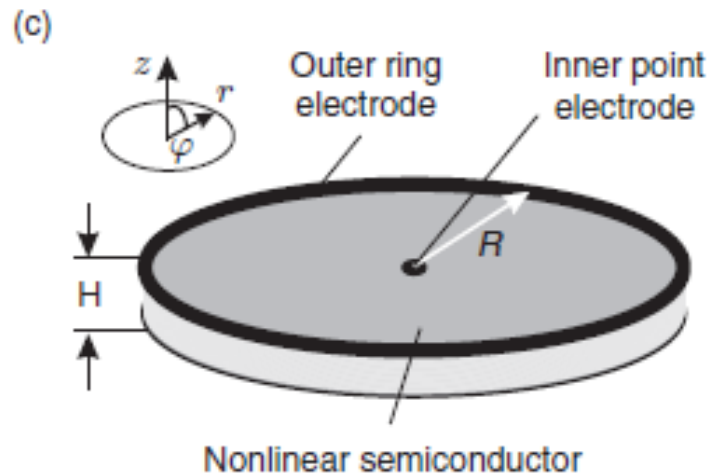
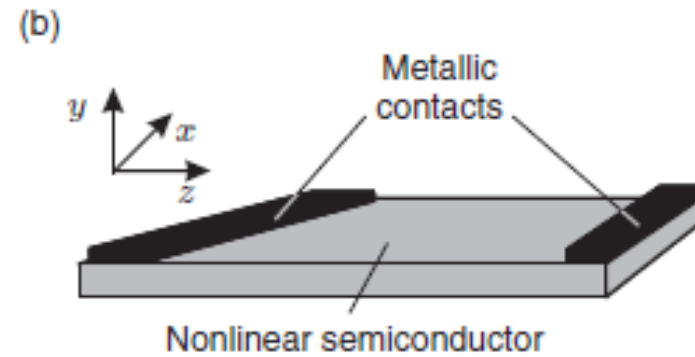
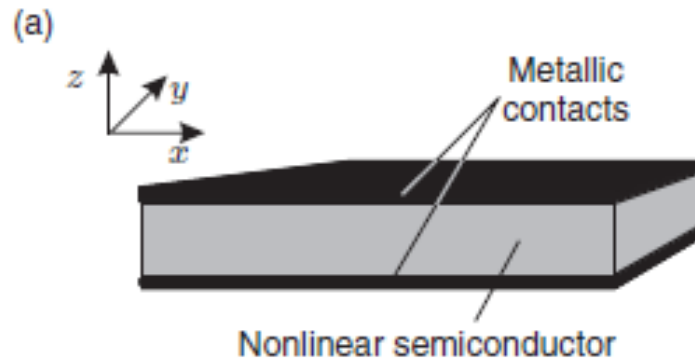




7. Experimental Results on Dissipative Solitons (DS) in Semiconductors



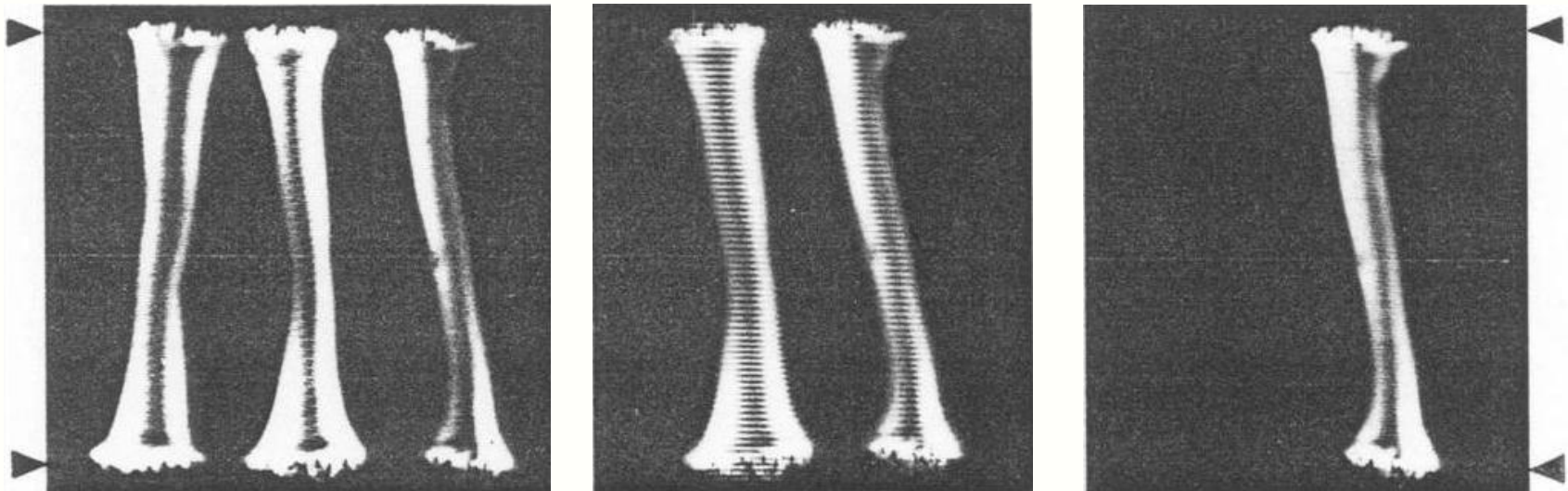
Current Filaments in Semiconductors I: Experimentally Investigated Configurations



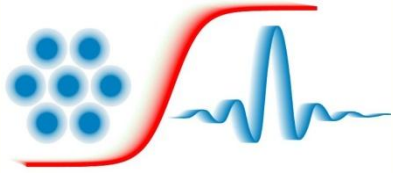


Current Filaments in Semiconductors II: Bifurcation Cascade of Stationary Current Filaments for Decreasing Driving Voltage

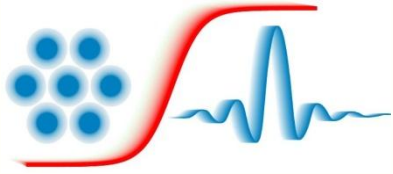
electron microscope image of a film of n-GaAs at $T = 4.2$ K



for Decreasing Driving Voltage in a Film

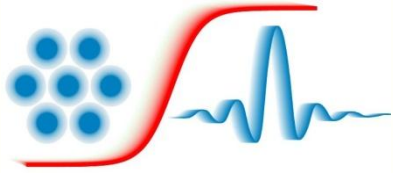


8. Quantitative Modelling of AC Gas-Discharge by Drift-Diffusion Equations



Typical Time Scales for Low Temperature Gas Discharge Systems (He, $p \approx 100\text{hPa}$, $T_e = 1\text{eV}$) I

time scale	value	comments
times for electrons		
• free path times	10^{-12}s	average time between collisions
• drift time	10^{-9}s	travel time between electrodes
• diffusion time	10^{-7}s	diffusion over electrode gap
times for ions		
• free path times	10^{-9}s	average time between collisions
• drift time	10^{-7}s	travel time between electrodes
• diffusion time	10^{-5}s	diffusion over electrode gap
times in experiment		
• ac driver period	$\approx 10^{-5}\text{s}$	
• transient time	$\lesssim 10^{-1}\text{s}$	



Drift-Diffusion Model for the Planar DC Semiconductor Gas-Discharge System I: Equations for the Gas

in the gas:

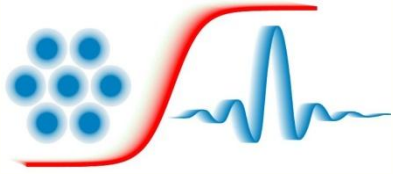
$$\partial_t \mathbf{n}_{e,i} = -\operatorname{div} \vec{\Gamma}_{e,i} + S_{e,i},$$

$$\vec{\Gamma}_{e,i} = \mp \mu_{e,i}(\mathbf{E}) \mathbf{n}_{e,i} \vec{\mathbf{E}} - \mathbf{D}_{e,i} \nabla \mathbf{n}_{e,i},$$

$$\Delta \varphi = (\mathbf{n}_i - \mathbf{n}_e)(|e| / \varepsilon_0 \varepsilon),$$

$$\vec{\mathbf{E}} = \nabla \varphi$$

$$S_e = S_i = \alpha(\mathbf{E})(\vec{\Gamma}_e / e)$$



Drift-Diffusion Model for the Planar DC Semiconductor Gas-Discharge System II: Boundary Gas – Semiconductor at $z=0$

$$\partial_t \sigma - D_S \Delta \sigma = \vec{e}_z (\vec{j}_g - \vec{j}_{SC})_{z=0},$$

$$\frac{1}{\epsilon_0} \sigma = (\epsilon \vec{e}_z \vec{E})_{z=+0} - (\vec{e}_z \vec{E})_{z=-0},$$

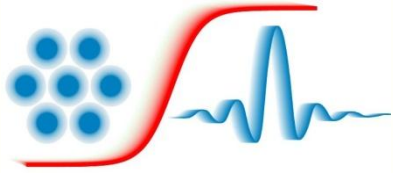
$$(\vec{\Gamma}_p \vec{e}_z)_{z=-d} = (\mu_p n_p \vec{E} \vec{e}_z + \frac{1}{4} n_p \langle v_p \rangle)_{z=-0},$$

$$(\vec{\Gamma}_e \vec{e}_z)_{z=-0} = (-\mu_e n_e \vec{E} \vec{e}_z + \frac{1}{4} n_e \langle v_e \rangle - \gamma \vec{\Gamma}_p \vec{e}_z)_{z=-0},$$

$$\langle v_{e,p} \rangle = \sqrt{8k T_{e,p} / \pi m_{e,p}}.$$

$$(\varphi)_{z=-d} - (\varphi)_{z=d_{SC}} = U.$$

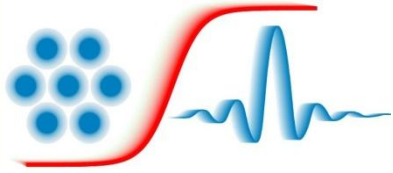
- σ surface charge, dependence on surface properties
- D_S diffusion constant of surface charge
- \vec{e}_z unity vector in z-direction
- ϵ dielectric constant of the semiconductor
- v_e, v_p thermal electron/ion speed
- γ γ -Townsend-coefficient
- k Boltzmann-constant
- m_e, m_p electron/ion mass
- U voltage drop at the component
- $z=0$ gas/semiconductor surface
- $z=-d$ metallic anode



Drift-Diffusion Model for the Planar DC Semiconductor Gas-Discharge System III: Semiconductor Wafer ($0 < z < d_{sc}$)

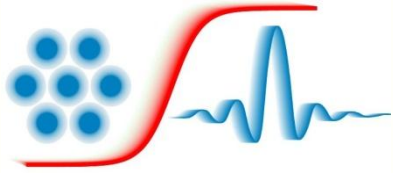
$$\begin{aligned}\vec{j}_{sc} &= \lambda \vec{E}, \\ \vec{E} &= -\vec{\nabla} \varphi, \\ \text{div}(\varepsilon \vec{\nabla} \varphi) &= 0.\end{aligned}$$

\vec{j}_{sc}	global electrical current density
λ	specific electrical conductivity
\vec{E}	electrical field
φ	electrical potential
ε	dielectric constant of the semiconductor



Drift-Diffusion Model for the Planar AC Semiconductor Gas-Discharge System I: Equations

equations similar to those of the DC system



Numerical Solutions of the Drift-Diffusion Equation for the Quasi-2-Dimensional AC Gas-Discharge System Hexagonal Pattern II: Parameters

experimental (theoretical) parameters

gas: He

$p = 300\text{hPa}$ (300)

$a = 0.5\text{mm}$ (0.5)

$d = 0.5\text{mm}$ (0.5)

$D = 8\text{mm}$ (8)

$f_{\sin} = 200\text{kHz}$ (200)

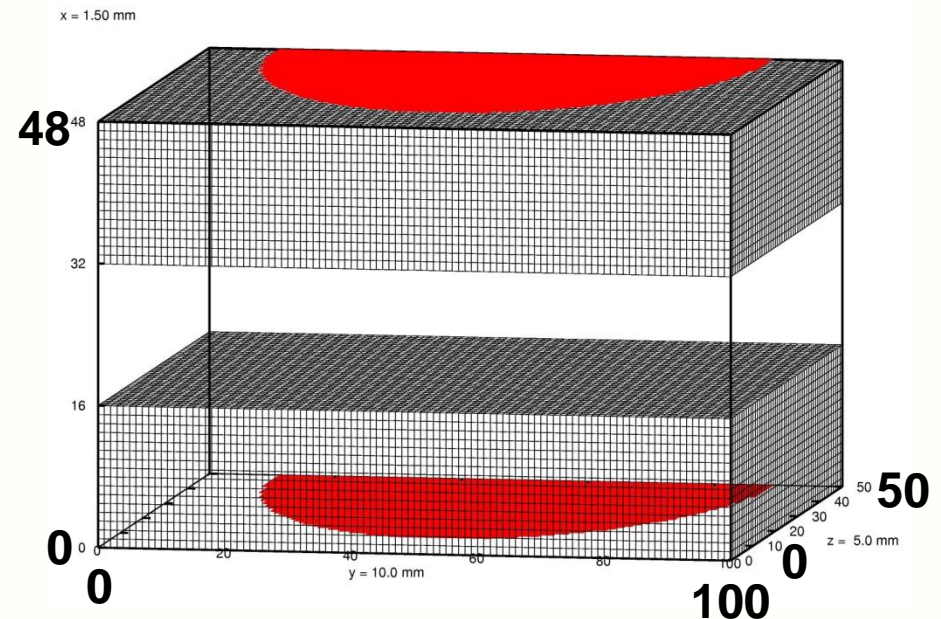
$\hat{U} = 675\text{V}$ (700)

