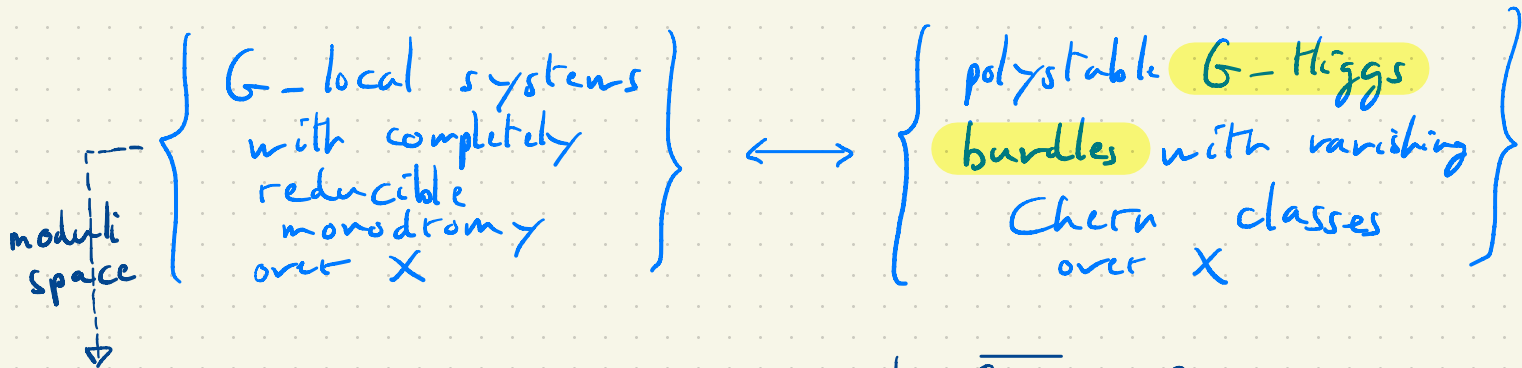


3. Higgs bundles for non-constant groups

a). Non-Abelian Hodge correspondence

Non-Abelian Hodge theory sets up a correspondence



$\text{Hom}(\pi_1 X; G) // G$: recall that $\overline{G \cdot \rho} = G \cdot \rho$
iff $\overline{\rho(\pi_1 X)}^{\text{zar}}$ is a reductive subgroup of G
Affine variety / G

X : nonsingular complex projective variety

G : complex reductive affine algebraic group

Equivalently, it is a correspondence

$$\left\{ \begin{array}{l} G\text{-bundles with} \\ \text{integrable connection} \\ \text{and completely reducible} \\ \text{monodromy over } X \end{array} \right\} \longleftrightarrow \left\{ \begin{array}{l} \text{polystable } G\text{-Higgs} \\ \text{bundles with vanishing} \\ \text{Chern classes} \\ \text{over } X \end{array} \right\}$$

$$(E, \mathcal{D}) \longleftrightarrow (\mathcal{E}, \Theta)$$

Henceforth, we shall focus on the case $G = GL(r; \mathbb{C})$.

ppal $GL(r; \mathbb{C})$ -bundle \longleftrightarrow rank r vector bundle
 locally free \mathcal{O}_X -module

Connection: $\mathcal{D}: E \rightarrow \Omega_X^1 \otimes E$ such that
 $\mathcal{D}(F\sigma) = dF \otimes \sigma + F \mathcal{D}\sigma$

Higgs Field: $\Theta: \mathcal{E} \rightarrow \Omega_X^1 \otimes \mathcal{E}$ such that $\Theta(F\sigma) = F\Theta(\sigma)$

Observations

(i) if $\text{rk}(\mathcal{E}) = 1$, a Higgs field is just a (holomorphic) 1-form

$$\theta \in \Omega_X^1 \otimes \underbrace{\text{End}(\mathcal{E})}_{\mathcal{E}^* \otimes \mathcal{E}}$$

$\mathcal{E}^* \otimes \mathcal{E} \cong \mathcal{O}_X$ if \mathcal{E} is a line bundle

(ii) Rank 1 local systems of vector spaces
Flat line bundles are parameterized by

$$H^1(X; \underline{\mathbb{C}}_X^*)$$

If $\dim_{\mathbb{C}} X = 1$, we have $(X \text{ compact, connected})$

$$H^1(X; \underline{\mathbb{C}}_X^*) \cong_{\text{hom-co}} \text{Hom}(\pi_1 X; \mathbb{C}^*) \cong (\mathbb{C}^*)^{2g}$$

where $g = \text{genus of } X$.

