

Tensor networks, commuting squares and higher relative commutants of subfactors

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Bi-unitary connections, subfactors and tensor networks

Physicists in condensed matter physics are recently interested in a certain family $(a_{ijkl})_{ijkl}$ of complex numbers labeled with 4 indices, called a **4-tensor**, in connection to **two-dimensional topological order**. They construct certain finite dimensional projections out of this and make physical studies of their ranges.

We first show that their 4-tensor corresponds to a **bi-unitary connection** giving a finite dimensional commuting square, labeled with 4 edges from the 4 Bratteli diagrams. Then our main result identifies the range of their projections with the **higher relative commutants** of the subfactor arising from such a commuting square.

A commuting square

Consider the following inclusions of four finite

dimensional C^* -algebras, $A \subset B$
 $C \subset D$, with a normalized
trace tr on D and $A = B \cap C$.

When the orthogonal projections onto subalgebras B, C with respect to the L^2 -norm arising from the trace commute on D , we say that the above is a **commuting square**. If we have $\text{span } BC = D$, then we say that the commuting square is **non-degenerate**. Finite dimensional non-degenerate commuting squares have been important and well-studied in subfactor theory of Jones over many years.

Repeated basic constructions

Starting with a finite dimensional non-degenerate commuting square, we can repeat **basic constructions** of Jones and get increasing sequences of finite dimensional algebras.

$$\begin{array}{ccccccc} A & \subset & B & \subset & B_1 & \subset & B_2 & \subset & \dots \\ \cap & & \cap & & \cap & & \cap & & \\ C & \subset & D & \subset & D_1 & \subset & D_2 & \subset & \dots \end{array}$$

We take the GNS-completions of the unions $\bigcup_{n=1}^{\infty} B_n \subset \bigcup_{n=1}^{\infty} D_n$ with respect to the trace to get $N \subset M$. Both N and M are **hyperfinite II_1 factors** and we get a **subfactor** of the finite Jones index. We can also repeat the basic construction vertically and get another subfactor $P \subset Q$.

An old question of Jones

When we have a subfactor $N \subset M$ with finite Jones index, we have the **Jones tower**

$N \subset M \subset M_1 \subset M_2 \subset \dots$ arising from the basic constructions. When we have only finitely many irreducible bimodules arising from ${}_N M_k N$, we say that $N \subset M$ has a **finite depth**. This is an important finiteness condition in connection to 3-dimensional topology and mathematical physics.

In 1995, Jones asked the following question.

When one of the two subfactors of $N \subset M$ and $P \subset Q$ has a finite depth, so does the other?

Sato gave a positive answer and a more detailed characterization of the relation between the two.

Strongly amenable subfactors and Popa's classification

From a subfactor $N \subset M$ with finite Jones index, we get the following sequence of commuting squares.

$$\begin{array}{ccccccc} M' \cap M & \subset & M' \cap M_1 & \subset & M' \cap M_2 & \subset & \cdots \\ \cap & & \cap & & \cap & & \\ N' \cap M & \subset & N' \cap M_1 & \subset & N' \cap M_2 & \subset & \cdots \end{array}$$

Popa proved that the subfactor $N \subset M$ is completely recovered from the above commuting squares if the subfactor satisfies a nice analytic property called **strong amenability**. If we have a finite depth and M is hyperfinite, a single commuting square

$$\begin{array}{ccc} M' \cap M_k & \subset & M' \cap M_{k+1} \\ \cap & & \cap \\ N' \cap M_k & \subset & N' \cap M_{k+1} \end{array} \quad \text{for a large } k \text{ suffices.}$$

A bi-unitary connection

We choose one edge each from the four Bratteli diagrams of a commuting square. Then we get an assignment W of a complex number to each such square with the following. This is called a **bi-unitary connection**.

$$\sum_{z, \xi_1, \xi_2} \begin{array}{ccc} x & \xi_4 & y \\ \xi_1 \downarrow & \boxed{W} & \downarrow \xi_3 \\ z & \xi_2 & w \end{array} \quad \overline{\begin{array}{ccc} x & \xi'_4 & y' \\ \xi_1 \downarrow & \boxed{W} & \downarrow \xi'_3 \\ z & \xi_2 & w \end{array}} = \delta_{\xi_3, \xi'_3} \delta_{\xi_4, \xi'_4}$$

$$\begin{array}{ccc} y & \tilde{\xi}_4 & x \\ \xi_3 \downarrow & \boxed{W'} & \downarrow \xi_1 \\ w & \tilde{\xi}_2 & z \end{array} = \sqrt{\frac{\mu_x \mu_w}{\mu_y \mu_z}} \overline{\begin{array}{ccc} x & \xi_4 & y \\ \xi_1 \downarrow & \boxed{W} & \downarrow \xi_3 \\ z & \xi_2 & w \end{array}}$$

