

Some Localised Pattern Problems

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Pioneering research
and skills

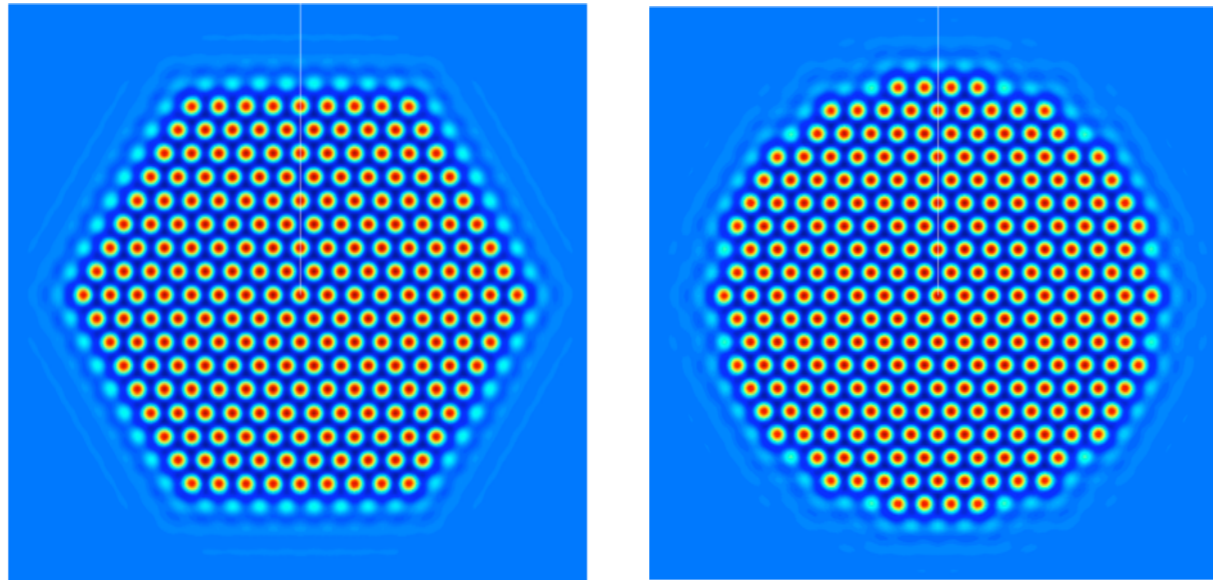


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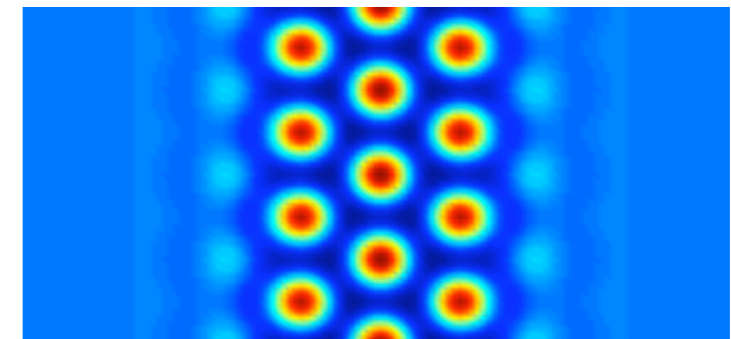
- **Wulff Structures and localised patches**
- **Snaking near the singular limit**
- **Ferrofluids**