$t_{\text{exp}} = 2\mu\text{s}$

($\gamma = 0,05$)

(α, μ from tables)

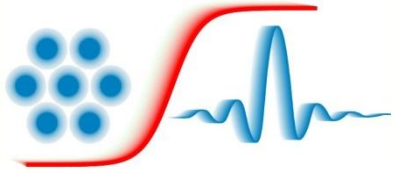
Numerical parameters



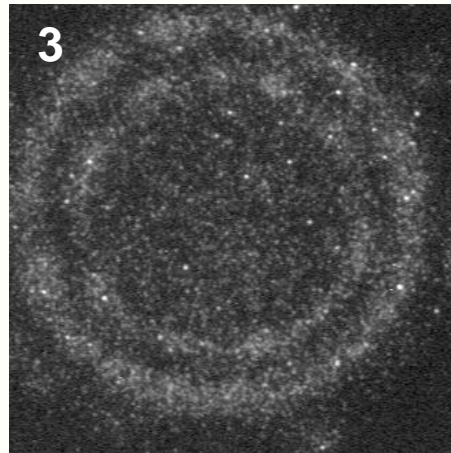
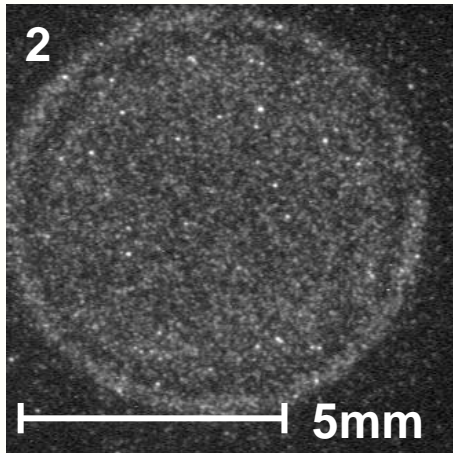
grid points 100 x 50 x 48

10mm x 5mm x 1,5mm

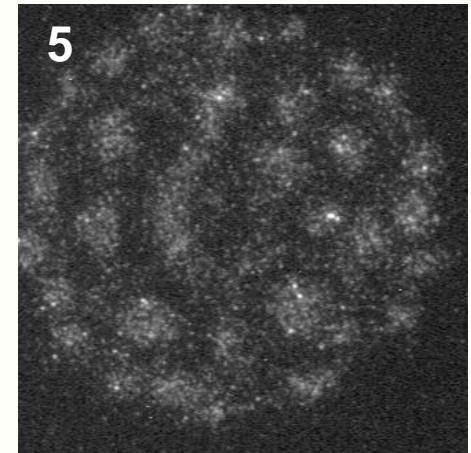
Neumann boundary conditions



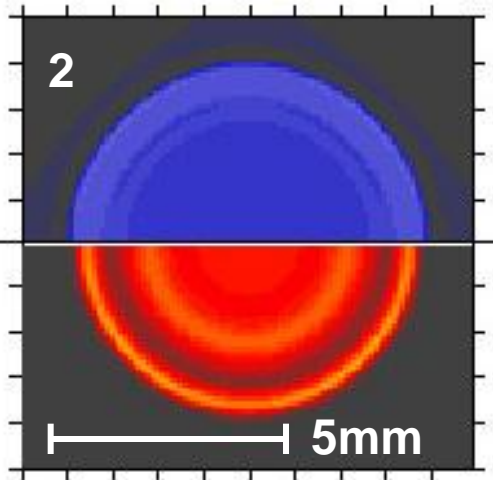
Numerical Solutions of the Drift-Diffusion Equation for the Quasi-2-Dimensional AC Gas-Discharge System Hexagonal Pattern III: Experiment versus Theory (a)



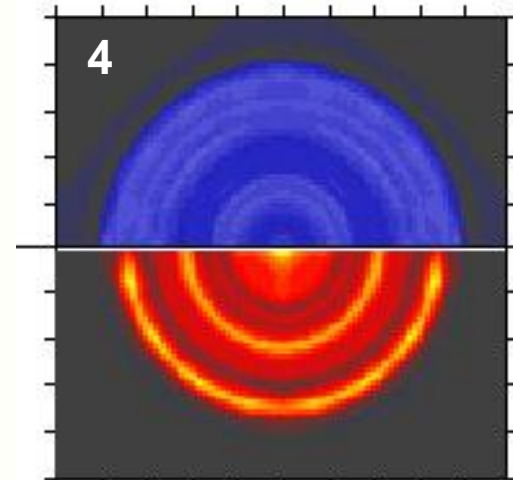
experiment

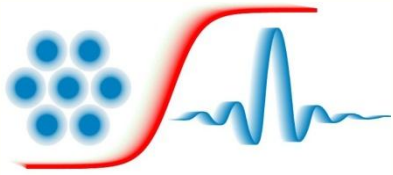


———— increasing number of breakdowns ———>



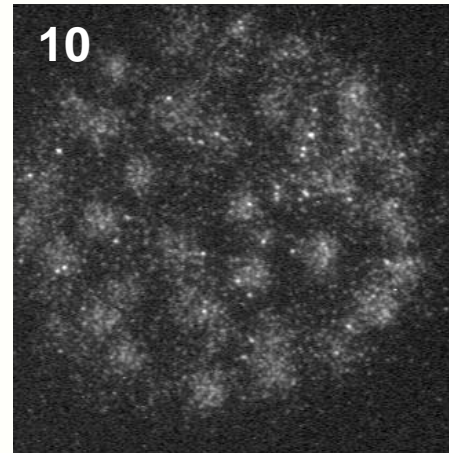
theory





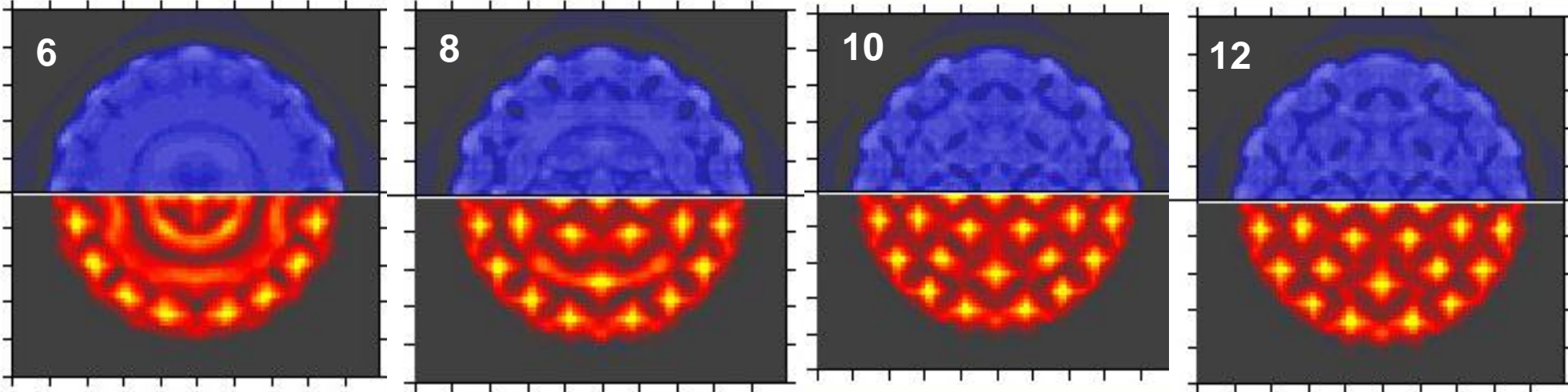
Numerical Solutions of the Drift-Diffusion Equation for the Quasi-2-Dimensional AC Gas-Discharge System Hexagonal Pattern IV: Experiment versus Theory (b)

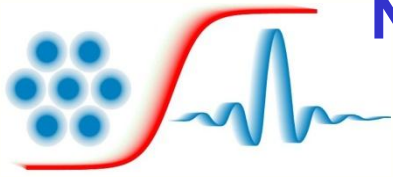
experiment →



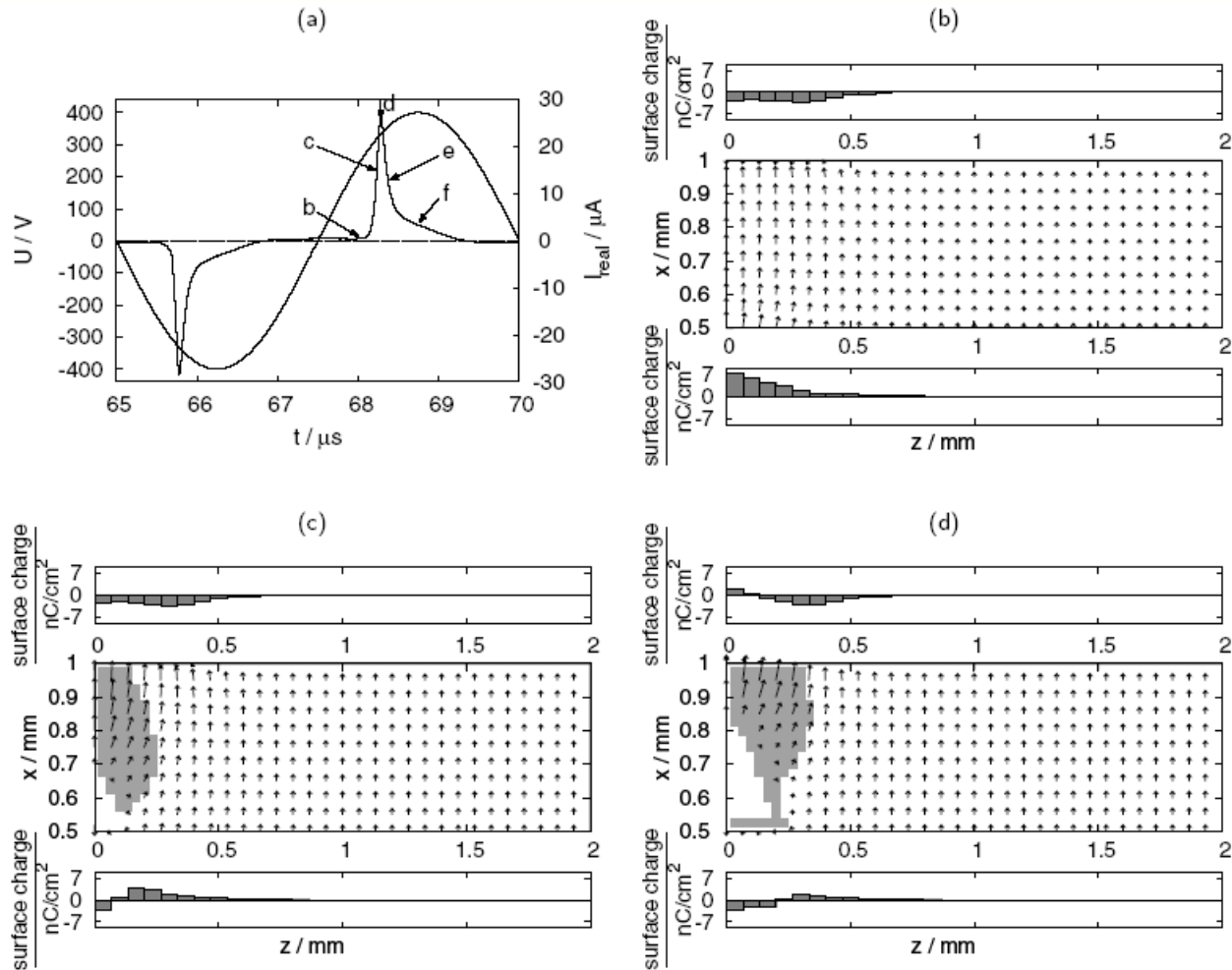
theory ↘

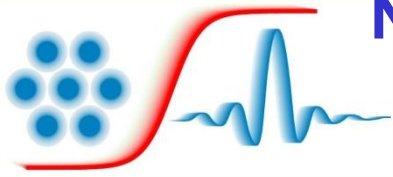
————— increasing number of breakdowns —————→



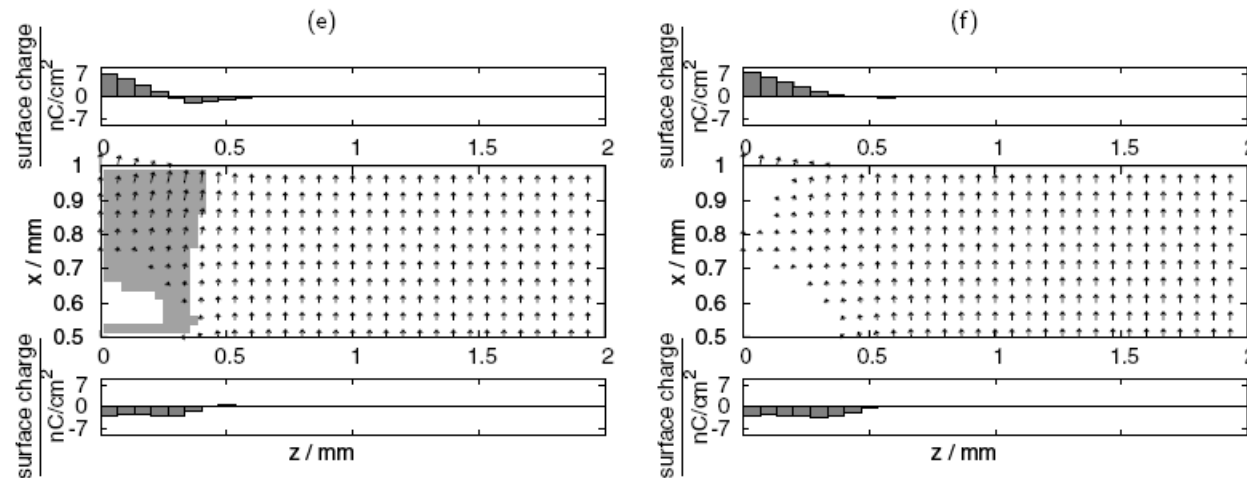


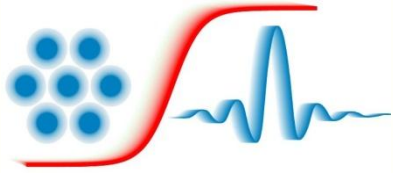
Numerical Solutions of the Drift-Diffusion Equation for the Quasi-2-Dimensional AC Gas-Discharge System: Evolution of a Stationary Filament I