Exponential exact sequence

$$(\dim_{\mathbb{C}} X = 1)$$

$$\begin{array}{ccccccc} 0 & \longrightarrow & \underline{\mathbb{Z}}_X & \longrightarrow & \underline{\mathbb{C}}_X & \xrightarrow{\exp} & \underline{\mathbb{C}}_X^* & \longrightarrow & 1 \\ & & \parallel & & \downarrow & & \downarrow & & \\ 0 & \longrightarrow & \underline{\mathbb{Z}}_X & \longrightarrow & \mathcal{O}_X & \xrightarrow{\exp} & \mathcal{O}_X^* & \longrightarrow & 1 \end{array}$$

$$f: U \rightarrow \mathbb{C}$$

$$\begin{array}{ccc} \exp(f) : U & \longrightarrow & \mathbb{C}^* \\ z & \longmapsto & e^{f(z)} \end{array}$$

induces

$$\begin{array}{ccccccc} 0 & \longrightarrow & H^1(X; \underline{\mathbb{Z}}_X) & \longrightarrow & H^1(X; \underline{\mathbb{C}}_X) & \longrightarrow & H^1(X; \underline{\mathbb{C}}_X^*) & \xrightarrow{0} & H^2(X; \underline{\mathbb{Z}}_X) \\ & & \parallel & & \downarrow & & \downarrow & & \parallel \\ 0 & \longrightarrow & H^1(X; \underline{\mathbb{Z}}_X) & \longrightarrow & H^1(X; \mathcal{O}_X) & \longrightarrow & \underbrace{H^1(X; \mathcal{O}_X^*)}_{\text{holomorphic line bundles on } X} & \xrightarrow{\deg} & \underbrace{H^2(X; \underline{\mathbb{Z}}_X)}_{\simeq \mathbb{Z}} \end{array}$$

flat line bundles
on X

holomorphic
line bundles on X

Line bundles of degree 0

The exactness of

$$0 \rightarrow H^1(X; \underline{\mathbb{Z}}_X) \rightarrow H^1(X; \mathcal{O}_X) \rightarrow H^1(X; \mathcal{O}_X^*) \xrightarrow{\deg} H^2(X; \underline{\mathbb{Z}}_X)$$

implies that

$$\text{Jac}(X) := \deg^{-1}(0) \simeq \frac{H^1(X; \mathcal{O}_X)}{H^2(X; \underline{\mathbb{Z}}_X)} \simeq \frac{\mathbb{C}^g}{\mathbb{Z}^{2g}}$$

Likewise, $H^2(X; \underline{\mathbb{C}}_X^*) \simeq \frac{H^1(X; \underline{\mathbb{C}}_X)}{H^2(X; \underline{\mathbb{Z}}_X)}$.

Harmonic forms

We seek to understand the map

$$H^2(X; \underline{\mathbb{C}}_X)$$



$$H^2(X; \mathcal{O}_X)$$

$$\frac{\{d\text{-closed } 1\text{-forms}\}}{\{d\text{-exact } 1\text{-forms}\}}$$

$$H^1(X; \underline{\mathbb{C}}_X) \simeq \underset{\text{De Rham}}{H^1_{dR}(X; \mathbb{C})} \simeq \underset{\text{Hodge}}{\text{Harm}^1_{\Delta_d}(X)} = \left\{ \omega \in \Omega^1_{\mathbb{C}}(X) \mid \Delta_d \omega = 0 \right\}$$

$\hookrightarrow dd^* + d^*d$

locally
constant

$$\frac{\{\bar{\partial}\text{-closed } (0,1)\text{-forms}\}}{\{\bar{\partial}\text{-exact } (0,1)\text{-forms}\}}$$

$$H^1(X; \mathcal{O}_X) \simeq \underset{\text{Dolbeault}}{H^{0,1}_{\text{Dol}}(X)} \simeq \underset{\text{Hodge}}{\text{Harm}^{0,1}_{\Delta_{\bar{\partial}}}(X)} = \left\{ \beta \in \Omega^{0,1}_{\mathbb{C}}(X) \mid \Delta_{\bar{\partial}} \beta = 0 \right\}$$