Basis change with a bi-unitary connection

Paths of length 2 on two Bratteli diagrams give an orthonormal basis $|\xi_1\xi_2\rangle$ of a (finite dimensional) Hilbert space. Those on the other two Bratteli diagrams give another basis $|\xi_4\xi_3\rangle$ of the same space, and a bi-unitary connection gives a basis change as follows.

$$\xi_1 \begin{array}{c} \text{---} \\ \xi_2 \end{array} = \sum_{\xi_3, \xi_4} \xi_1 \begin{array}{c} \xi_4 \\ \boxed{W} \\ \xi_2 \end{array} \xi_3 \begin{array}{c} \xi_4 \\ \text{---} \\ \xi_3 \end{array}$$

Namely, the bi-unitary connection W gives a unitary matrix $\langle \xi_1\xi_2 | \xi_4\xi_3 \rangle$ on this Hilbert space. This unitarity is a “half” of bi-unitarity.

The string algebra construction

Suppose we have a series of Bratteli diagrams for inclusions $\mathbb{C} = \mathbf{A}_0 \subset \mathbf{A}_1 \subset \mathbf{A}_2 \subset \mathbf{A}_3 \subset \mathbf{A}_4 \subset \dots$.

We have a model for these inclusions as follows. Let (ξ_1, ξ_2) be a pair of path of the same length on this Bratteli diagram with a common starting vertex and a **common ending vertex** at some stage. We call such a pair a **string** and they span a finite dimensional \mathbb{C} -vector space. A string (ξ, η) really means an operator $|\xi\rangle\langle\eta|$ in the **bracket** notation, and this gives an algebra structure among strings of the same length.

We make an embedding of a string (ξ_1, ξ_2) of length k into the next row as $\sum_{\eta} (\xi_1 \cdot \eta, \xi_2 \cdot \eta)$, where η is a path of length 1 and \cdot stands for concatenation of paths.

Construction of a subfactor

We use four Bratteli diagrams and their reflections to obtain doubly indexed string algebras A_{kl} . Since the bi-unitary connection gives a basis change of paths of length 2, it also gives a basis change of strings of length

$$\begin{array}{ccc} A_{kl} & \subset & A_{k,l+1} \\ \cap & & \cap \\ A_{k+1,l} & \subset & A_{k+1,l+1}. \end{array}$$

This is a **commuting square** due to the other “half” of bi-unitarity.

Taking the GNS-completions, we have the limit algebras $A_{k,\infty}$ and $A_{\infty,l}$, and they are **hyperfinite II_1 factors**. We naturally have two **subfactors** $A_{0,\infty} \subset A_{1,\infty}$ and $A_{\infty,0} \subset A_{\infty,1}$, like $N \subset M$ and $P \subset Q$ before.

An example for the Dynkin diagrams

We give an example of a bi-unitary connection as follows. Fix one of the *A-D-E* Dynkin diagram and use it for the four Bratteli diagrams. Let n be its Coxeter number and set $\varepsilon = \sqrt{-1} \exp \frac{\pi \sqrt{-1}}{2(n+1)}$. We write μ_x for the Perron-Frobenius eigenvector entry for a vertex x . Then a bi-unitary connection is given as follows.

$$\begin{array}{ccc} j & & k \\ & \boxed{W} & \\ l & & m \end{array} = \delta_{kl} \varepsilon + \sqrt{\frac{\mu_k \mu_l}{\mu_j \mu_m}} \delta_{jm} \bar{\varepsilon}$$

Figure: A bi-unitary connection on the Dynkin diagram

A subfactor and a fusion category

Suppose the subfactor $A_{0,\infty} \subset A_{1,\infty}$ (or equivalently $A_{\infty,0} \subset A_{\infty,1}$) has a finite depth. Consider the bimodules ${}_{A_{0,\infty}}(A_{k,\infty})_{A_{0,\infty}}$ and their irreducible decompositions. We get only finitely many irreducible bimodules in this way and we have a **fusion category** of bimodules. We have a **relative tensor product** of bimodules and **dual bimodules** there.

We also have corresponding tensor products and irreducible decompositions at the level of bi-unitary connections. We then have an equivalent fusion category of bi-unitary connections. This correspondence is given by the **open string bimodule** construction, due to Asaeda-Haagerup.