Large scale patch shape

- **Hexagon patches related to Hexagon fronts**

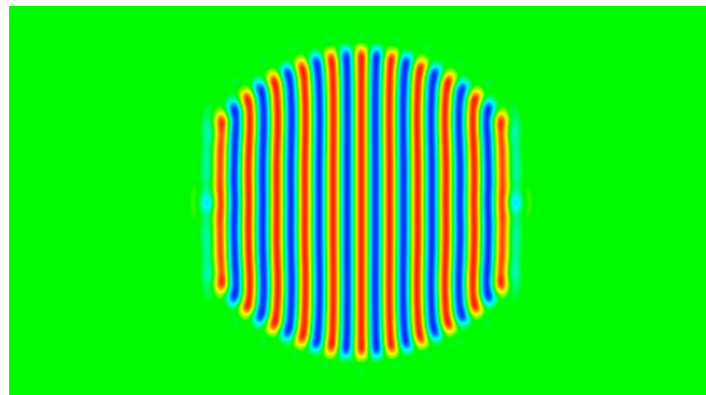


Localised hexagon patches



hexagon pulses

- **What about other types of patches?**

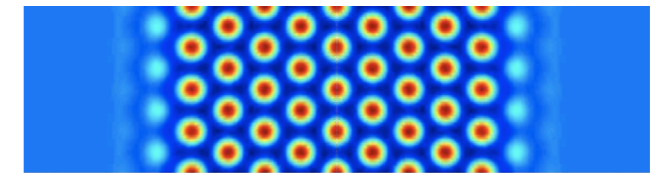