$\hookrightarrow \bar{\partial}\bar{\partial}^* + \bar{\partial}^*\bar{\partial}$

coherent

Kähler identities and the Hodge decomposition theorem

Recall that $d = \partial + \bar{\partial}$, $= d^{1,0} + d^{0,1}$ (for 1-forms)

and that $\Omega_{\mathbb{C}}^1(X; \mathbb{C}) = \Omega_{\mathbb{C}}^{1,0}(X; \mathbb{C}) \oplus \Omega_{\mathbb{C}}^{0,1}(X; \mathbb{C})$.

The **Kähler identities** on X imply that

$$\Delta_d = 2 \Delta_{\bar{\partial}}$$

In particular $\omega = \underbrace{\alpha}_{(1,0)} + \underbrace{\beta}_{(0,1)}$ satisfies $\Delta_d = 0$ iff $\Delta_{\bar{\partial}} \alpha = 0$ and $\Delta_{\bar{\partial}} \beta = 0$

Hodge decomposition:

$$\begin{aligned} H^1(X; \underline{\mathbb{C}}_X) &\simeq H_{dR}^1(X; \mathbb{C}) \simeq H^{1,0}(X) \oplus H^{0,1}(X) \\ &\downarrow \text{projection} \\ H^1(X; \mathcal{O}_X) &\simeq H^0(X; \Omega_X^1) \oplus H^1(X; \mathcal{O}_X) \end{aligned}$$

X Kähler
 \updownarrow

The Abelian case of the NAHC (in dimension 1)

From the Hodge decomposition, one gets:

$$\begin{aligned}
 H^1(X; \underline{\mathbb{C}}_X^*) &\simeq \frac{H^1(X; \underline{\mathbb{C}}_X)}{H^1(X; \underline{\mathbb{R}}_X)} \stackrel{\text{homeo}}{\simeq} H^0(X; \Omega_X^1) \oplus \frac{H^1(X; \underline{\mathbb{O}}_X)}{H^1(X; \underline{\mathbb{R}}_X)} \\
 &\simeq \text{Jac}(X) \times H^0(X; \Omega_X^1)
 \end{aligned}$$

holomorphic 1-form

c.e.

a correspondence

$$\left\{ \text{flat line bundles } (L, \nabla) \right\} \leftrightarrow \left\{ \text{rank 1 Higgs bundles } (Z, \theta) \right\}$$

$$\begin{aligned}
 &\stackrel{\text{homeo}}{\simeq} (\mathbb{C}^*)^{2g} \\
 &\quad \downarrow \\
 &S_1 \times \mathbb{R} \quad (\text{polar decomposition})
 \end{aligned}$$

$$\begin{aligned}
 &\frac{\mathbb{C}^g}{\mathbb{Z}^{2g}} \times \mathbb{C}^g \\
 &\stackrel{\text{homeo}}{\simeq} S_1^{2g} \times \mathbb{R}^{2g}
 \end{aligned}$$

Observation

The rank 1 case of the NHT of a curve X yields a homeomorphism:

$$\underbrace{(\mathbb{C}^*)^2}_{\text{affine algebraic variety (Betti space)}} \simeq \text{Hom}(\Pi_1 X; \mathbb{C}^*) \underset{\text{homeo}}{\simeq} \left\{ \begin{array}{l} \text{rank 1 degree 0} \\ \text{Higgs bundles} \end{array} (Z, \theta) \right\}$$

$$\simeq \underbrace{\text{Jac}(X) \times H^0(X; \Omega_X^1)}$$

---> these two moduli spaces are homeomorphic but not isomorphic as algebraic varieties or complex analytic manifolds

this is the cotangent bundle of $\text{Jac}(X)$
(Dolbeault space)