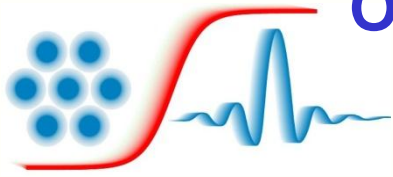


Numerical Solutions of the Drift-Diffusion Equation for the Quasi-2-Dimensional AC Gas-Discharge System: Evolution of a Stationary Filament II

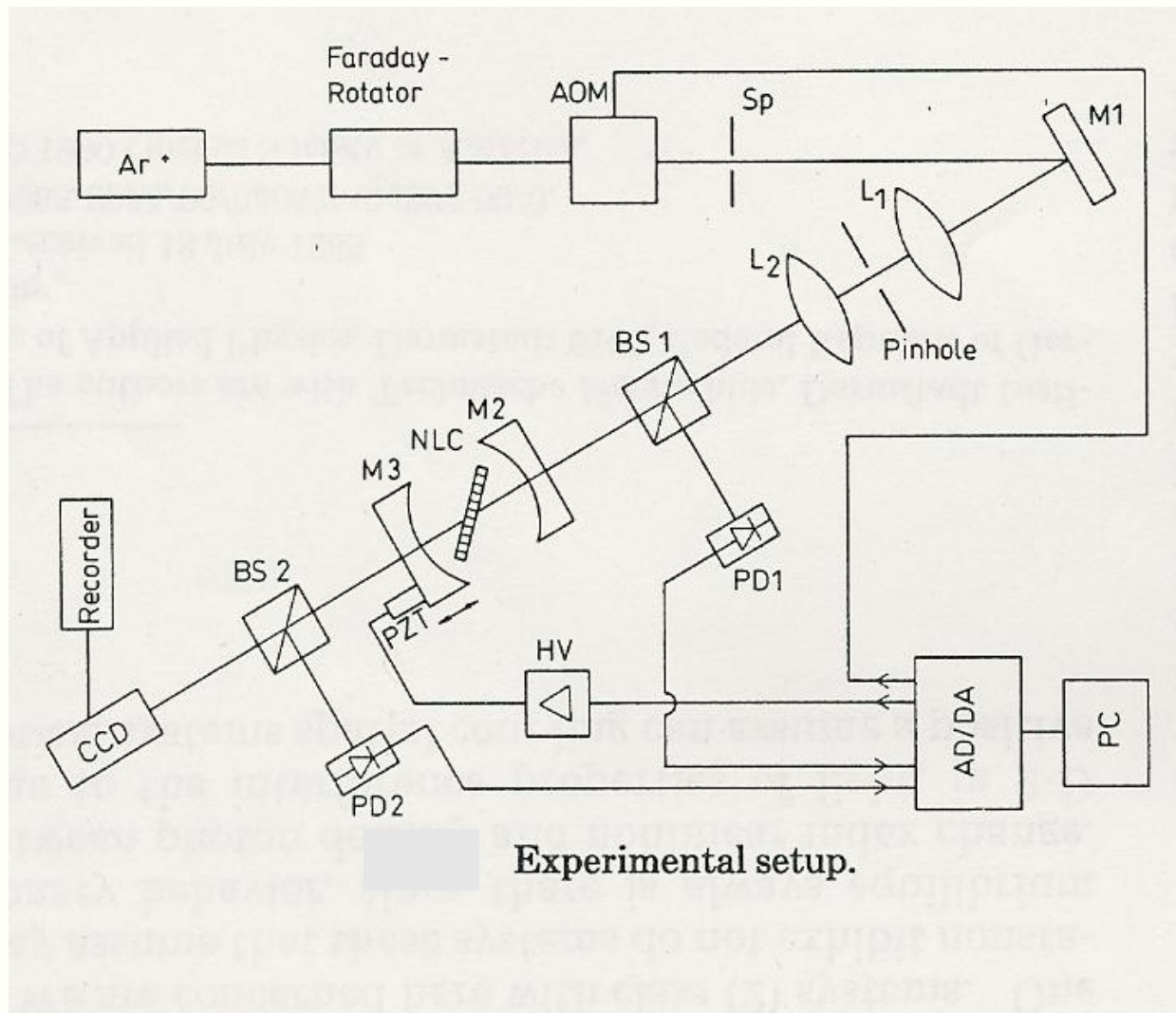


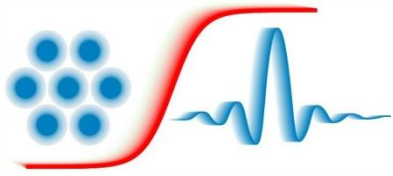


9. Various Dissipative Soliton Carrying Experimental Systems: Other than Electrical Transport Systems



Optical Filaments in a Fabry-Perot Resonator with a Liquid Crystal as Kerr Medium I: Experimental Set-Up

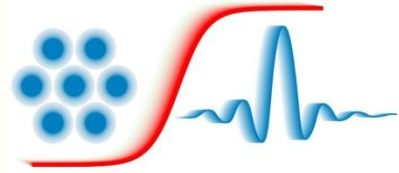




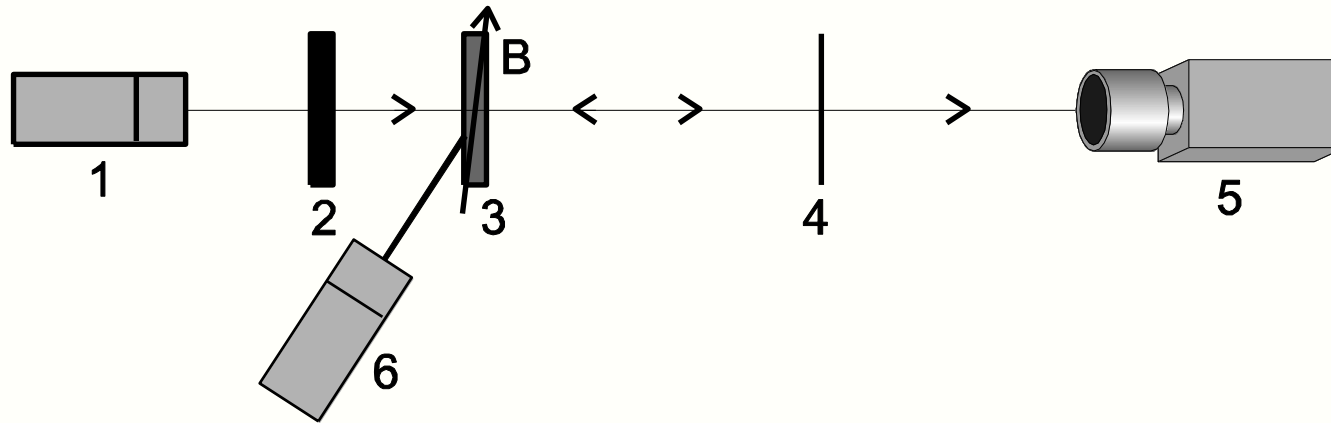
Optical Filaments in a Fabry-Perot Resonator with a Liquid Crystal as Kerr Medium II: Observation of Localized Structures



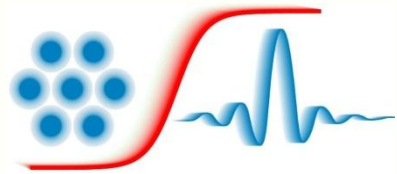
Sequence of patterns from a structure-forming bistable element.



Optical Filaments in a Single Mirror Feedback Device with Na Vapour as Nonlinear Medium I: Experimental Set-Up

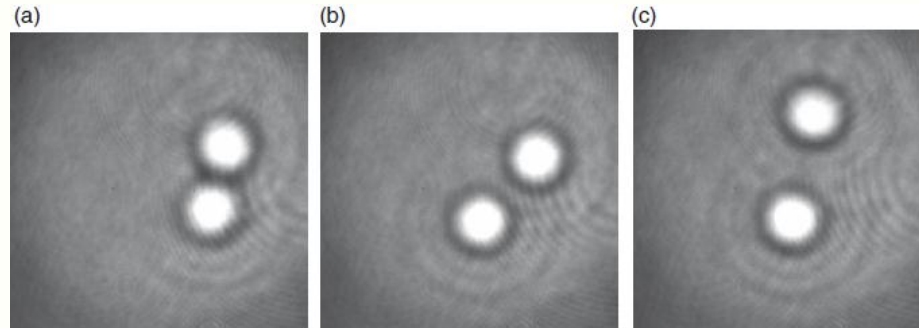


1. laser, 2. polarizing element, 3. cell with Na vapour being subject to a magnetic field, 4. semitransparent mirror, 5. camera, 6. laser for control beam

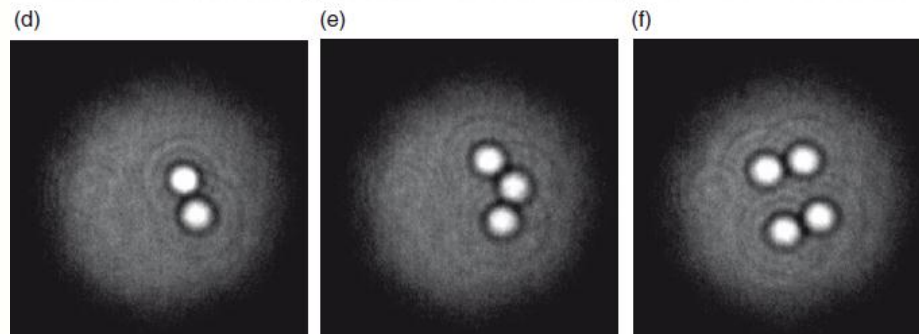


Optical Filaments in a Single Mirror Feedback Device with Na Vapour as Nonlinear Medium II: Multistability of Localized Structures

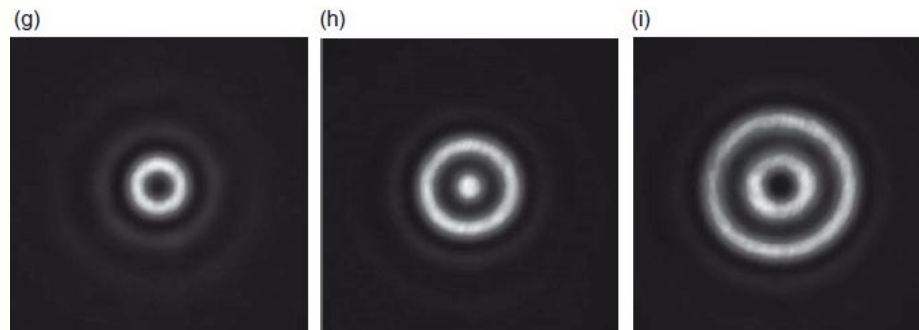
at a given set of parameters filaments may lock in at different distances forming clusters („molecules“)

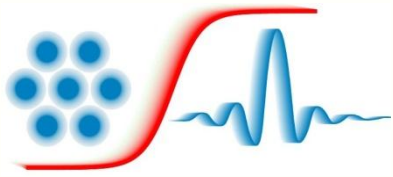


at a given set of parameters different number of can exist



at a given set of parameters different kinds of filament can exist



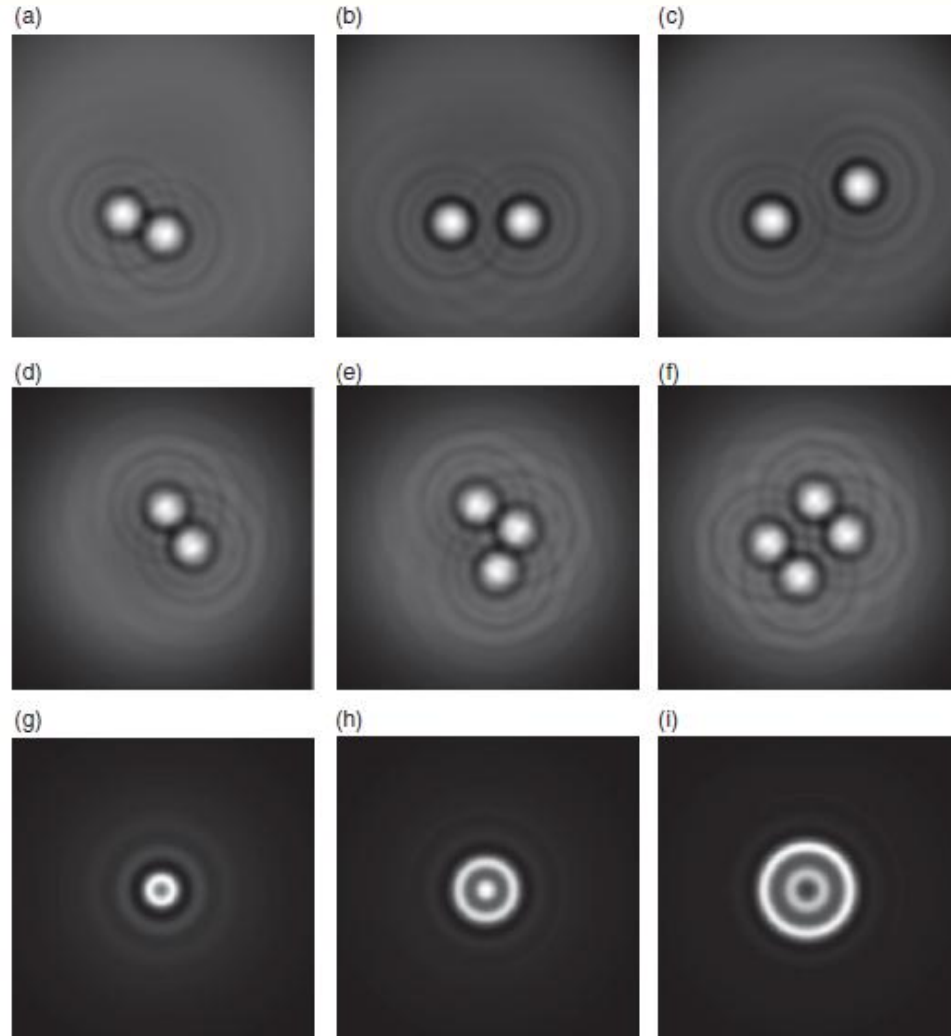


Optical Filaments in a Single Mirror Feedback Device with Na Vapour as Nonlinear Medium III: Numerical Solutions of the Model Equations

at a given set of parameters filaments may lock in at different distances forming clusters („molecules“)