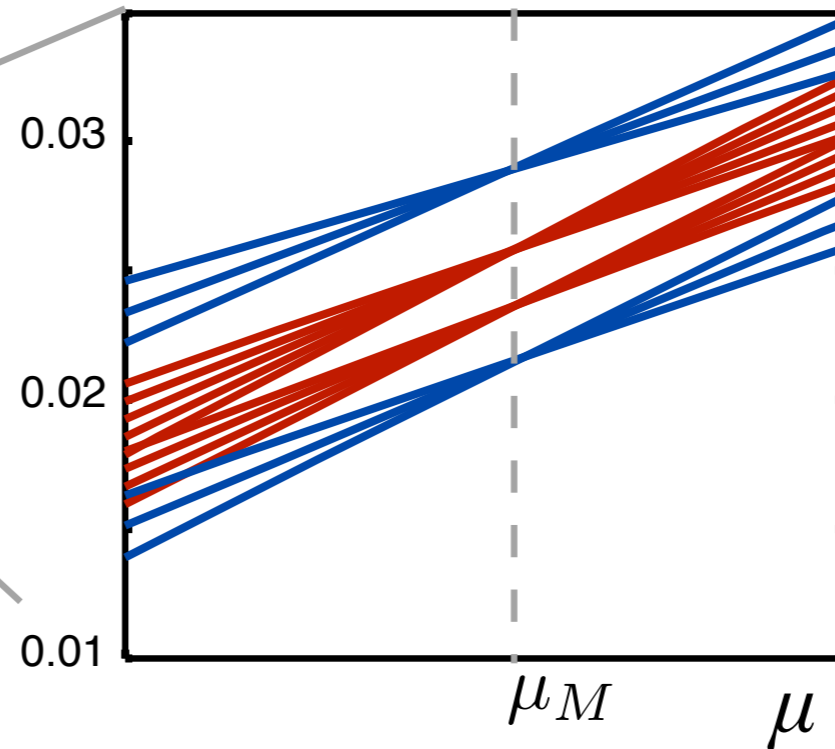
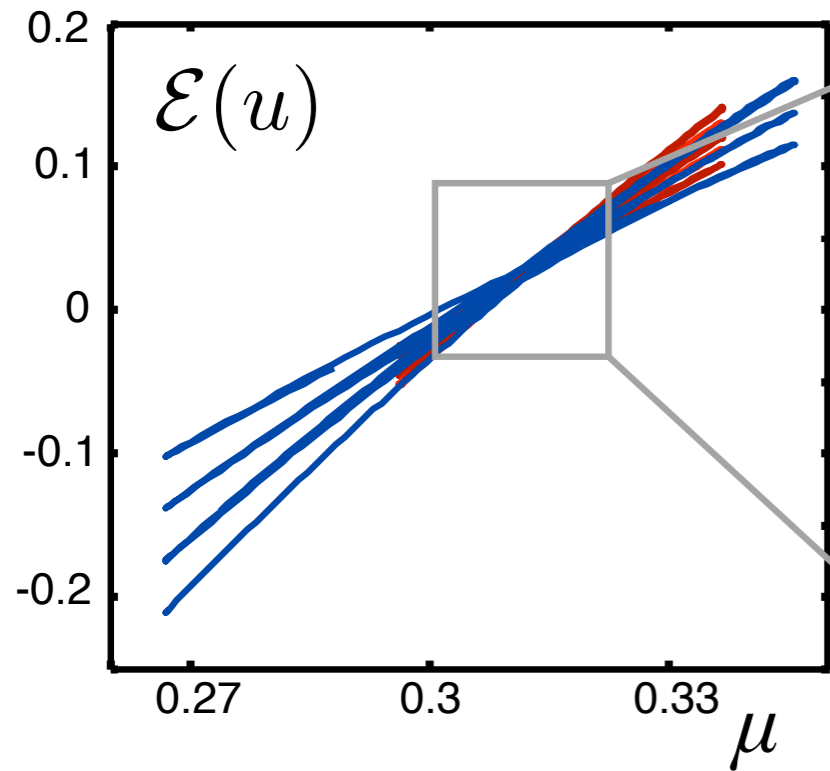


Localised roll patches

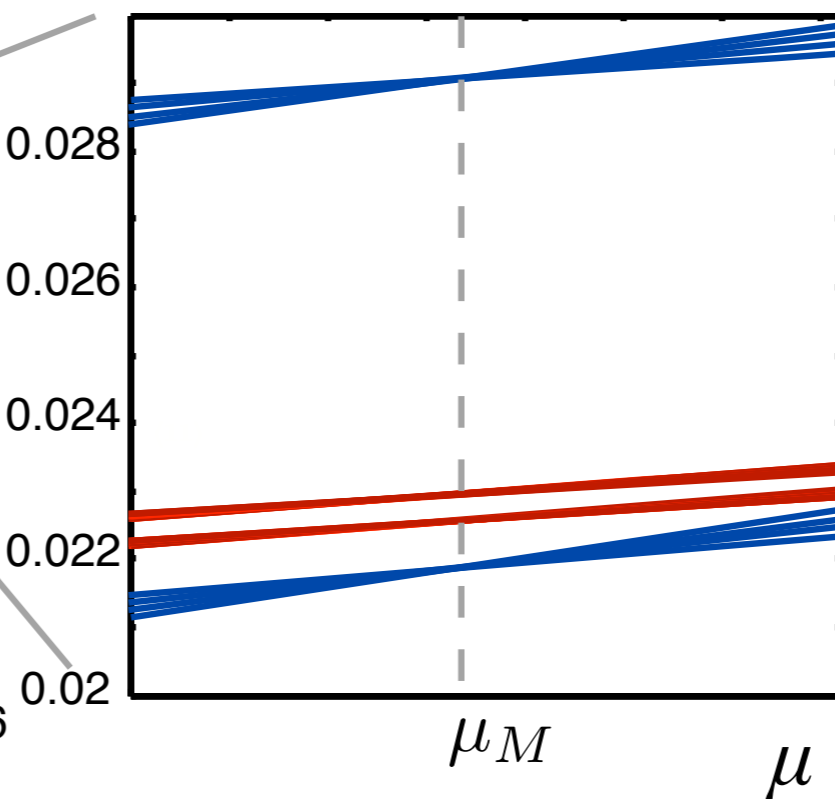
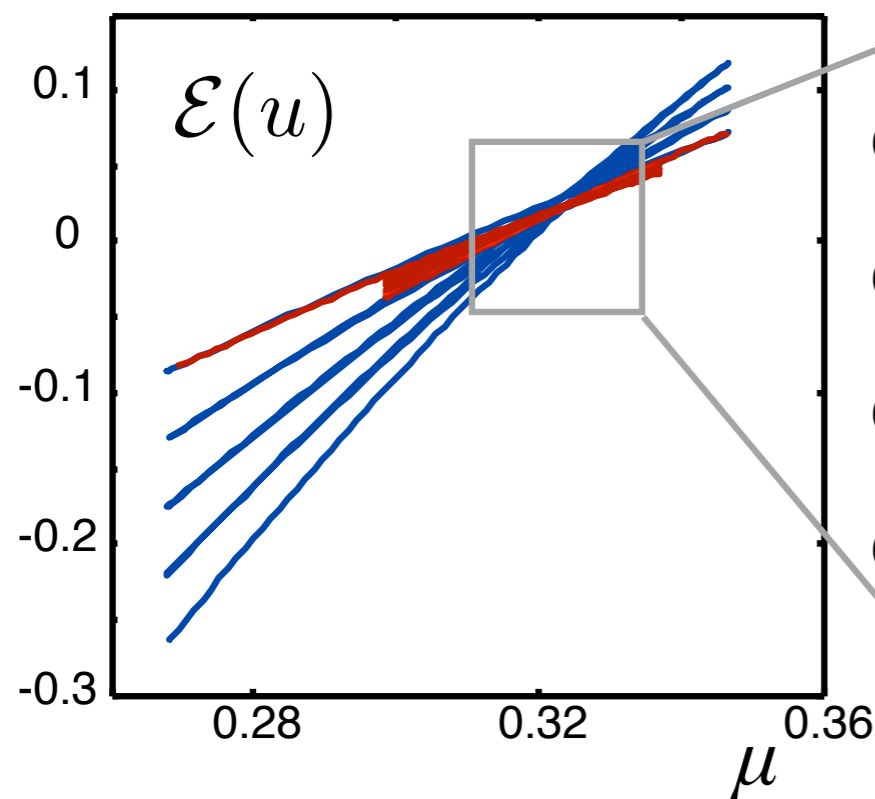
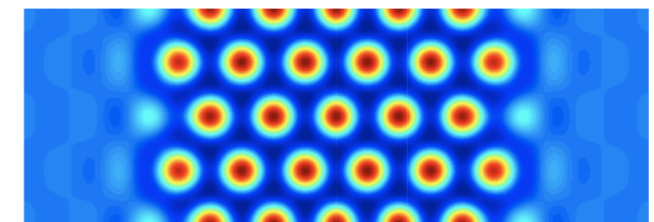