The higher rank case Building on work by Kobayashi,

Hitchin, Donaldson, Corlette and Simpson
 $r=2, \dim_{\mathbb{C}} X = 1$ $r, \dim X$ arbitrary

have generalized Hodge theory to the
higher rank case.

$$\mathbb{C}^* \xrightarrow[\text{replaced by}]{} GL(r; \mathbb{C})$$

$$\left\{ \begin{array}{l} \text{Flat vector bundles} \\ \text{of rank } r \end{array} \right\} \xleftrightarrow{?} \left\{ \begin{array}{l} \text{semistable Higgs bundles} \\ \text{of rank } r \text{ \& degree } 0 \end{array} \right\}$$

Harmonic bundles

A harmonic bundle is a quadruple (E, h, A, ψ)

where:

- $E \rightarrow X$ is a C^∞ complex vector bundle
- h is a Hermitian metric on E
- A is a unitary connection on (E, h)
- $\psi \in \Omega_{C^\infty}^1(X; \text{Herm}(E, h))$

such that

Hitchin equations

$$\left\{ \begin{array}{l} F_A + \frac{1}{2} [\psi \wedge \psi] = 0 \\ d_A \psi = 0 \\ d_A^* \psi = 0 \end{array} \right.$$

From harmonic bundles to Higgs and flat bundles

{ harmonic bundles }
(E, h, A, ψ)

$\nabla = A + \psi$

{ Flat vector bundles }
(E, D)

$$\left. \begin{aligned} F_A + \frac{1}{2}[\psi, \psi] = 0 \\ d_A \psi = 0 \end{aligned} \right\} \Leftrightarrow F_D = 0$$

$d_A^* \psi = 0 \Rightarrow (E, D) \text{ has } \underline{\text{completely reducible holonomy}}$

$\Sigma = (E, k_A^{0,1})$
 $\Theta = \psi^{1,0}$

{ Higgs bundles with vanishing Chern classes }
(Σ, Θ)

$$\left. \begin{aligned} k_A \psi = 0 \\ d_A^* \psi = 0 \end{aligned} \right\} \Leftrightarrow \left\{ \begin{aligned} d_A^{0,1} \Theta = 0 \\ \Theta \wedge \Theta = 0 \end{aligned} \right.$$

$$\left. \begin{aligned} F_A + \frac{1}{2}[\psi, \psi] = 0 \\ \Leftrightarrow F_A + [\Theta \wedge \Theta^*] = 0 \end{aligned} \right\} \Rightarrow (\Sigma, \Theta) \text{ is } \underline{\text{polystable}}$$

Harmonic metrics

(i) on a flat bundle (E, \mathcal{D}) :

a metric h induces a decomposition

$$\mathcal{D} = A + \psi$$

with A h -unitary
and ψ Hermitian

(satisfying $F_A + \frac{1}{2} [\psi, \psi] = 0$
and $d_A \psi = 0$)

Theorem (Donaldson - Corlette)

If \mathcal{D} has **completely reducible holonomy**,

then (E, \mathcal{D}) admits h such that $d_A^{*h} \psi = 0$.

c.c. (E, \mathcal{D}) comes from a harmonic bundle.

Harmonic metrics

(ii) on a Higgs bundle (E, θ) :

Since Σ is holomorphic, a metric h determines a Chern connection A , which is unitary, and a Hermitian 1-form $\psi := \theta + \theta^{*h}$ (satisfying $d_A \psi = 0$ and $d_A^{*h} \psi = 0$).

Theorem (Hitchin-Simpson)

If $(E, \theta) \simeq (E_1, \theta_1) \oplus \dots \oplus (E_k, \theta_k)$ with each (E_i, θ_i) stable with vanishing Chern classes, then (E, θ) admits h such that $F_A + [\theta \wedge \theta^{*h}] = 0$
i.e. (E, θ) comes from a harmonic bundle.

Stability for degree 0 Higgs bundles

$$(\dim_{\mathbb{C}} X = 1)$$

Let (\mathcal{E}, θ) be a degree 0 Higgs bundle.

Then (\mathcal{E}, θ) is called stable if

$\forall (0 \subsetneq \mathcal{F} \subsetneq \mathcal{E})$ such that $\theta(\mathcal{F}) \subset \Omega_X^1 \otimes \mathcal{F}$,
one has $\deg \mathcal{F} < 0$.