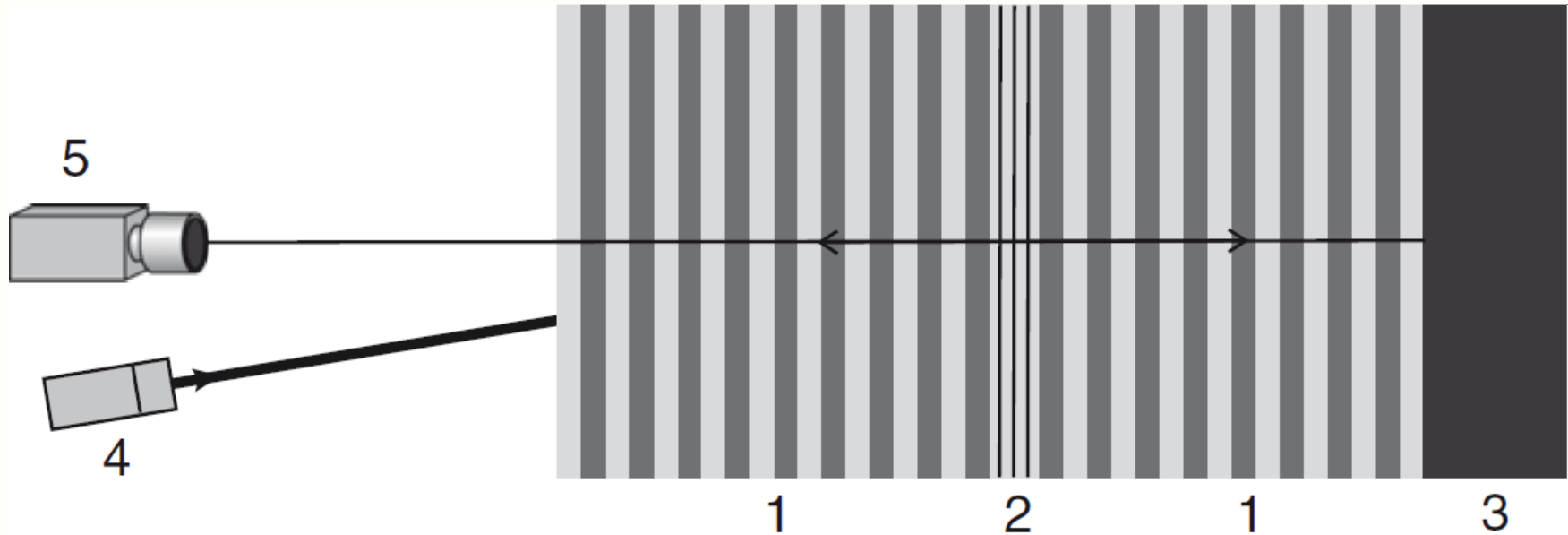
at a given set of parameters different number of can exist

at a given set of parameters different kinds of filament can exist

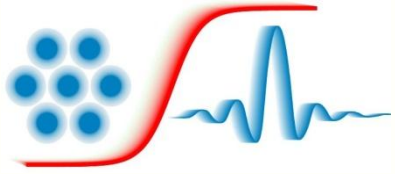




Optical Semiconductor Vertical Cavity Surface Emitting Lasers (VCSELs) I: Experimental Set-Up

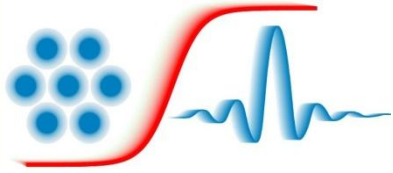


- 1 Bragg reflectors
- 2 layer of quantum wells
- 3 substrate
- 4 Laser
- 5 camera

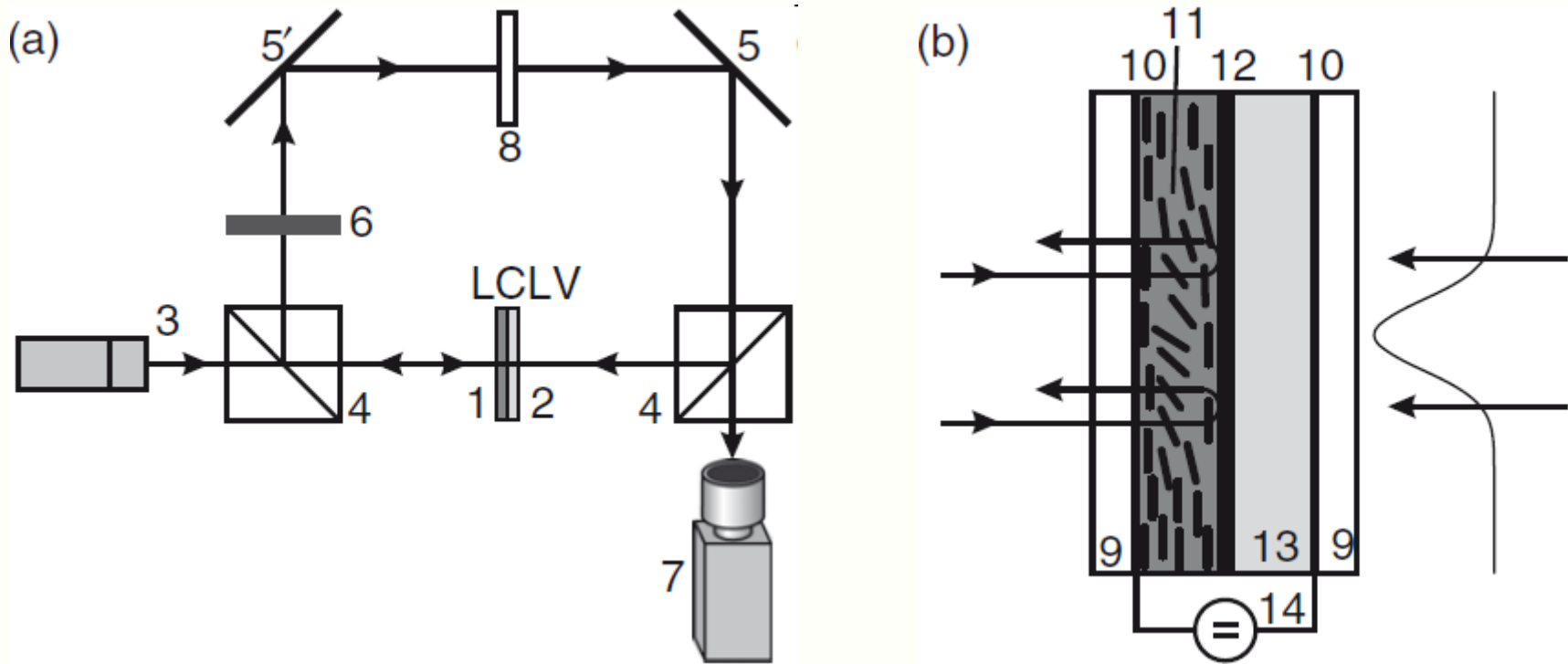


Optical Semiconductor Vertical Cavity Surface Emitting Lasers (VCSELs) II: Stationary Localized structures in the Plane





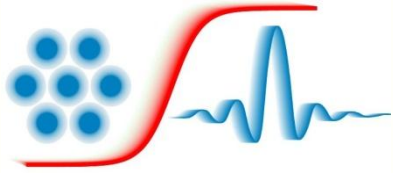
Optical Liquid Crystal Light Valves S(LCLVs) I: Experimental Set-Up



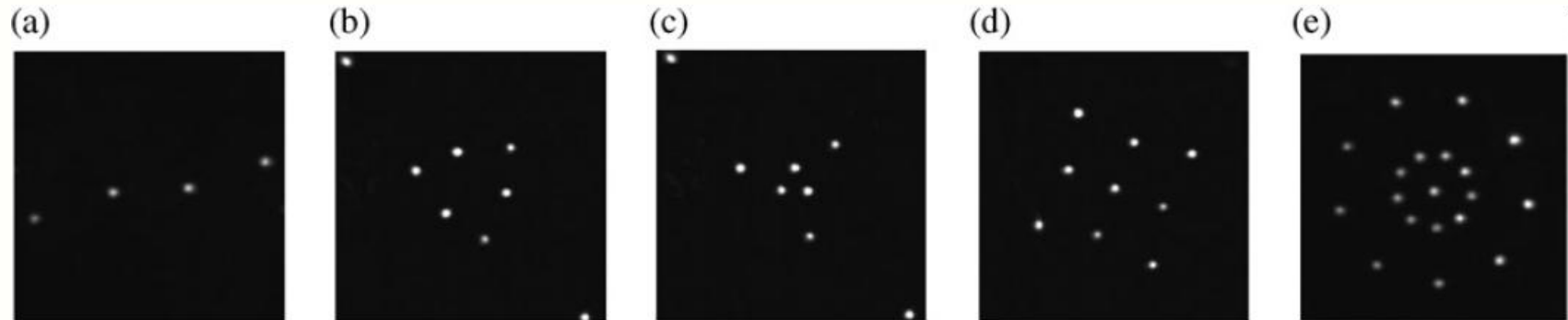
1 write side of the LCLV
2 read side of the LCLV
3 laser
4 beam splitter
5 mirror

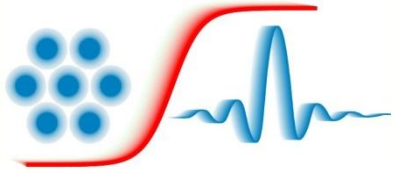
6 polarising elements
7 camera
8 additional element for
manipulation
9 glass plates

10 electrically cond. layer
11 liquid crystal layer
12 mirror
13 photocond. layer
14 voltage

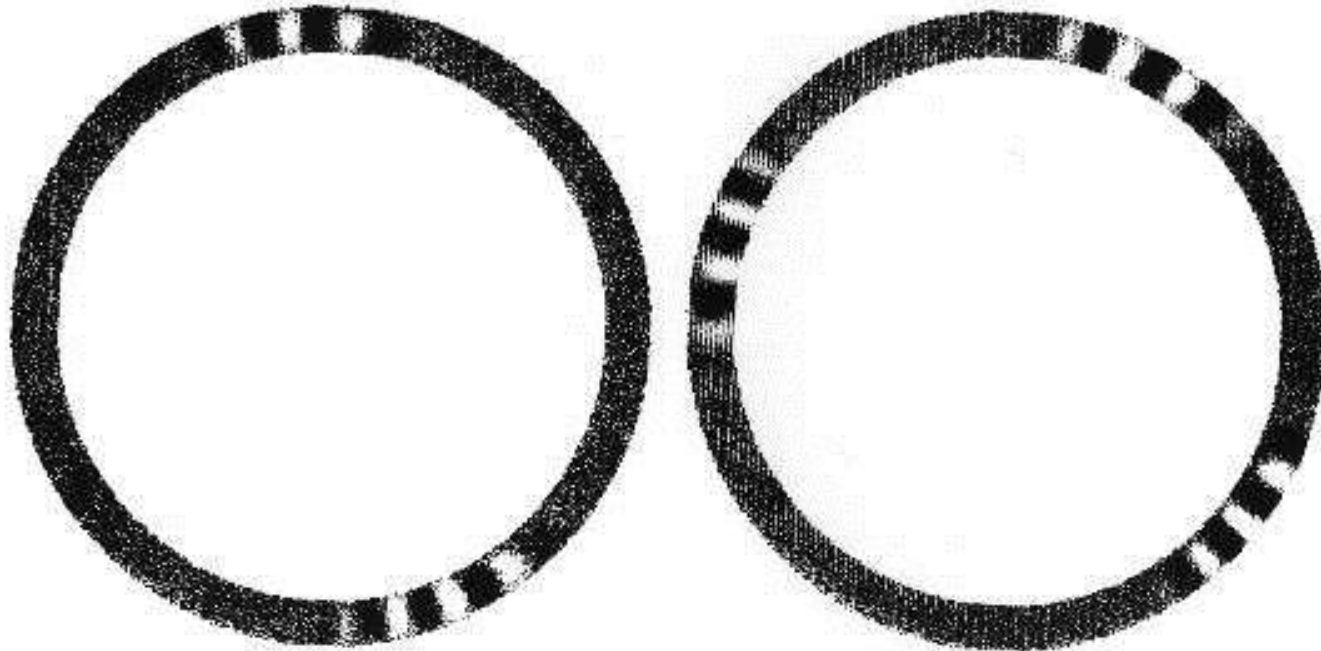


Optical Liquid Crystal Light Valves S(LCLVs) II: Stationary Localized Structures

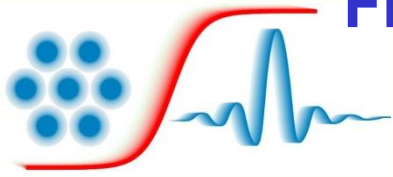




Intermediate Localization of Patterns in Hydrodynamics: Quasi 1-Dimensional Rayleigh-Benard Experiment