Localised square patches

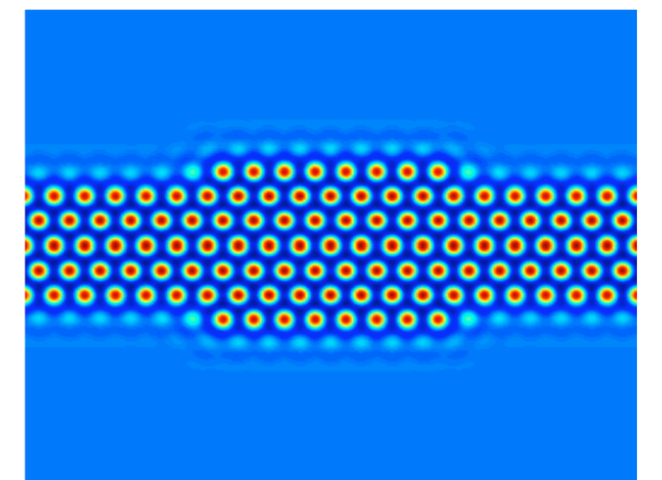
Energy of planar hexagons



[10] hexagon front
[11] hexagon front



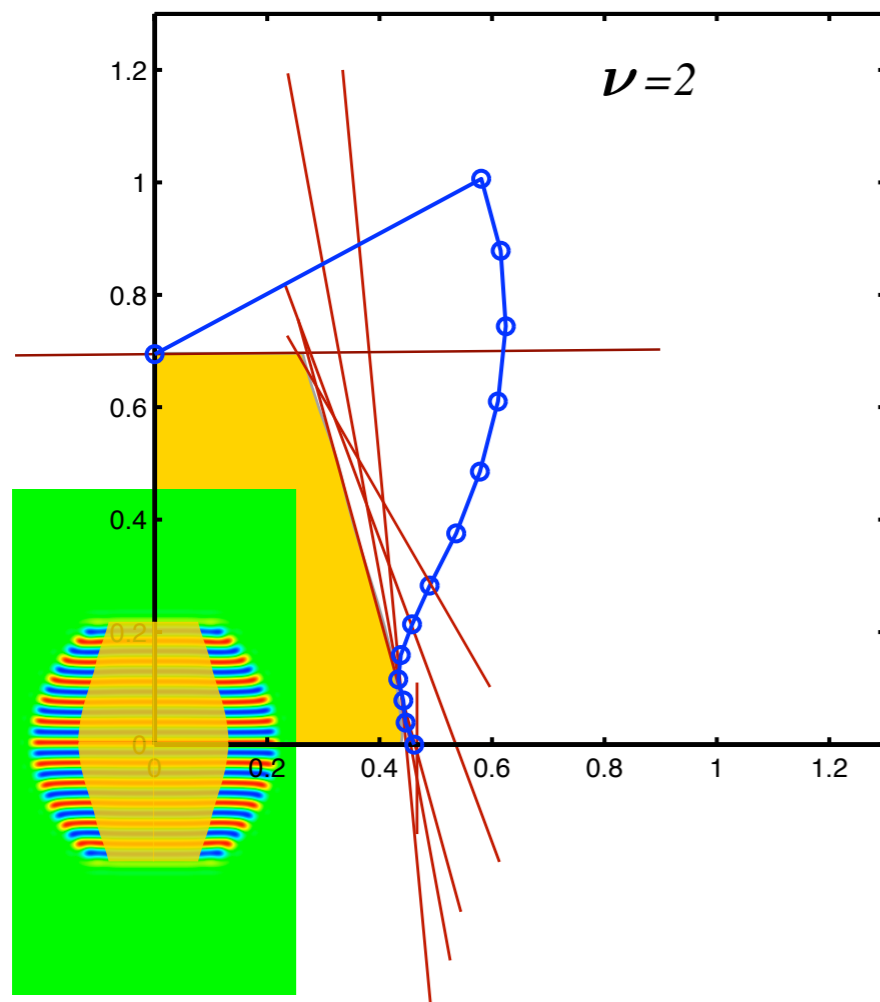
Hexagon front [10]
Almost [10] front



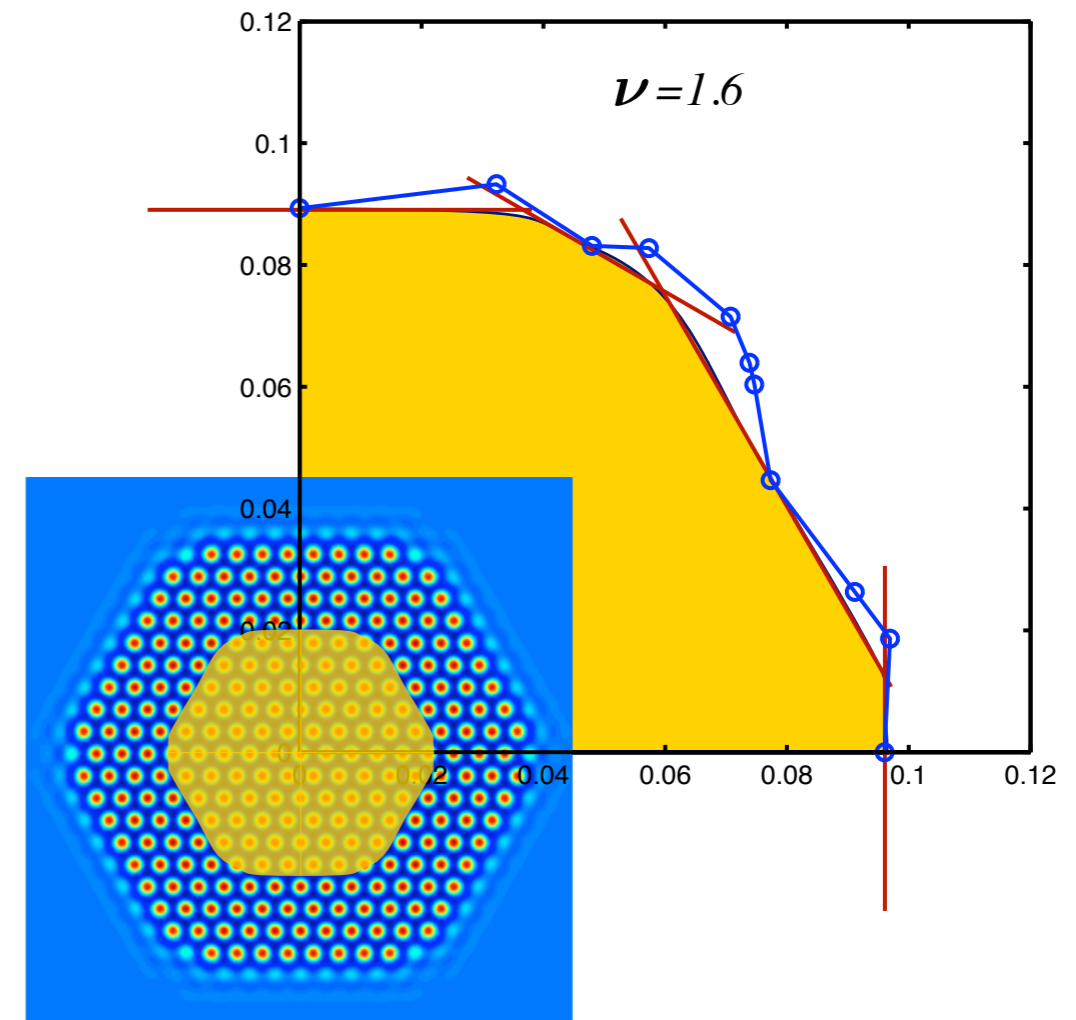
Wulff Constructions

- Compute fronts in Skew coordinates

Wulff construction for rolls



Wulff construction for hexagons



- Near co-dim 2 point, asymptotics \rightarrow hexagon Wulff construction circle
- What happens for other localised patches?

Snaking near the singular limit

- **Crime hotspot model:**

- Short et al. (Math. Models Meth. Appl. Sci. 1249 (18) 2008)

- **Non-dimensional Mean field equations with:**

$$A_t = \eta \nabla^2 A - A + A_0 + \rho A,$$
$$\rho_t = \nabla \cdot \left(\nabla \rho - \frac{2\rho}{A} \nabla A \right) - \rho A + \bar{A} - A_0$$

- **Spatially Homogenous equilibrium:**

$$A = \bar{A}, \quad \rho = 1 - \frac{A_0}{\bar{A}} \quad \text{Stable: } A_0 > A_0^{\text{crit}}$$

- **Turing instability:** $A_0^{\text{crit}} = \frac{2}{3}\bar{A} - \frac{1}{3}\eta\bar{A}^2 - \frac{2}{3}\bar{A}\sqrt{\eta\bar{A}}$

- **1D Subcritical:** $\sqrt{\eta\bar{A}} \lesssim 0.157$

- **Steady problem can be re-written as a 4th-order PDE:**

$$-\eta B_{xxxxx} - 2B_{xx} + 2\frac{B_x^2}{B+\bar{A}} + 4\eta\frac{B_x B_{xxx}}{B+\bar{A}} - 6\eta\frac{B_x^2 B_{xx}}{(B+\bar{A})^2} - 6A_0\frac{B_x^2}{(B+\bar{A})^2} + 3\eta\frac{B_{xx}^2}{B+\bar{A}} + 3A_0\frac{B_{xx}}{B+\bar{A}} - \bar{A}B - B^2 + (B+\bar{A})\eta B_{xx} = 0$$