Examples

(i) rank 1 Higgs bundles (\mathcal{L}, θ)

(ii) $\mathcal{E} = \Omega_X^{-1/2} \oplus \Omega_X^{1/2}$, $\theta = \begin{pmatrix} 0 & q_2 \\ 1 & 0 \end{pmatrix}$

$$\theta: \Omega_X^{1/2} \oplus \boxed{\Omega_X^{-1/2}} \rightarrow \boxed{\Omega_X^{3/2}} \oplus \Omega_X^{1/2} \quad q_2 \in \Gamma(\text{Hom}(\Omega_X^{-1/2}, \Omega_X^{3/2})) \\ \cong \Omega_X^{\otimes 2}$$

b). The non-constant case

G : complex reductive group

A principal G -Higgs bundle
is a pair (\mathcal{E}, Θ)

where:

- \mathcal{E} is a principal G -bundle
- $\Theta \in H^0(X; \Omega_X^1 \otimes \text{ad}(\mathcal{E}))$
such that $[\partial_n \Theta] = 0$

$$\left(\text{ad}(\mathcal{E}) = \left(\mathcal{E} \times \text{Lie}(G) \right) / G \right)$$

Example $G = GL(n, \mathbb{C})$ $\text{ad}(\mathcal{E}) \cong \text{End}(V)$
 $\mathcal{E} \cong V \rightarrow$ classical Higgs bundle

$\mathfrak{g} \rightarrow X$ a group bundle on X

A \mathfrak{g} -Higgs torsor
is a pair (\mathcal{E}, Θ)

where:

- \mathcal{E} is a \mathfrak{g} -torsor
- $\Theta \in H^0(X; \Omega_X^1 \otimes \text{ad}(\mathcal{E}))$
such that $-\partial_n \Theta = 0$

$$\left(\text{ad}(\mathcal{E}) = \left(\mathcal{E} \times_{\mathfrak{g}} \text{Lie}(\mathfrak{g}) \right) / \mathfrak{g} \right)$$

Example $\mathfrak{g} = X \times G \rightarrow$ principal G -Higgs bundles

Relation to equivariant principal bundles

Assume that $\mathcal{G} = \mathcal{G}_\rho := (\tilde{X} \times G) / \pi_1 X$ For $\varphi: \pi_1 X \rightarrow \text{Aut}(G)$.

[the only case in which we can have
a correspondence with twisted local systems]

Define a cover $\underbrace{X_\rho \xrightarrow{p} X}$ by $\pi_1 X_\rho := \ker \varphi \triangleleft \pi_1 X$.

Galois cover, finite if $F := \text{Im } \varphi$ is finite

$p^* \mathcal{G}_\rho = G_{X_\rho} = X_\rho \times G$ because $\pi_1 X_\rho \cong \text{Aut}_X(X_\rho)$

$\pi_1 X_\rho$ acts trivially on G

\mathcal{G}_ρ -torsors (on X) \longleftrightarrow F -equivariant principal G -bundles

Equivariant Higgs Fields

(E, θ) : a \mathfrak{g}_T -Higgs torsor

$E := \rho^* E$ has an F -equivariant structure:

$$\begin{array}{ccc} E & \xrightarrow{\tau_F} & E \\ \downarrow & & \downarrow \\ X_T & \xrightarrow{F} & X_T \end{array}$$

$$\tau_{1_F} = \text{id}_E$$

$$\tau_{f_1 \circ f_2} = \tau_{f_1} \circ \tau_{f_2}$$

$$\tau_f(m \cdot g) = \tau_f(m) \cdot \varphi_f(g)$$

$$\theta \in H^0(X; \Omega_X^{-1} \otimes \text{ad}(E))$$

$$\begin{array}{c} \updownarrow \\ \tilde{\theta} \in \text{Fix}_F H^0(X_T; \Omega_{X_T}^{-1} \otimes \text{ad}(E)) \end{array}$$

$$\begin{array}{ccc} TX_T & \xrightarrow{\tilde{\theta}} & \text{ad}(E) = (E \times \text{Lie}(G)) / G \ni [m, \xi] \\ \downarrow F & & \downarrow \tau_F \quad \downarrow \\ TX_T & \xrightarrow{\tilde{\theta}} & \text{ad}(E) = (E \times \text{Lie}(G)) / G \ni [\tau_F(m), \rho_F(\xi)] \end{array}$$