A 4-tensor from a bi-unitary connection

Suppose we have a bi-unitary connection W_a . We then define a 4-tensor a as follows.

$$\begin{array}{c} \xi_6 \cdot \xi_5 \\ \xi_1 \text{---} \textcircled{a} \text{---} \xi_4 \\ \xi_2 \cdot \xi_3 \end{array} = \sqrt[4]{\frac{\mu_x \mu_w}{\mu_y \mu_z}} \begin{array}{c} x \quad \xi_6 \quad \xi_5 \quad y \\ \xi_1 \quad \boxed{W_a \quad W'_a} \quad \xi_4 \\ z \quad \xi_2 \quad \xi_3 \quad w \end{array}$$

Here W'_a stands for the horizontal **reflection** of W_a . We also use the vertical reflection so that we can concatenate 4-tensors as usual. The reflection corresponds to **basic construction** and the vertical concatenation of 4-tensors corresponds to the product of bi-unitary connections.

A matrix product operator algebra

Suppose we have a 4-tensor corresponding to a commuting square giving a subfactor of finite depth.

Bultinck-Mariën-Williamson-Şahinoğlu-Haegeman-Verstraete gave an anyon algebra, a finite dimensional C^* -algebra, in this setting and argued that its minimal central projections give **anyons** describing a **two-dimensional topological order**. Here an anyon is a new type of quasi-particle more general than a boson and a fermion and it is expected to be useful for constructing a topological quantum computer.

We proved that this anyon algebra is isomorphic to the **tube algebra** of Ocneanu and anyons correspond to irreducible objects of the **Drinfel'd center**.

A projector matrix product operator

We define a **matrix product operator** O_a^k as follows.

$$\sum \left(\begin{array}{c} \xi_1 \quad \xi_2 \quad \dots \quad \xi_k \\ | \\ \textcircled{a} \text{---} \textcircled{a} \text{---} \dots \text{---} \textcircled{a} \\ | \\ \eta_1 \quad \eta_2 \quad \dots \quad \eta_k \end{array} \right) | \xi_1 \xi_2 \dots \xi_k \rangle \langle \eta_1 \eta_2 \dots \eta_k |$$

We then set $P^k = \sum_a \frac{d_a}{w} O_a^k$ like Bultinck-Mariën-

Williamson-Şahinoğlu-Haegeman-Verstraete. This is a **projector matrix product operator (PMPO)** and it acts on certain **projected entangled pair state (PEPS)**.

Higher relative commutants of a subfactor

The range of the projector matrix product operator P^k plays an important role in theory of two-dimensional topological order, and we identify it with the **higher relative commutant** $A'_{\infty,0} \cap A_{\infty,k}$ of the subfactor. This is equal to $A_{0,k}$ if (and only if) the original bi-unitary connection is **flat**, but we do **not** assume this flatness here.

We have the inclusion $A'_{\infty,0} \cap A_{\infty,k} \subset A_{0,k}$ due to Ocneanu's **compactness argument** and he proved that an element in $A'_{\infty,0} \cap A_{\infty,k}$ is characterized as a **flat field of strings** of length k . A field of strings is an element in a certain string algebra and it is flat if and only if it does not change the form under **parallel transport** of length 2.

A sketch of a proof

We sketch a proof of the above identification.

It is not difficult to show that if we have a flat field of strings, then it is preserved under the projector matrix product operator P^k because a flat field does not change the form under a parallel transport.

Conversely, take an element in the range of the projector matrix product operator P^k . Then we construct an element $x_m \in A'_{m,0} \cap A_{m,k}$ in a simple manner. Using the Perron-Frobenius theorem, we show that $\{x_m\}_m$ is a Cauchy sequence in the L^2 -norm, so it converges to some x in $A'_{\infty,0} \cap A_{\infty,k}$ and gives a **flat field of strings**. We next show that all x_m are actually equal to x .

The above two maps are actually mutual inverses.

The Drinfel'd center and Morita equivalence

For getting a fusion category, we used the subfactor $A_{0,\infty} \subset A_{1,\infty}$, but now for the range of the projector matrix product operator, we used the higher relative commutants of the **other** subfactor $A_{\infty,0} \subset A_{\infty,1}$. The former is used to get a modular tensor category through the tube algebra and we have description of **anyons**. The latter produces a series of Hilbert spaces on which Hamiltonians act.

These two subfactors can be quite different, but still the relation between the two is characterized as being **opposite Morita equivalent**. In particular, they produce complex conjugate **topological quantum field theory** (TQFT) and have the same **Drinfel'd center**.