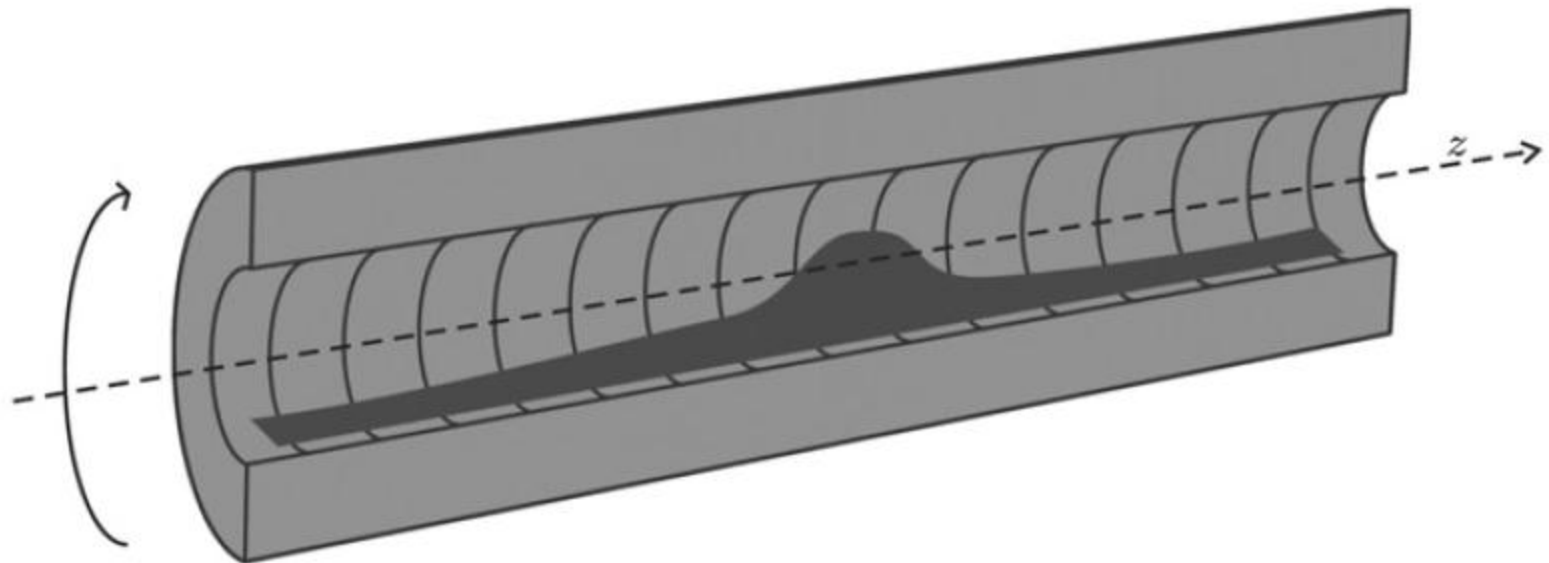


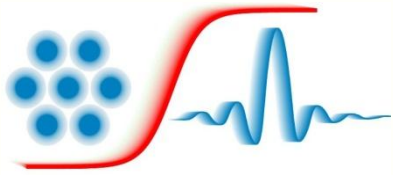
localized structures in the convection of an annulus filled with a mixture of ethanol and water and heated from below, the pulses are surrounded by non-convecting fluid



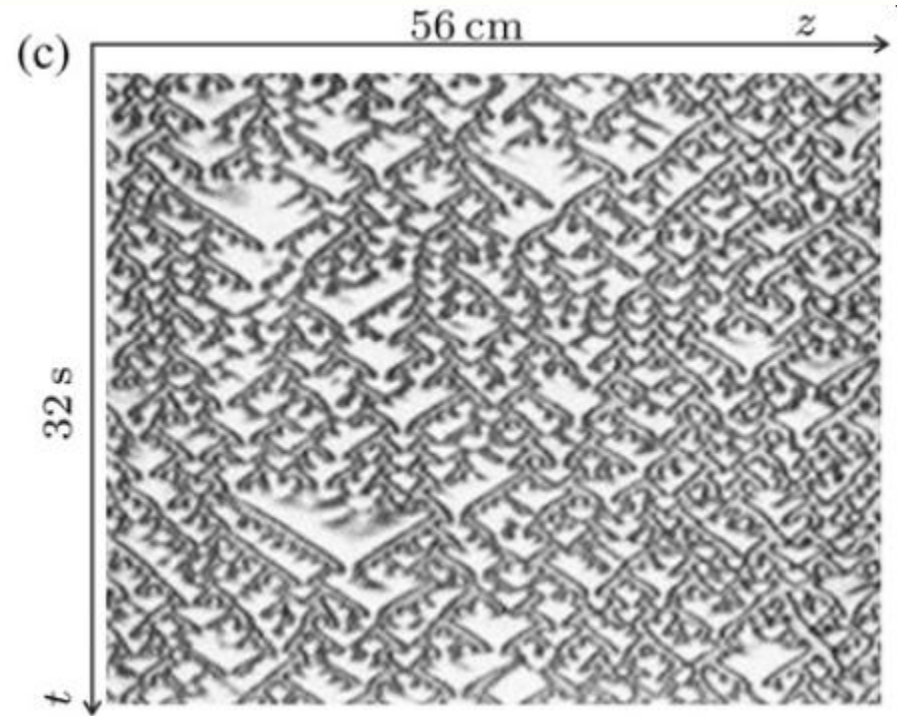
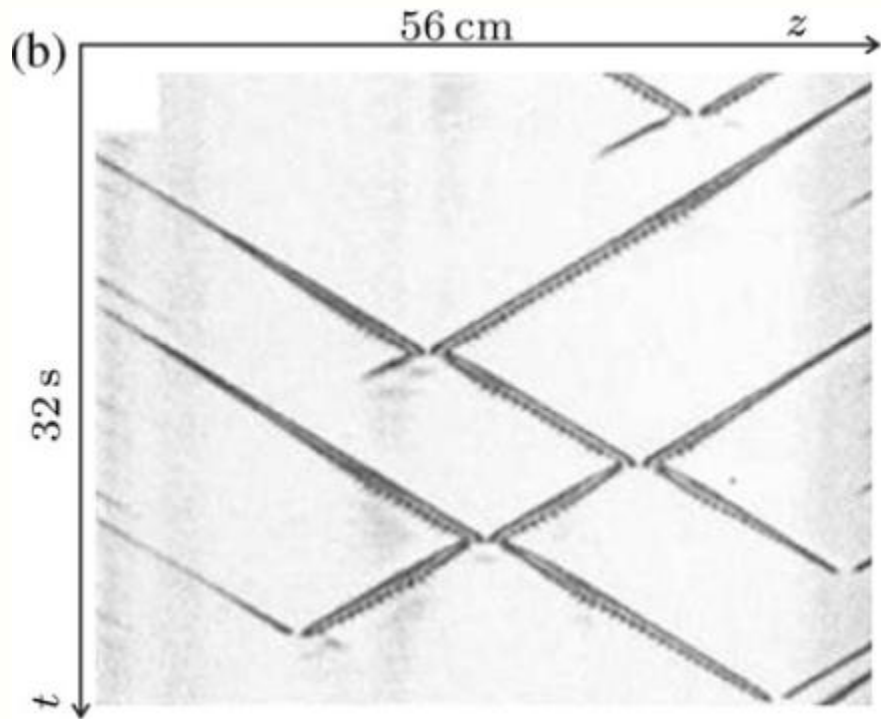
Film Dragging in a Rotating Tube Partly Filled with a Fluid I: Experimental Set-Up with Indication of a Localized Structure

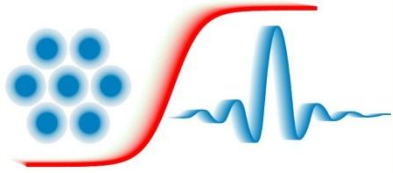
(a)





Film Dragging in a Rotating Tube Partly Filled with a Fluid II: Travelling Localized Structures



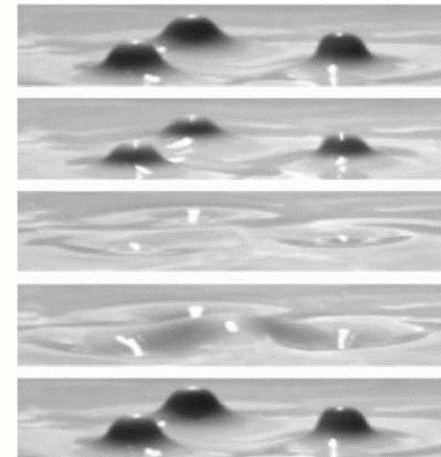


Vertically Driven Systems: Granular Material and Fluid



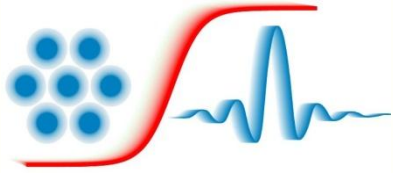
**vertically driven granular
medium**

Umbanhowar et al., Nature 382 p. 793. (1996)

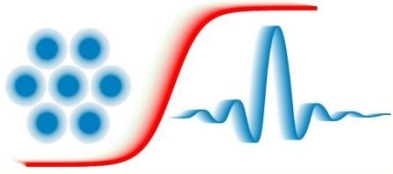


**vertically driven
liquid**

Lioubashevski et al., Phys. Rev. Lett. 83, p. 3190 (1999)

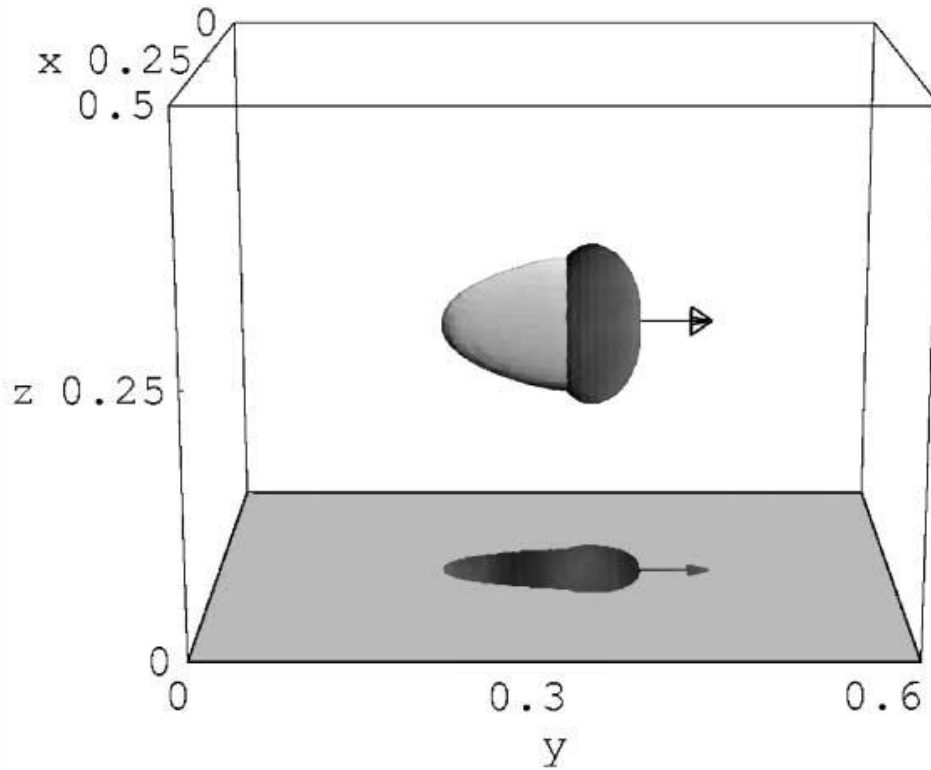


10. Dissipative Soliton (DS) Solutions of the Generalize FitzHugh-Nagumo (FN) Equation In 3-Dimensional Space



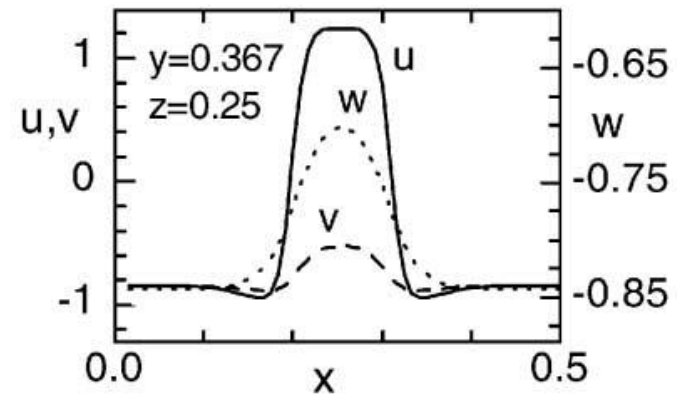
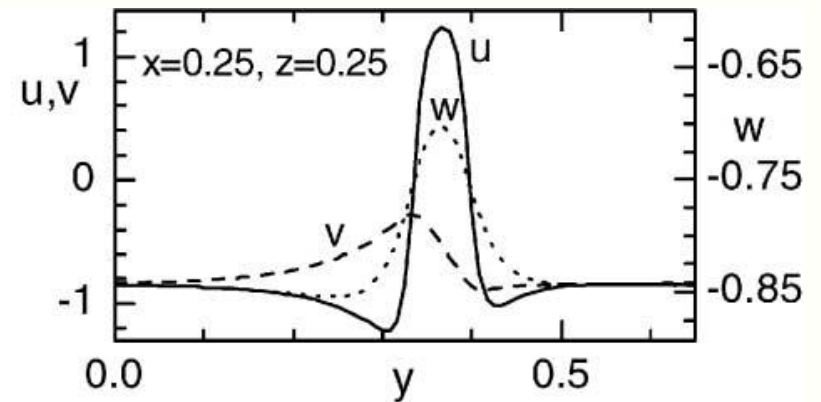
Numerical Solution of the Generalized FN Equation in 3-Dimensional Space: Isolated Travelling DS I