Equivariant NAHC

$\left\{ \begin{array}{l} \text{F-equivariant} \\ \text{harmonic bundles} \\ (E, h, A, \psi, \tau) \end{array} \right\}$

$\swarrow \mathbb{D} = A + \psi$

$\left\{ \begin{array}{l} \text{F-equivariant} \\ \text{flat vector bundles} \\ (E, \mathbb{D}, \tau) \end{array} \right\}$

$\searrow \begin{array}{l} \Sigma = (E, L_A^{0,1}) \\ \Theta = \psi^{1,0} \end{array}$

$\left\{ \begin{array}{l} \text{F-equivariant} \\ \text{Higgs bundles with} \\ \text{vanishing Chern classes} \\ (\Sigma, \Theta, \tau) \end{array} \right\}$

\leadsto study the existence of **F-invariant solutions** to Hitchin's equations

Stability

(à la Ramanan, when $\dim_{\mathbb{C}} X = 1$)

An F -equivariant principal G -Higgs bundle (E, θ, τ) on X_{φ} is called stable if, for all F -invariant

parabolic subgroup $P \subset G$ and all F -equivariant

reduction of structure group $s: X_{\varphi} \rightarrow E/P$

such that $\theta(T_{X_{\varphi}}) \subset \text{ad}(\underbrace{s^*E}_{=: E_P})$,

one has

$$\deg \text{ad}(E_P) < 0.$$

→ OK in all our examples

Example (no stability issue here)

Recall the case of a Riemann surface with involution (Y, σ) : σ acts on \mathbb{C}^* via $z \mapsto z^{-1}$ and we can look at twisted local systems for the non-constant group $\mathfrak{g}_\sigma := (\mathbb{C}^* \times \mathbb{C}^*) / \langle \sigma \rangle$ on $X := Y / \langle \sigma \rangle$.

[anti-invariant rank 1 local systems]

NAHC on Y : $\left\{ \begin{array}{l} \text{flat } \mathbb{C}^* \text{-bundles} \\ \text{on } Y \end{array} \right\} \leftrightarrow \left\{ \begin{array}{l} \text{rank 1 line bundles of degree 0} \\ \text{with a holomorphic 1-form} \end{array} \right\}$

equivariant version $\left\{ \begin{array}{l} \text{anti-invariant rank 1} \\ \text{local systems on } Y \end{array} \right\} \leftrightarrow \left\{ \begin{array}{l} \text{anti-invariant} \\ \text{line bundle} \\ (\sigma^* L)^* = L \end{array} \right\} \times \left\{ \begin{array}{l} \text{1-form } \omega \\ \text{such that} \\ \sigma^* \omega = -\omega \end{array} \right\}$

σ -twisted \mathbb{C}^* -local systems on X \mathfrak{g}_σ -Higgs torsors on X

Observations

(i) g_σ -Higgs torsors are not classical Higgs bundles on X (as opposed to the case of **invariant** local systems on Y , giving rise to local systems and classical Higgs bundles on X)

(ii) The NAHC on $X/\langle \sigma \rangle$ yields a homeomorphism

$$\left[\text{Hom}_\sigma(\pi_1 X; \mathbb{C}^* \langle \sigma \rangle) / \mathbb{C}^* \right] \cong T^* \text{Ban}_{g_\sigma} = T^* \text{Jac}(X) / \langle \sigma \rangle$$

which is a twisted version of

$$\text{Hom}(\pi_1 X; \mathbb{C}^*) \cong T^* \text{Jac}(X).$$

(iii) When the space X is allowed to be an orbifold, then the total space of the cover $X_\rho \rightarrow X$ is not a manifold in general. So, to apply the equivariant version of the NAHC, one needs to pass to a higher degree cover, which is not canonical in general.

Applications of this point of view include the study of Hitchin components for cocompact Fuchsian groups $\Gamma \subset \mathrm{PGL}(2, \mathbb{R})$.