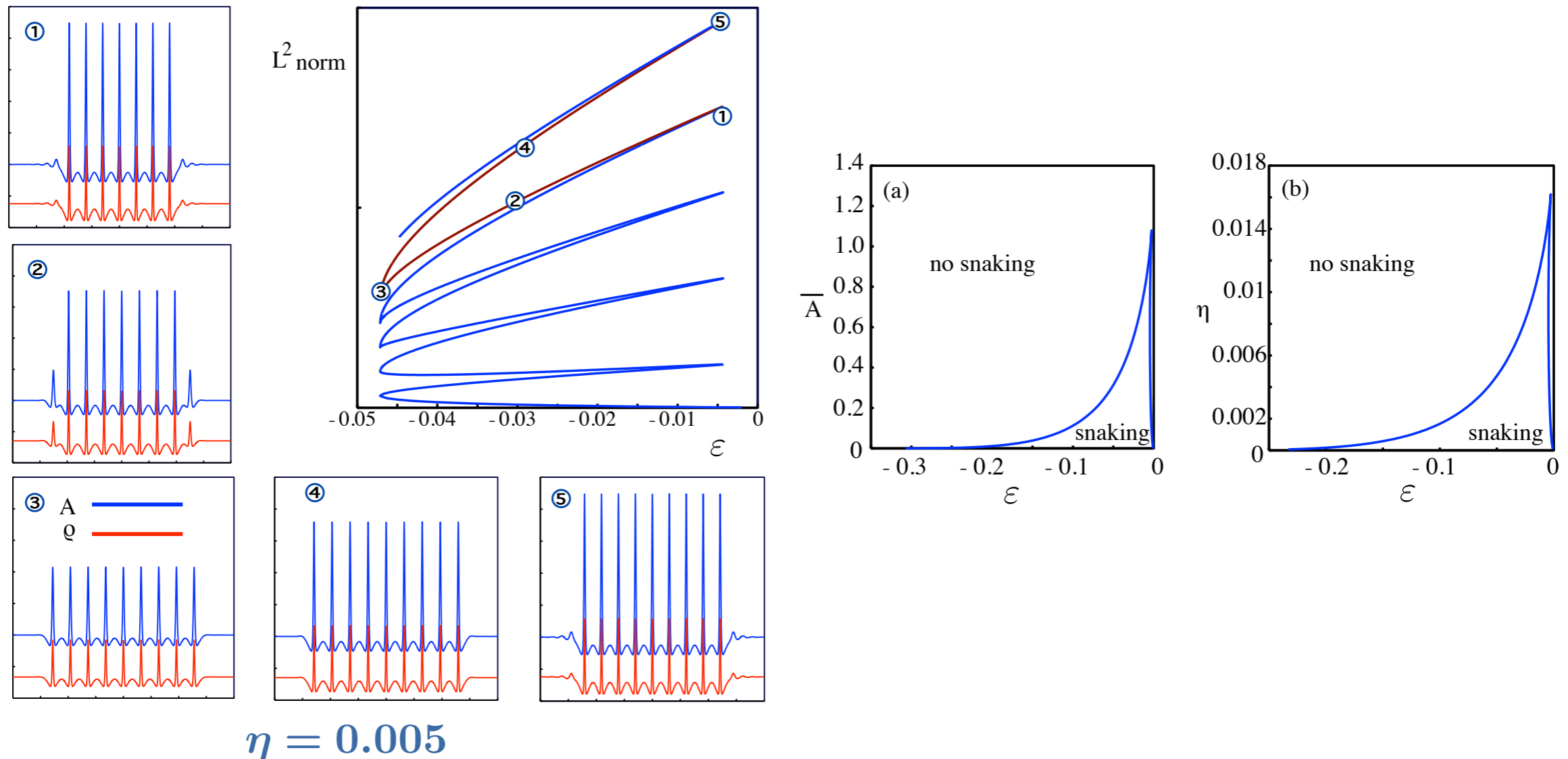
Homoclinic Snaking

- **Weakly nonlinear analysis:** $A(x,t) \approx \bar{A} + \epsilon^{\frac{1}{2}} P(X,T)e^{ix} + c.c.$

$$P_T = \eta P_{XX} + \epsilon P - \nu |P|^2 P$$

where: $A_0 = A_0^{\text{crit}} - \epsilon \bar{A}$

- **Two pulses bifurcate off equilibrium...**



Singular Limit Analysis

- **Change of coordinate:** $v = \frac{\rho}{A^2}$

- **Rescale:** $\eta \rightarrow \eta^2$, $A \rightarrow \frac{A}{D}$, $v \rightarrow D^2 v$, $\bar{A} \rightarrow \frac{\bar{A}}{D}$, $A_0 \rightarrow \frac{A_0}{D}$

$$\eta^2 A_{xx} - A + vA^3 + A_0 = 0,$$

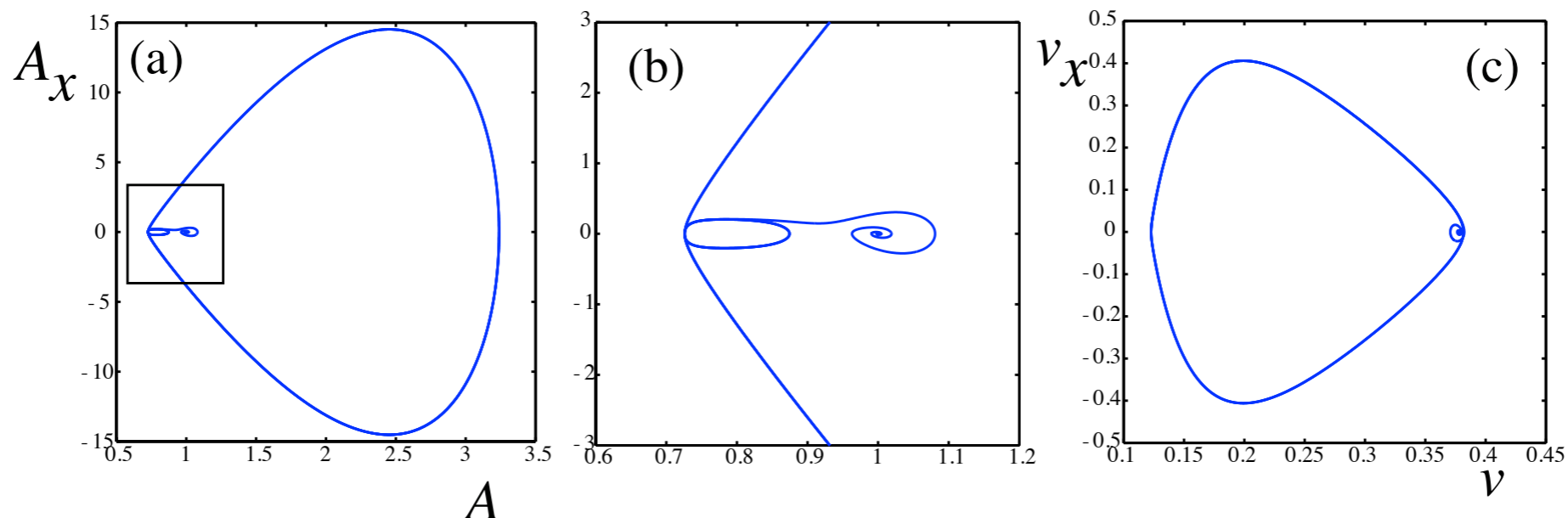
$$D(A^2 v_x)_x - vA^3 + \bar{A} - A_0 = 0$$

- **Kolokolnikov/Ward/Sun (2011) look at $D \gg 1$**

$$A(x) \sim \begin{cases} a \operatorname{sech}(x/\eta) & x = \mathcal{O}(\eta), \\ A_0 & x \gg \mathcal{O}(\eta) \end{cases} \quad v(x) \sim v_0 \ll 1$$