pseudo 3-dimensional representation

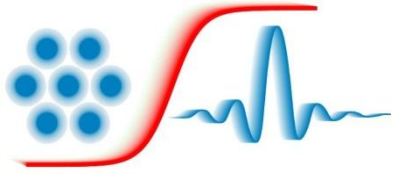


(a)

longitudinal and transverse cross section

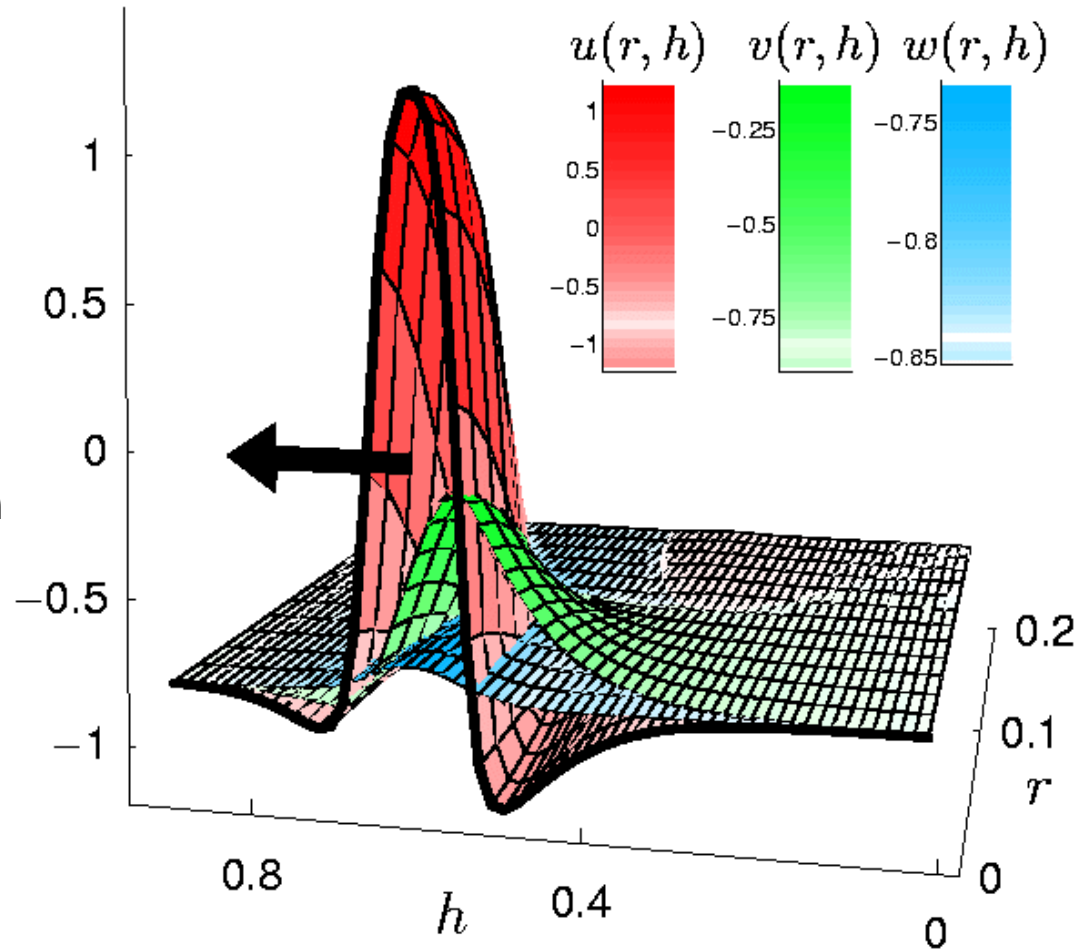


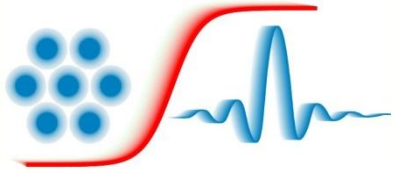
(b)



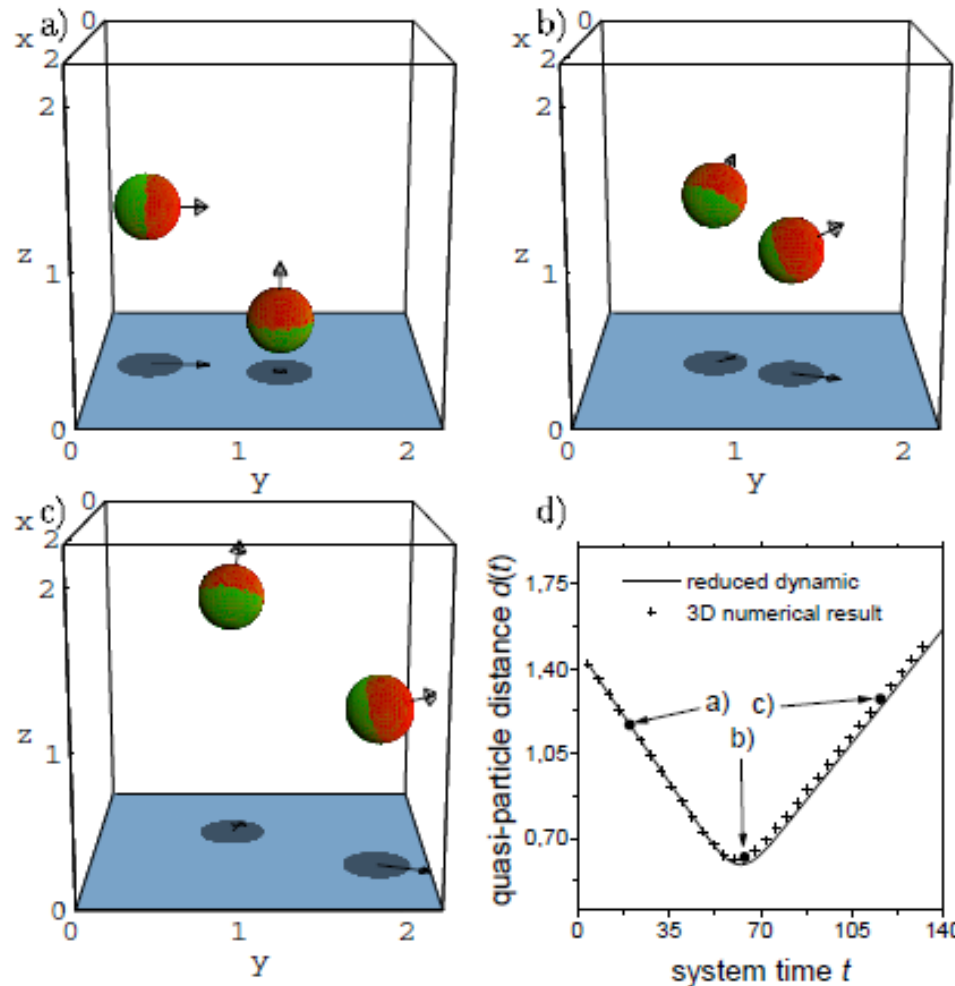
Numerical Solution of the Generalized FN Equation in 3-Dimensional Space: Isolated Travelling DS II

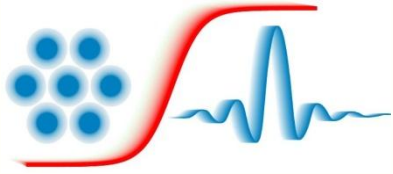
pseudo
3-dimensional
representation



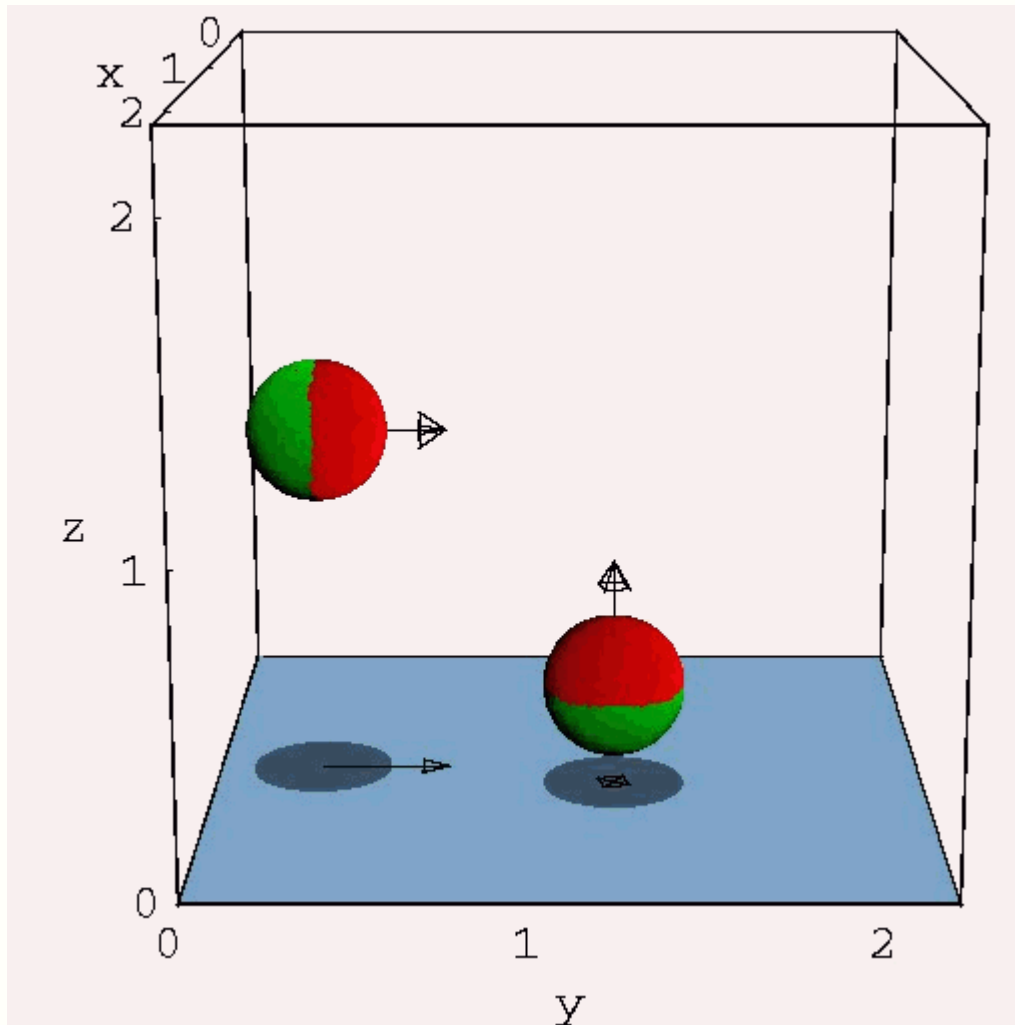


Numerical Solution of the Generalized FN Equation in 3-Dimensional Space: Weak Scattering of 2 DSs I

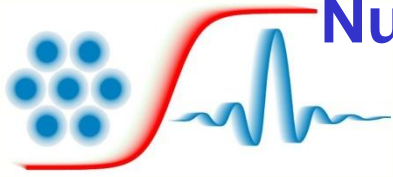




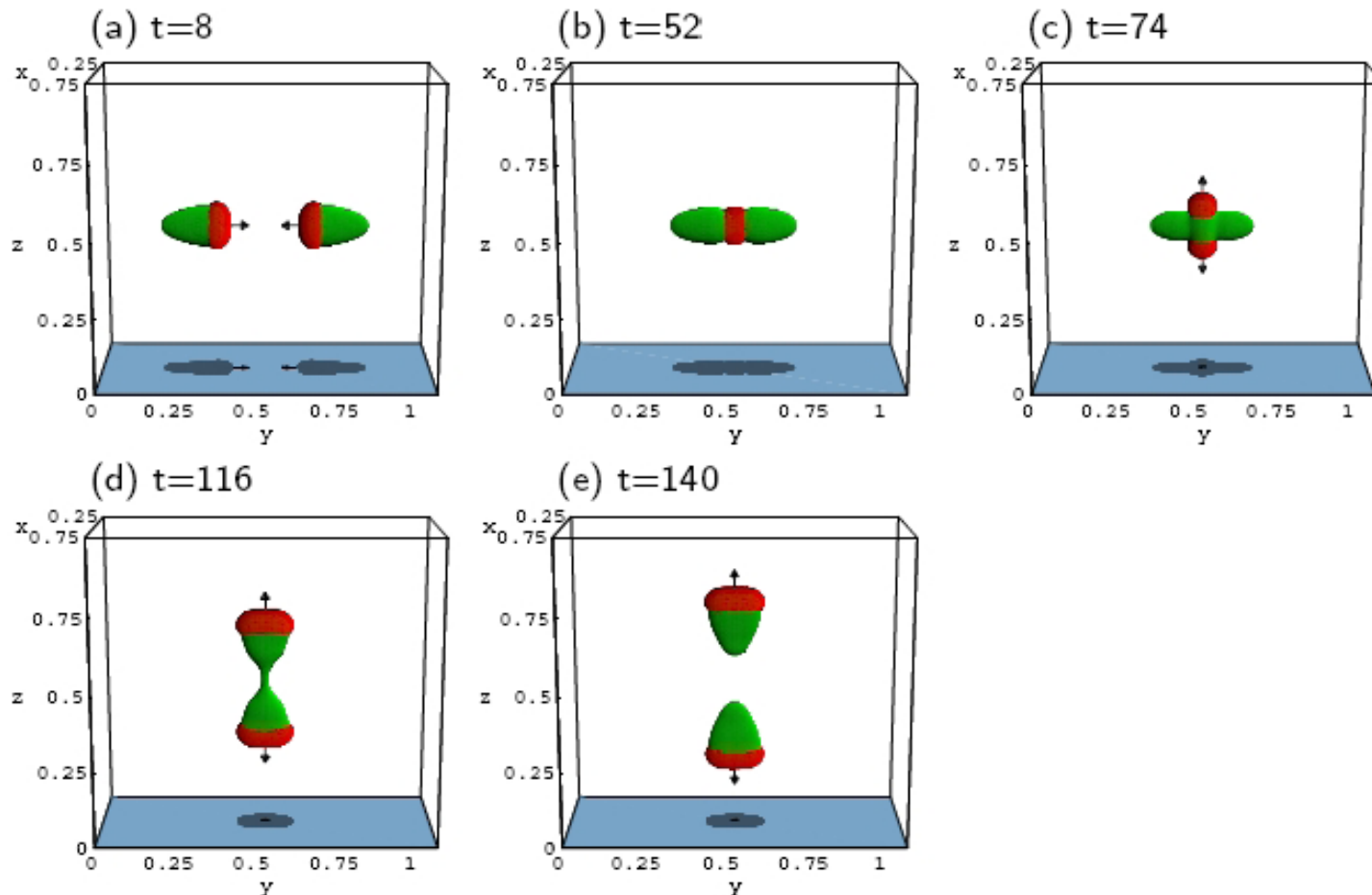
Numerical Solution of the Generalized FN Equation in 3-Dimensional Space: Weak Scattering of 2 DSs II (Movie)

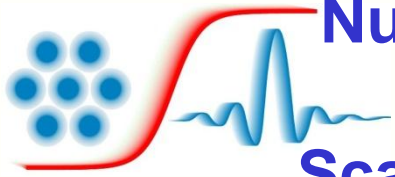


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in the folder

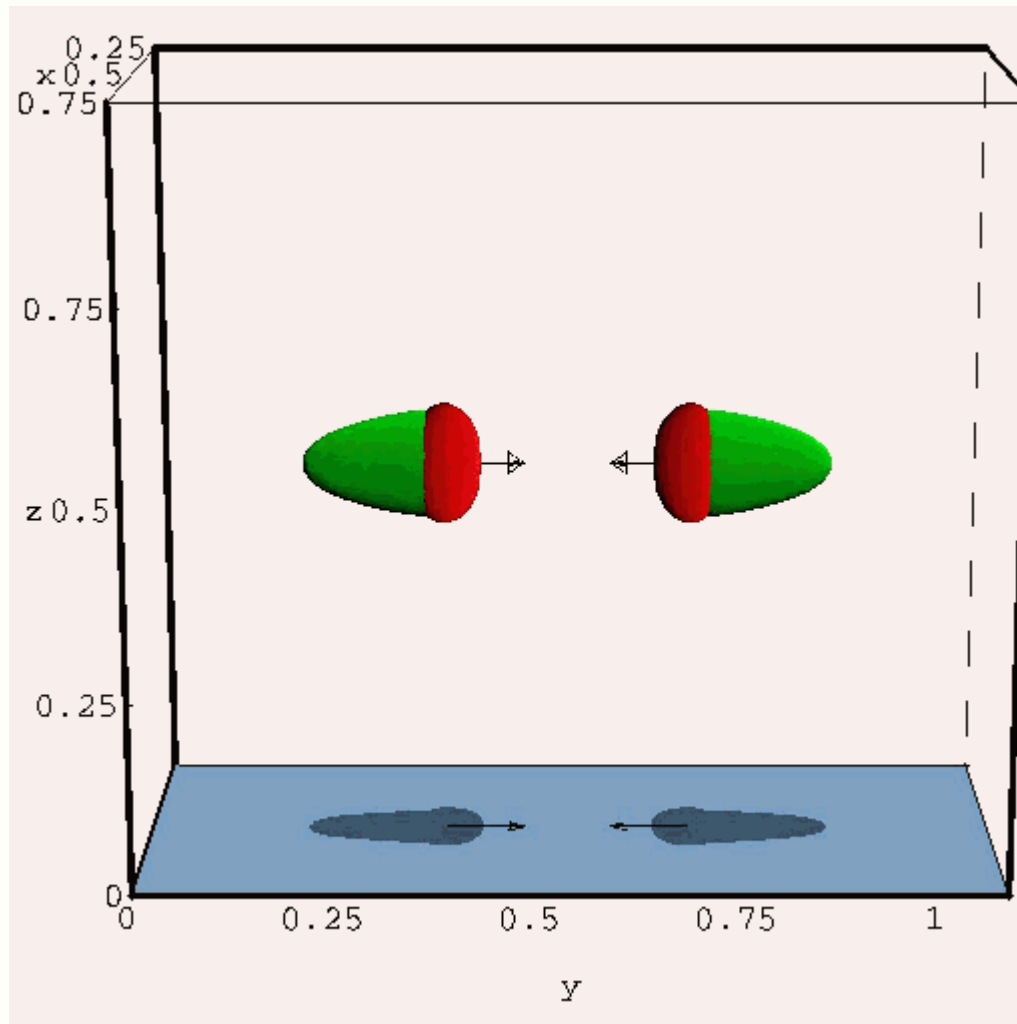


Numerical Solution of the Generalized FN Equation in 3-Dimensional Space: Scattering of 2 DSs with Strong Interaction I

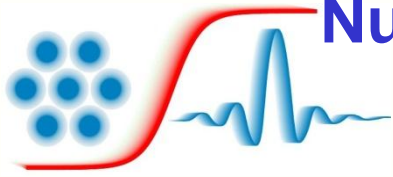




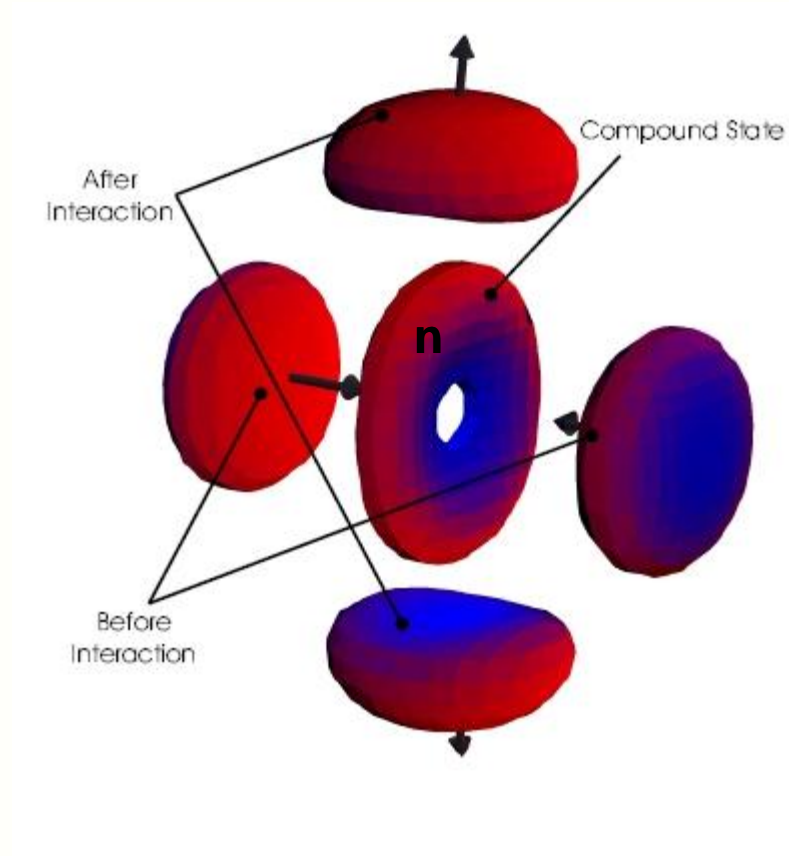
Numerical Solution of the Generalized FN Equation in 3-Dimensional Space: Scattering of 2 DSs with Strong Interaction II (Movie)

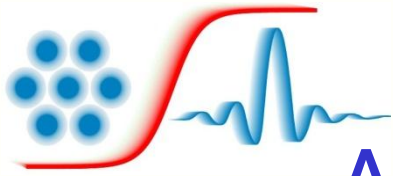


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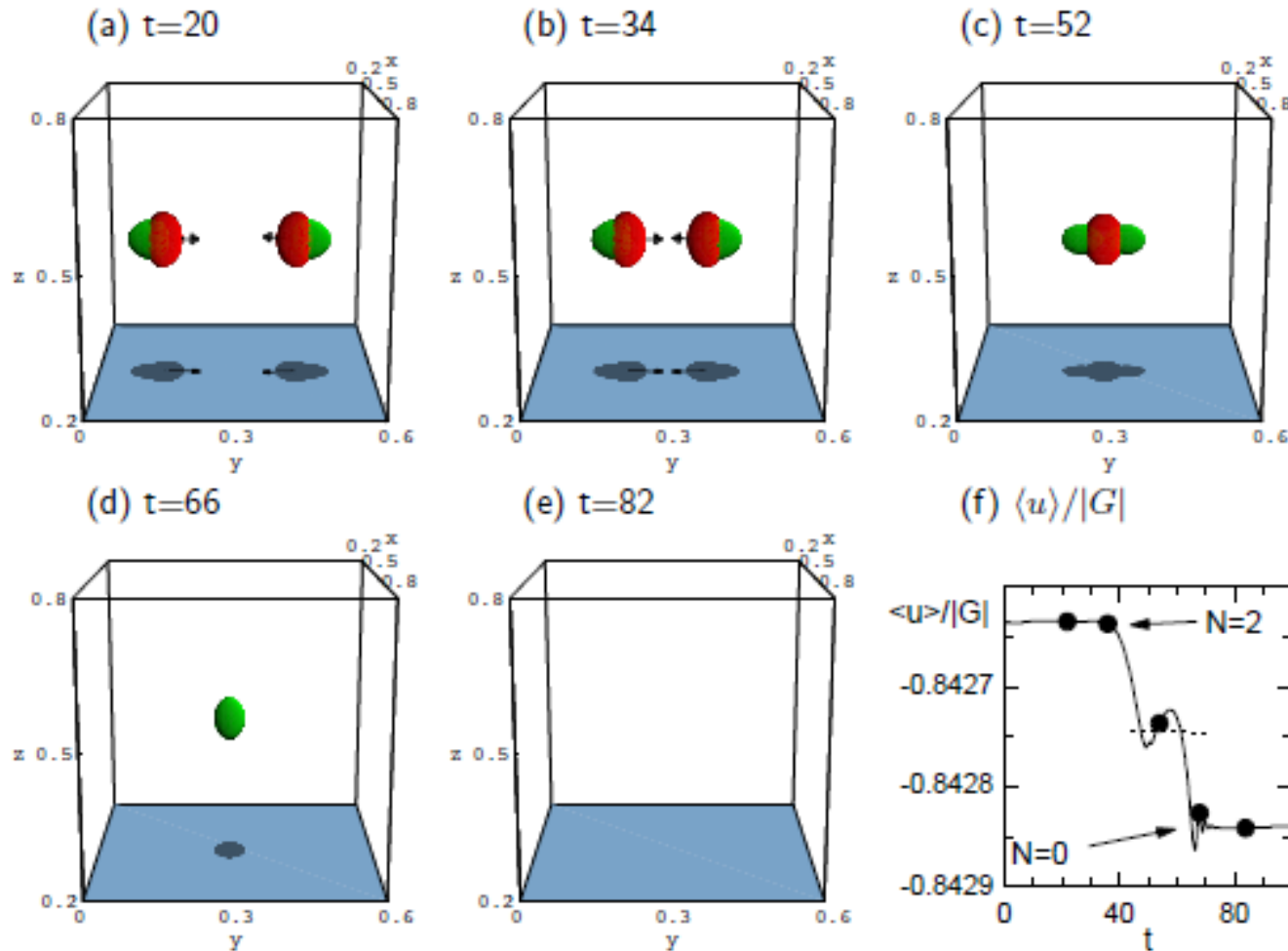


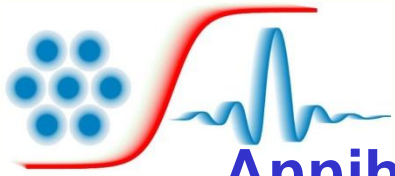
Numerical Solution of the Generalized FN Equation in 3-Dimensional Space: Scattering of 2 DSs with Strong Interaction III