- **Interesting limit $D \ll 1$**

ID Subcritical Turing: $\sqrt{\frac{\eta^2 \bar{A}}{D}} \lesssim 0.157$



Slow limit

- **Slow system:**

$$\eta A_x = B,$$

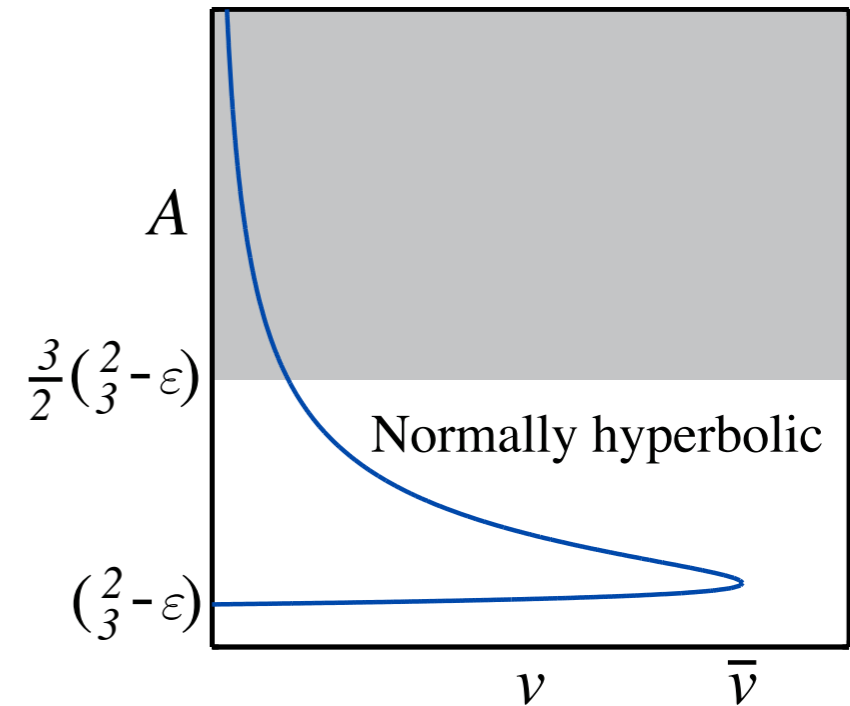
$$\eta B_x = A - vA^3 - A_0,$$

$$v_x = \frac{w}{A^2},$$

$$w_x = vA^3 - \bar{A} + A_0,$$

- **Setting $\eta = 0$**

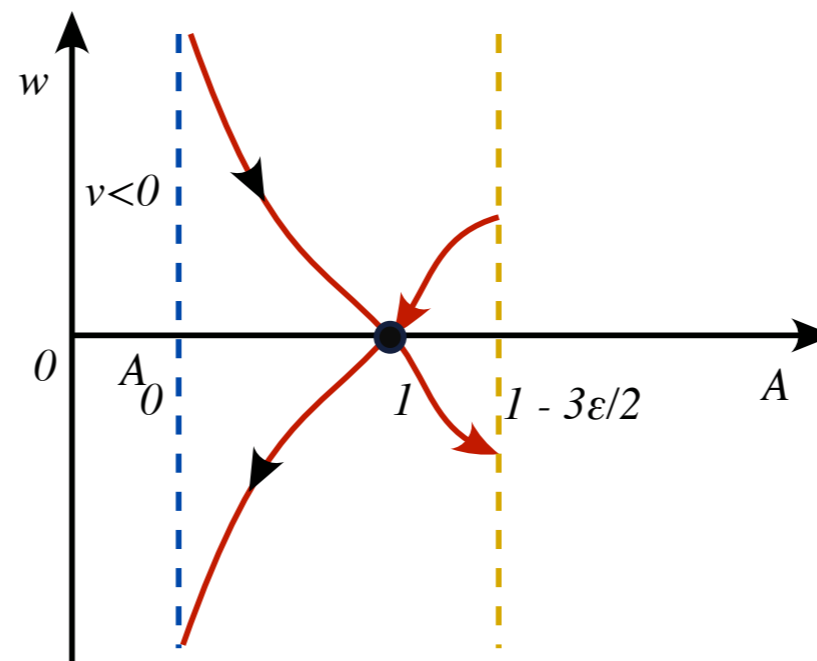
$$B = 0, \quad A - vA^3 - \left(\frac{2}{3} - \epsilon\right)\bar{A} = 0$$



- **Re-writing slow flow:**

$$A_x = \frac{A^2 w}{(2 - 3\epsilon - 2A)},$$

$$w_x = A - 1,$$



Fast limit

- Fast system:**

$$A_x = B,$$

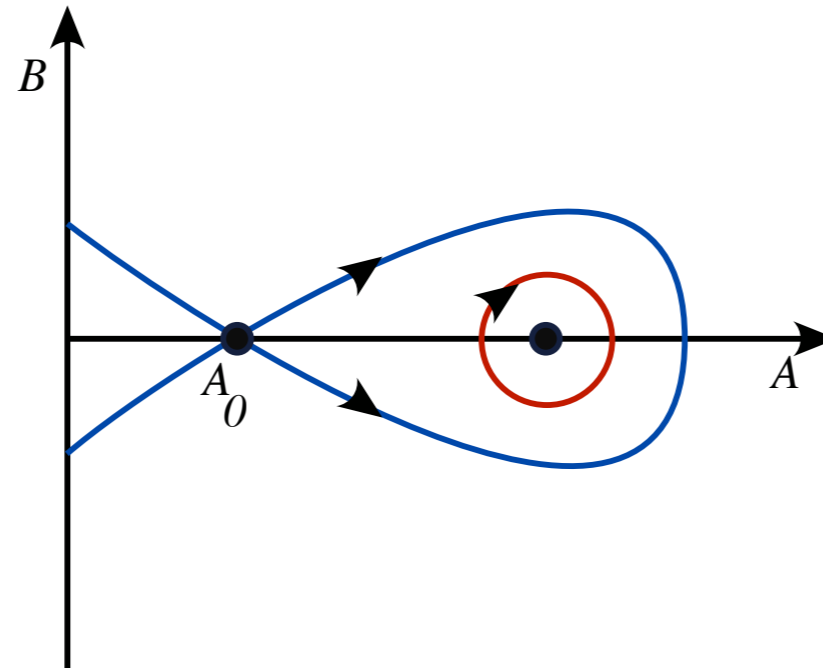
$$B_x = A - vA^3 - A_0,$$

$$v_x = \eta \frac{w}{A^2},$$

$$w_x = \eta v A^3 - \eta(\bar{A} - A_0),$$

- Setting:** $\eta = 0$

$$v(y) = v_0 \quad w(y) = w_0 \quad A_{yy} = A - v_0 A^3 - \left(\frac{2}{3} - \epsilon\right) \bar{A}$$