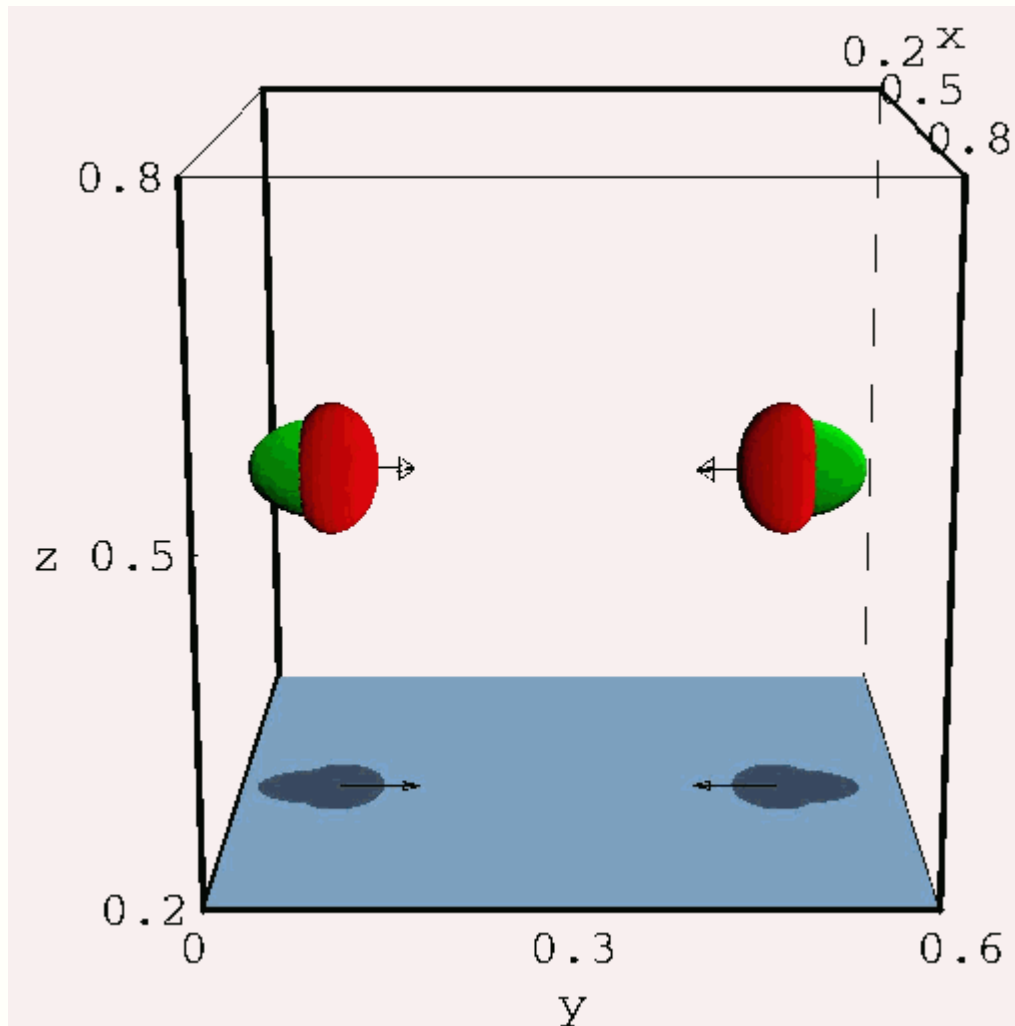


Numerical Solution of the Generalized FN Equation in 3-Dimensional Space: Annihilation in the Course of Collision of 2 DSs

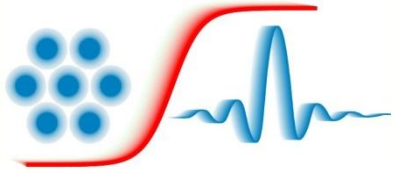




Numerical Solution of the Generalized FN Equation in 3-Dimensional Space: Annihilation in the Course of Collision of 2 DSs (Movie)

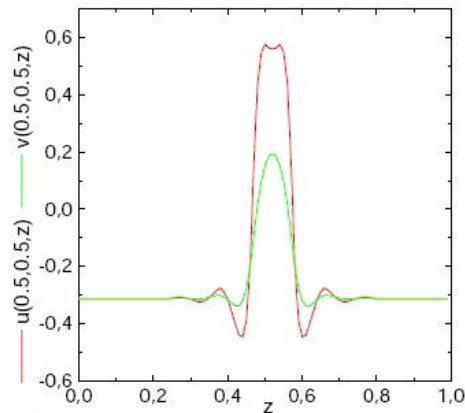


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ung.avi"
in the folder

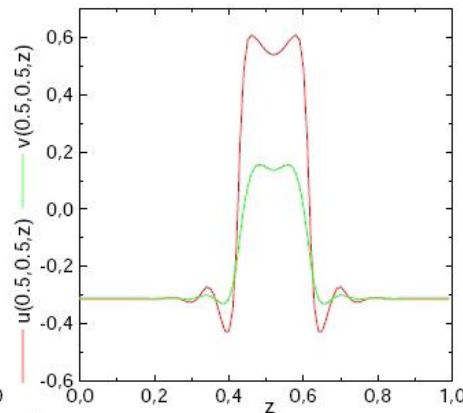


Numerical Solution of the FN-Equation in 3-Dimensional Space: Self-Completion I

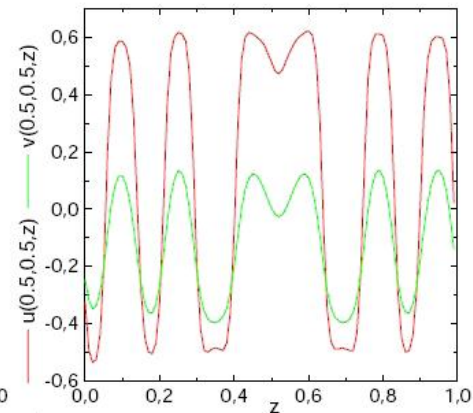
a) $t=0$



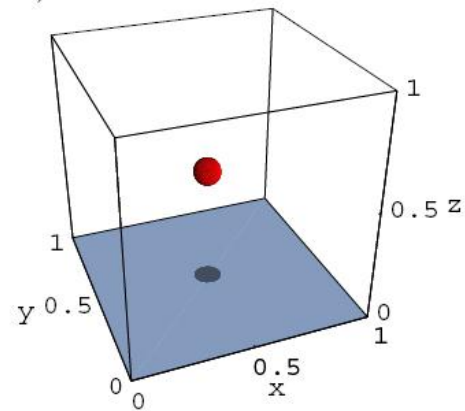
b) $t=230$



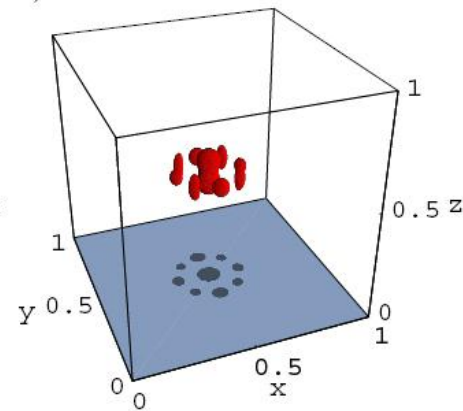
c) $t=600$



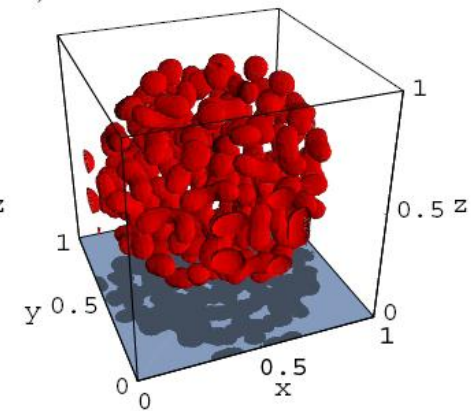
d) $t=0$

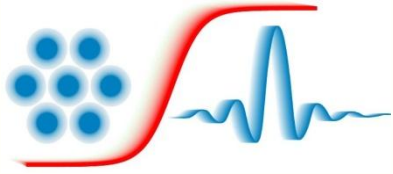


e) $t=230$

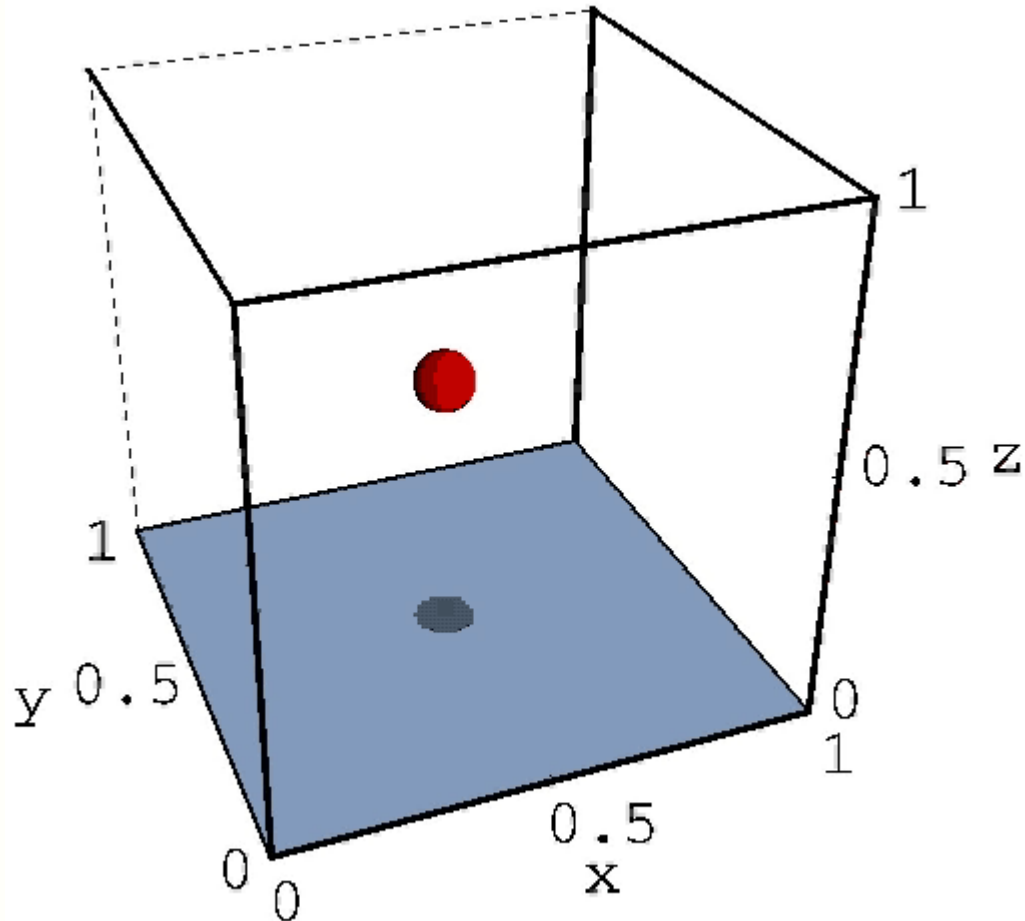


f) $t=600$

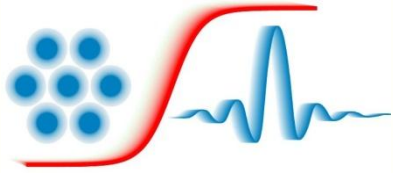




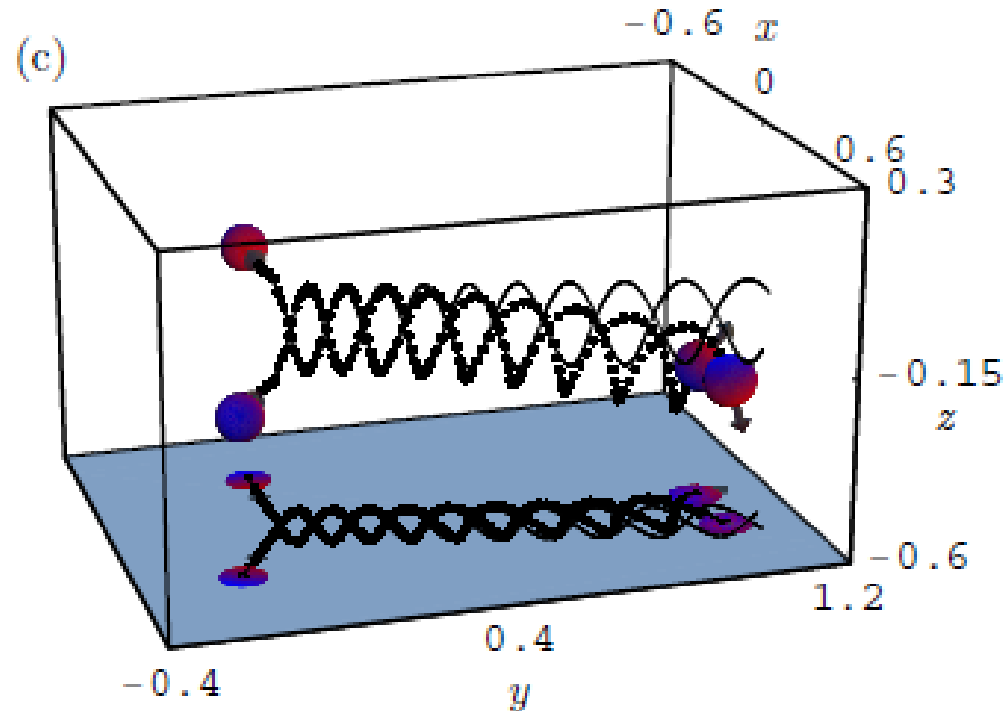
Numerical Solution of the FN-Equation in 3-Dimensional Space: Self-Completion II (Movie)



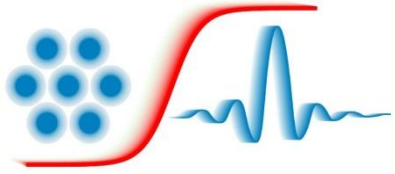
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Destabilisierung.avi”
in the folder



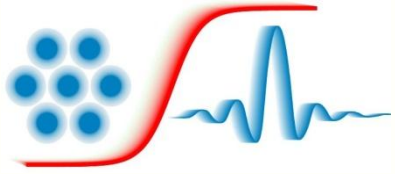
Numerical Solution of the FN-Equation in 3-Dimensional Space: Dynamic Cluster of 2 DSs (“Molecule”)



the 2-DS cluster undergoes propagation along the axis that is vertical to the line connecting the centers of the individual DSs, simultaneously the cluster rotates around this axis



11. Summary



The End