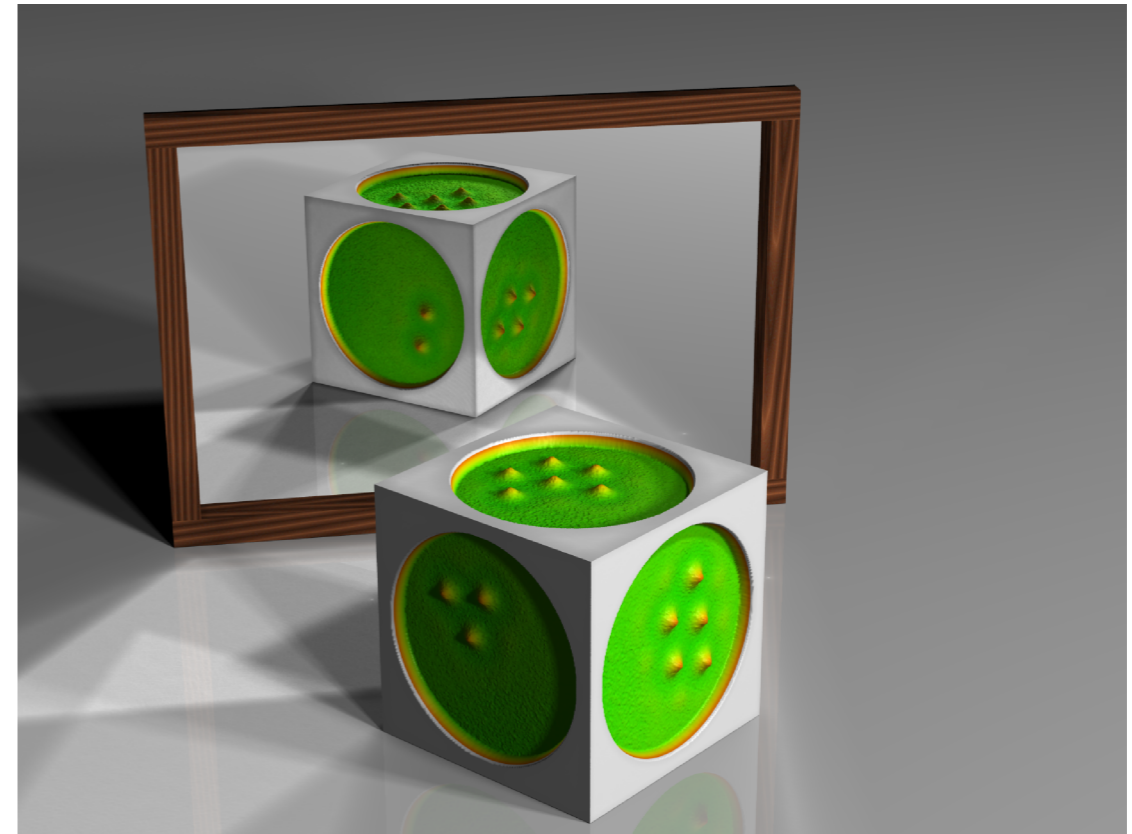


- Matching? cf. Gray Scott problem**



Ferrosolitons

[Richter & Barashenkov]



Ferro-hexagon patches

[Gollwitzer, Rehberg & Richter]

Q: Weakly interacting spots or hexagon patches?

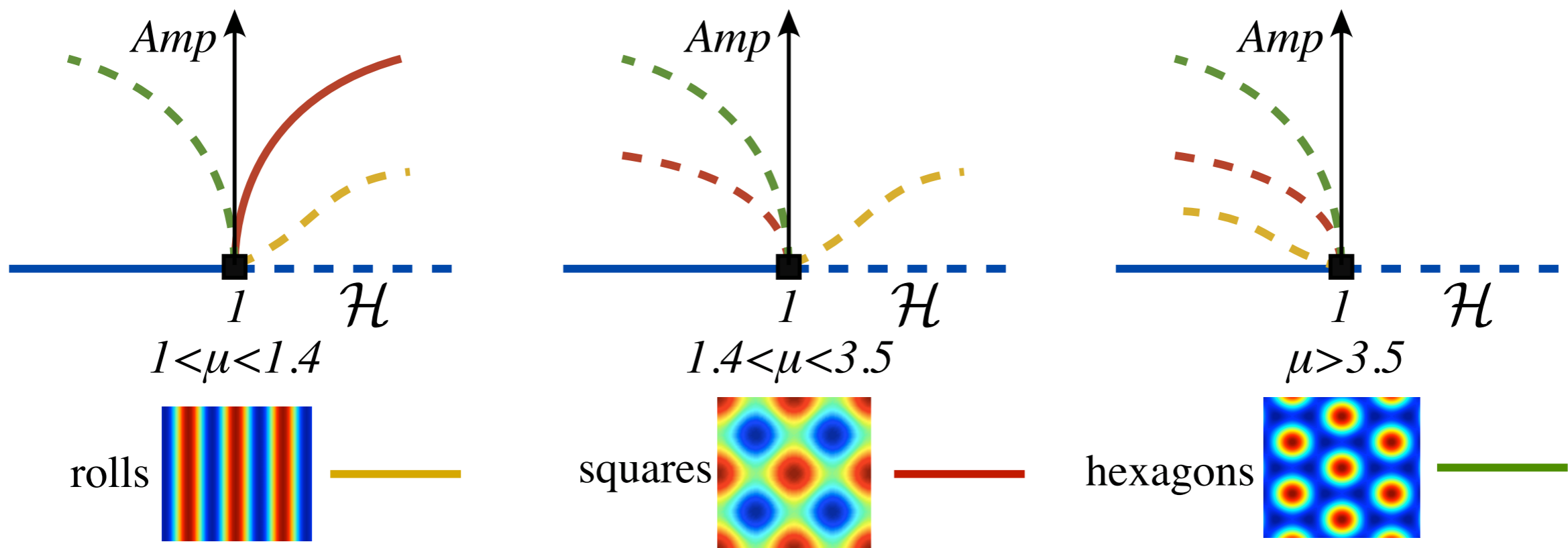
Periodic Ferro-patterns

Linear Magnetisation Law

2 non-dimensional params:

\mathcal{H} - rescaled uniform magnetic field $\Rightarrow \mathcal{H} = 1$ onset of instability

μ - permeability of ferrofluid



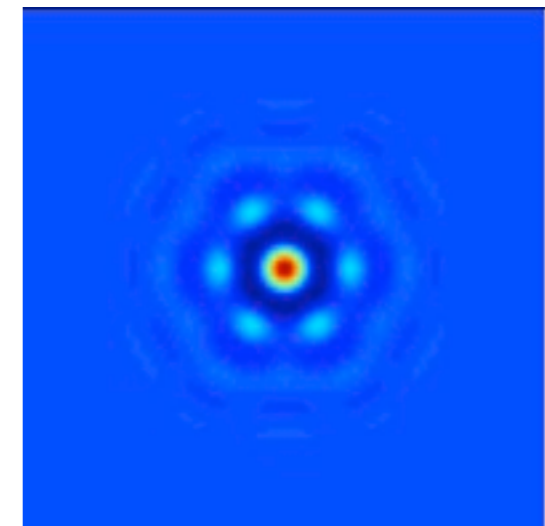
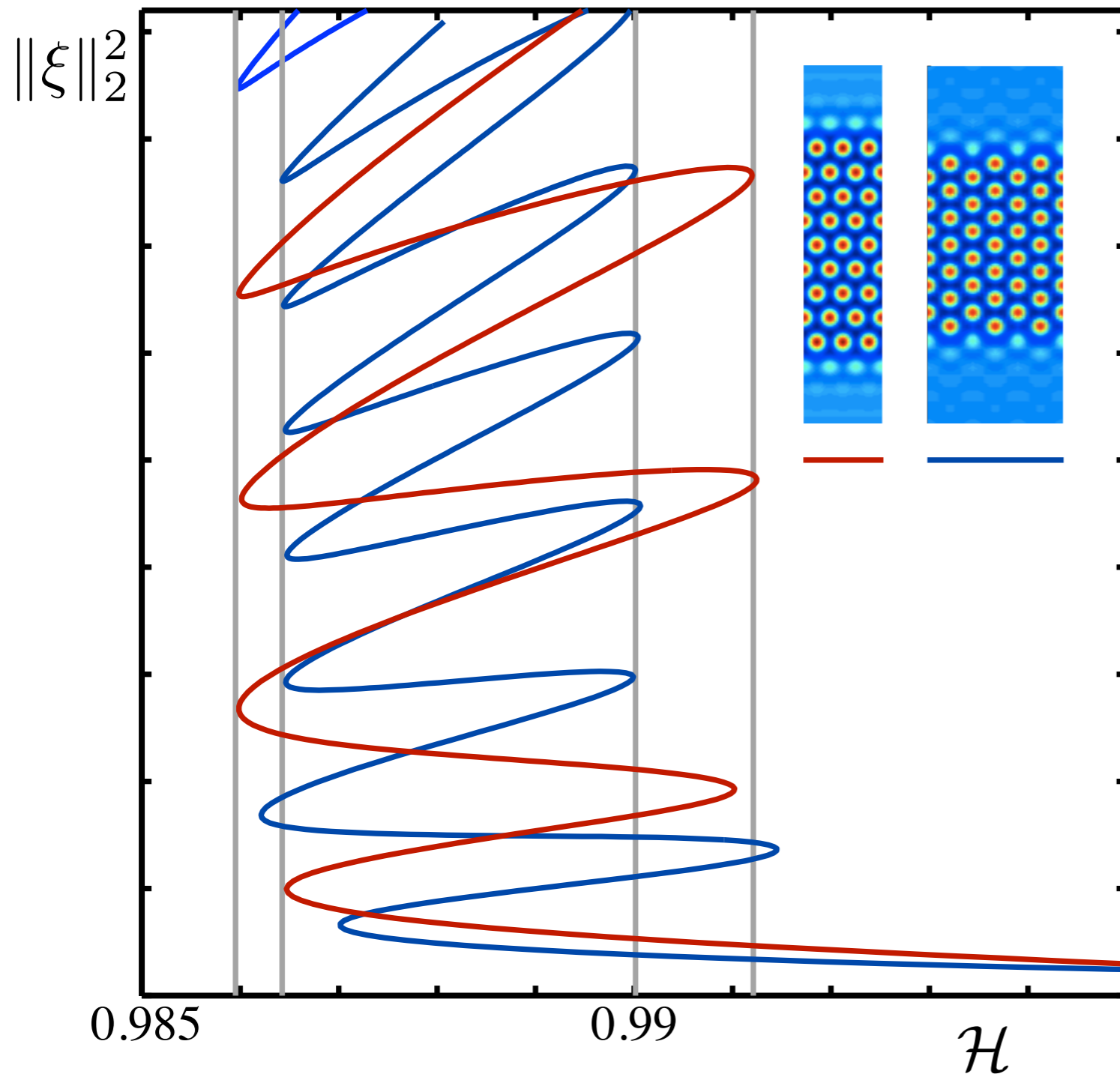
[Silber & Knobloch]

Richter & Barashenkov ferrofluid $\mu \sim 3.2$

Mystery: Spots should be unstable from SHEqn...

Hexagon Ferro-fronts

$$\mu = 2.5$$



Patch bifurcation off spot

$$\mathcal{H} = 0.995$$

Challenge:

Understand hexagon patches

- Snaking breaks down $\mu \geq 3$

Ferrofluid: Spatial Dynamics

- **Energy:**
$$\mathcal{E}(\phi, \chi, \xi) = \int \int \left[\frac{H_0^2}{2\mu_0} \left(\int_{\xi(x,y)}^{\infty} |\nabla \chi|^2 dz + \frac{1}{\mu} \int_{-\infty}^{\xi(x,y)} |\nabla \phi|^2 dz \right) + \frac{1}{2} \rho g \xi^2 + \tau \sqrt{1 + |\nabla \xi|^2} \right] dx dy$$

Bcs: $\phi = \mu \chi, \quad \phi_z - \chi_z = (\nabla \phi - \nabla \chi) \cdot \nabla \xi.$

- **Flatten out the free-surface:** $\tilde{z} = z - \xi(x, y)$

- **Legendre transformation -> Hamilton's equations**

$$\begin{aligned} \xi_x &= \frac{\tau W}{\sqrt{\tau^2 - W^2}}, \\ \eta_x &= \rho g \xi, \\ \phi_x &= \frac{\mu \mu_0}{H_0^2} \alpha + \frac{H_0^2}{\mu \mu_0} \phi_{\tilde{z}} \xi_x, \\ \chi_x &= \frac{\mu_0}{H_0^2} \beta + \frac{H_0^2}{\mu_0} \chi_{\tilde{z}} \xi_x, \\ \alpha_x &= -\frac{\mu \mu_0}{H_0^2} \phi_{\tilde{z}\tilde{z}} - \frac{\tau H_0^2}{\mu_0 \mu} \alpha_{\tilde{z}} \xi_x, \\ \beta_x &= -\frac{\mu_0}{H_0^2} \chi_{\tilde{z}\tilde{z}} - \frac{\tau H_0^2}{\mu_0} \beta_{\tilde{z}} \xi_x \end{aligned}$$

Bcs:

$$\begin{aligned} \phi &= \mu \chi \\ \phi_z - \chi_z &= \xi_x \frac{\mu_0}{H_0^2} (\mu \alpha - \beta) \end{aligned}$$

$$W = \eta + \frac{H_0^2}{\mu_0} \left[\int_0^{\infty} \beta \chi_{\tilde{z}} d\tilde{z} + \frac{1}{\mu} \int_{-\infty}^0 \alpha \phi_{\tilde{z}} d\tilde{z} \right].$$

- **Equations similar to two-layer interface water waves...**

Conclusion

- **Justification of Wulff construction for PDEs? (Peletier..)**
- **Snaking near the singular limit?**
- **Ferrofluids:**
 - **Matching Experiments and Numerics**
 - **1D/2D pulses interface near onset - normal form?**